Introduction to hadronic collisions: theoretical concepts and practical tools for the LHC

Lecture 4

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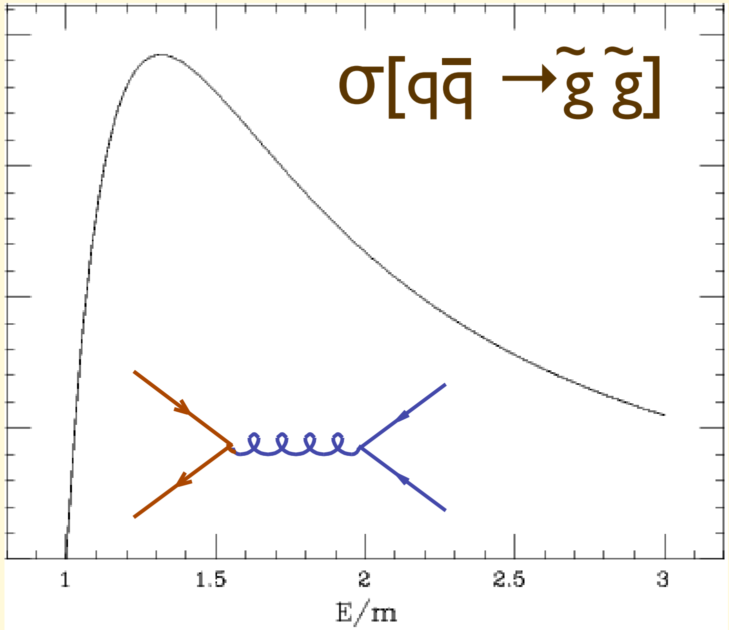
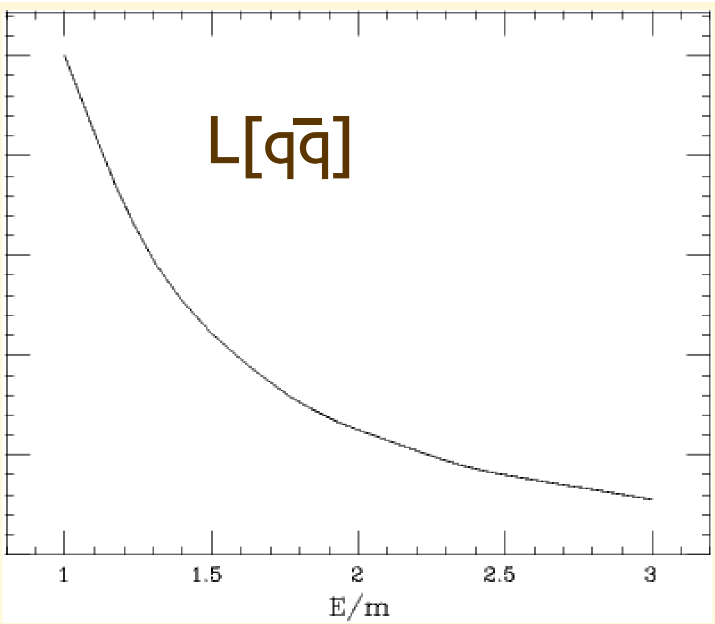
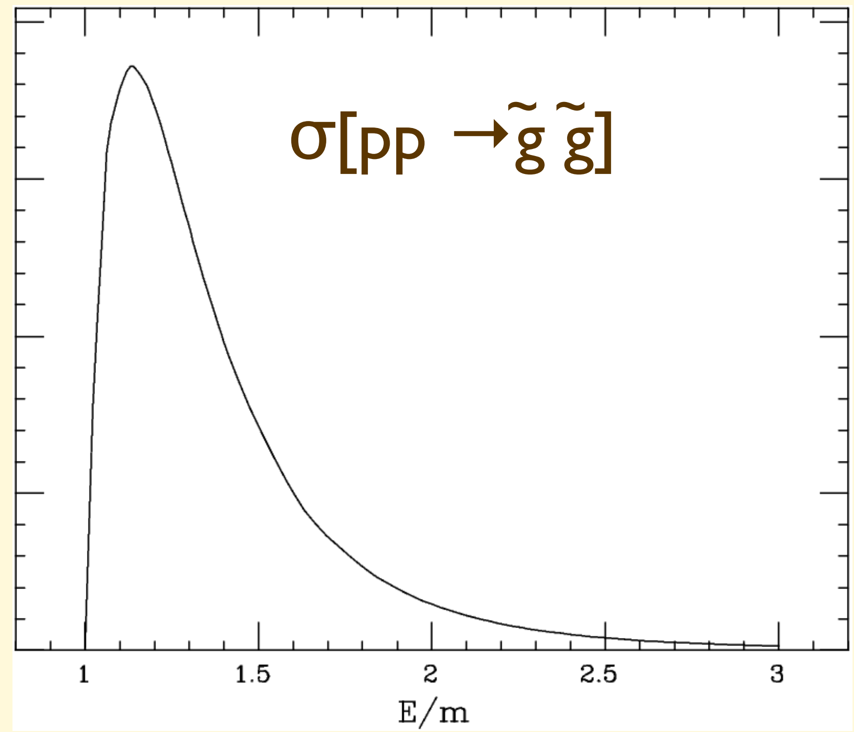
Ex: Gluino pair production

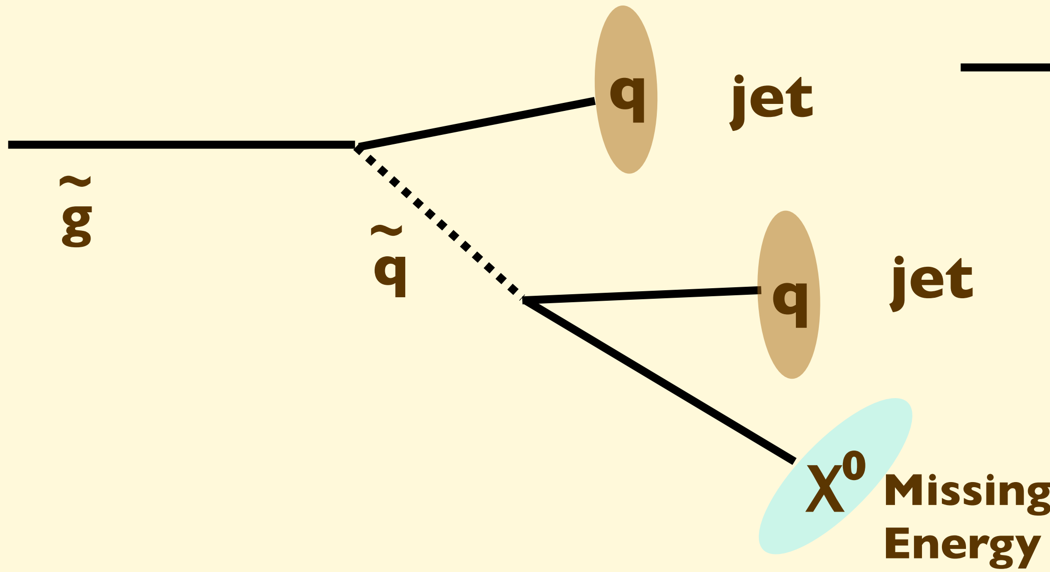


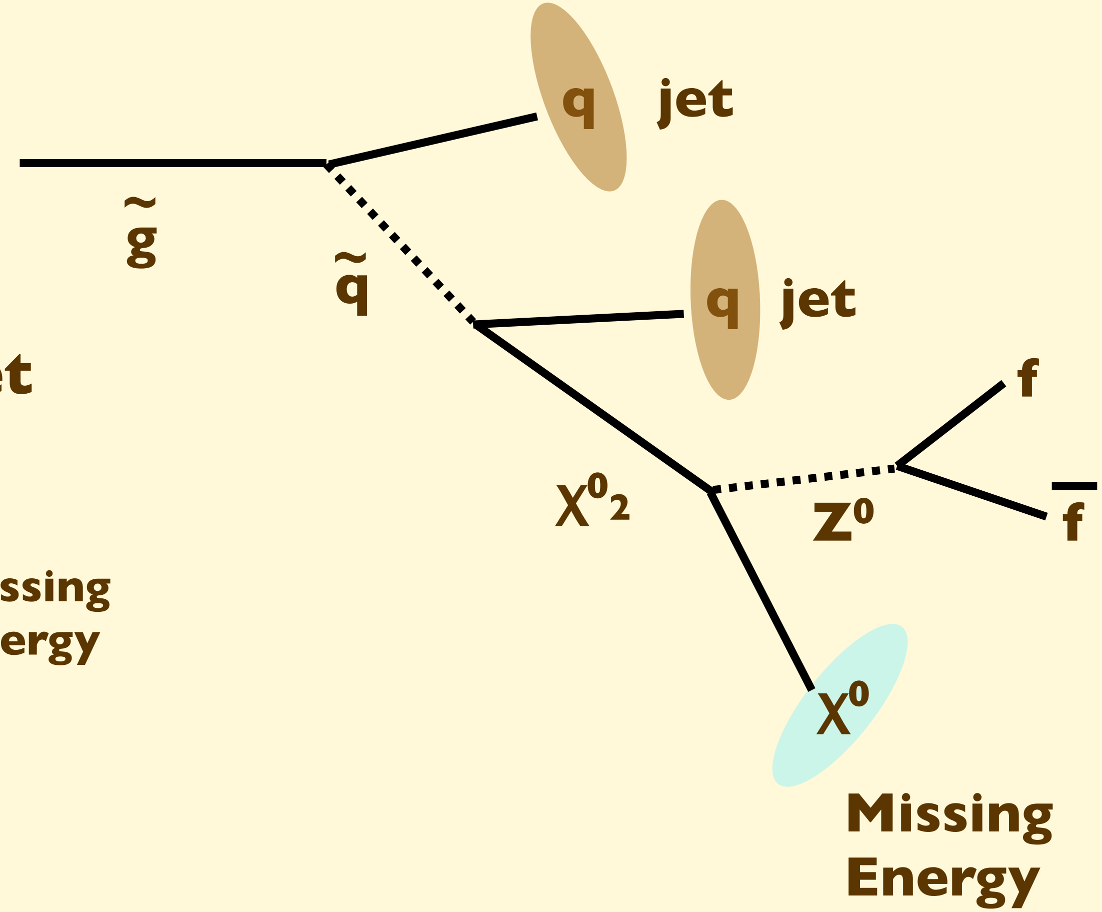


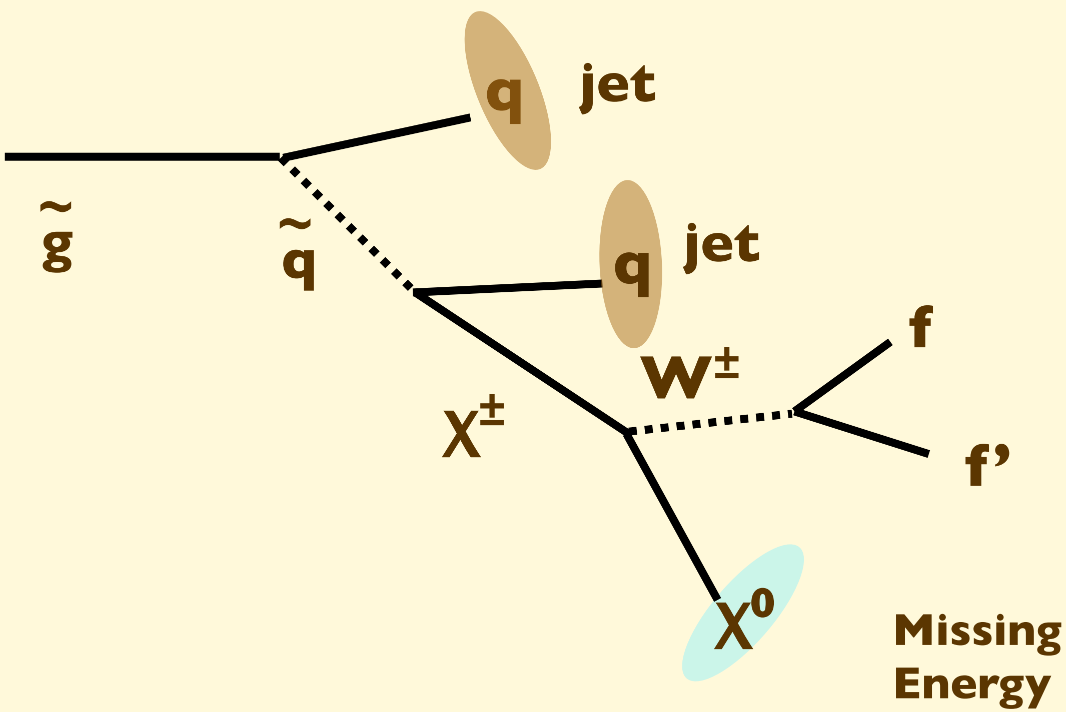


⇒ slow gluinos, $\beta \sim 0.5$

Gluino final states

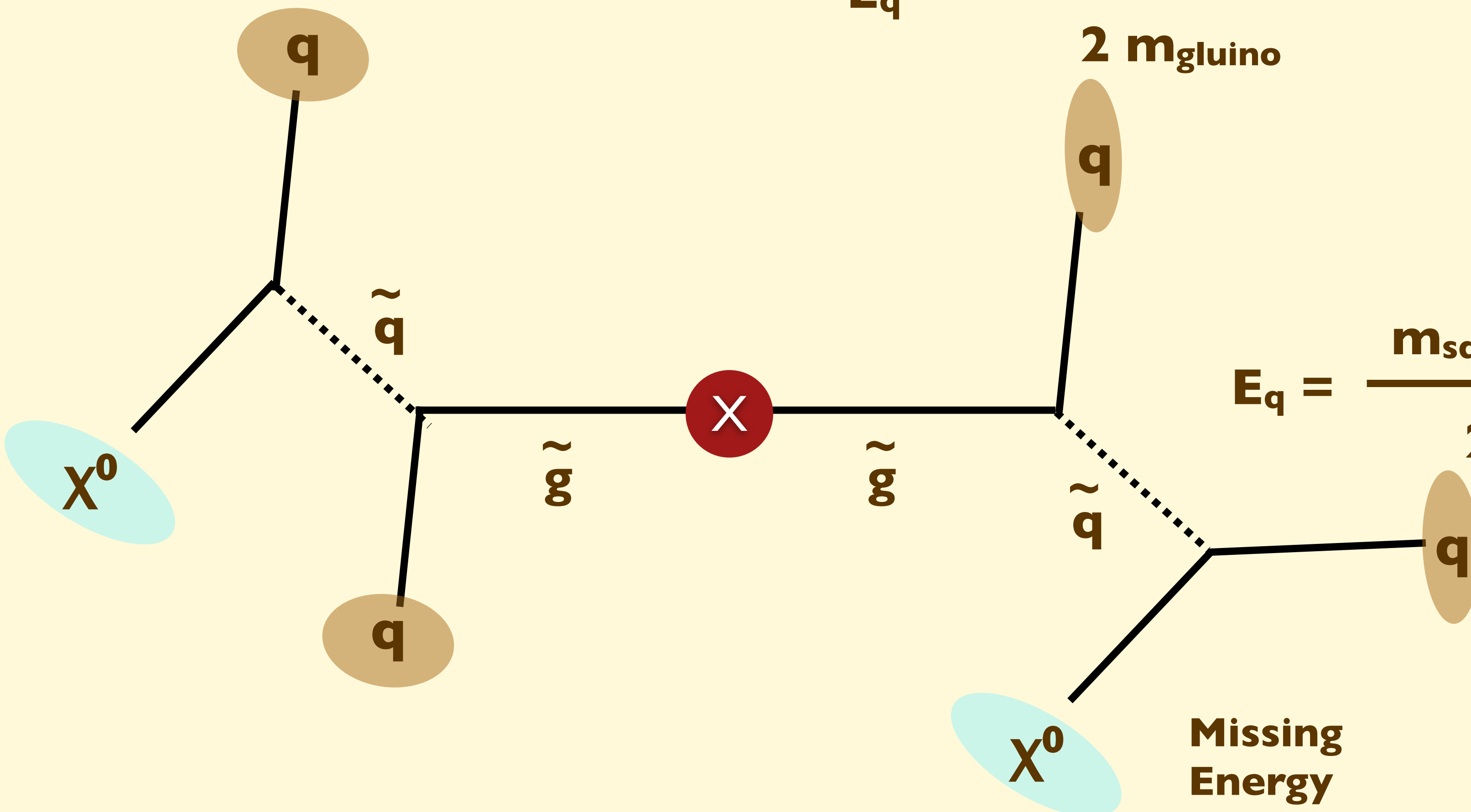






$\tilde{g} \rightarrow \geq 2$ **jet + MET +** $\geq 0\ \ell^{\pm}$

$\tilde{g}\,\tilde{g} \rightarrow$ 4 jet + MET



$$E_q = \frac{m_{gluino}^2 - m_{squark}^2}{2\, m_{gluino}}$$

$$E_q = \frac{m_{squark}^2 - m_X^2}{2\, m_{squark}}$$

Widely-spaced jets, no significant hierarchy in transverse energies and missing E_T

Typical analysis cuts (ATLAS):

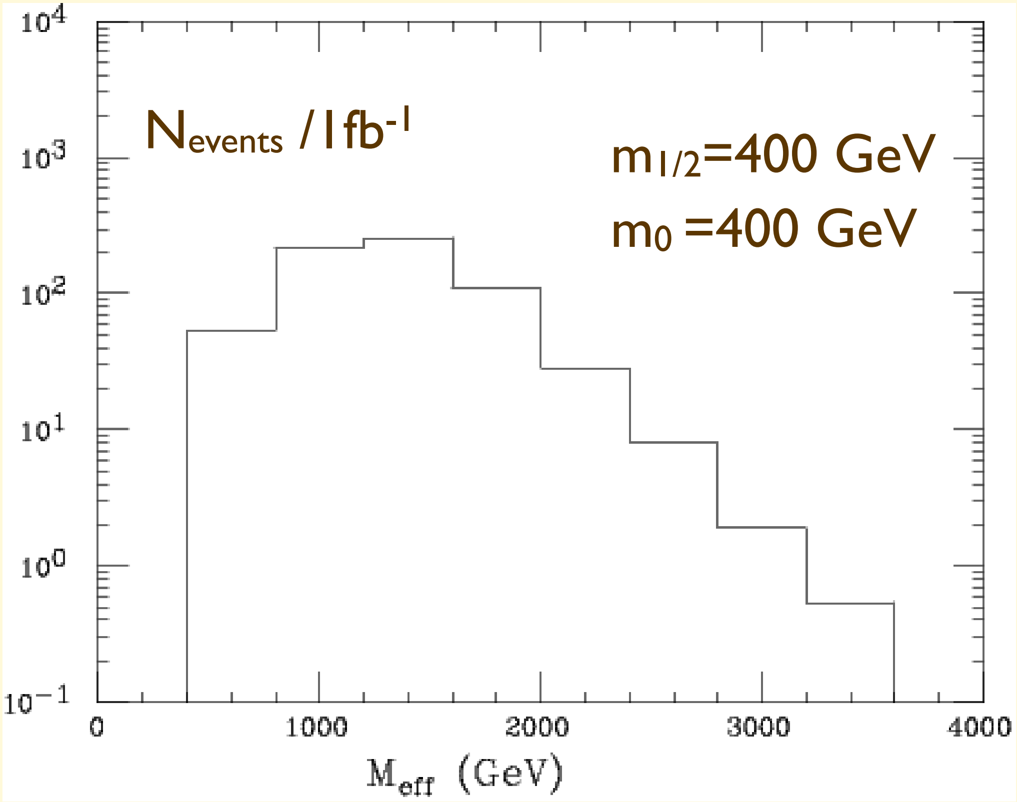
$\geq$4jets, E_T>50 GeV leading jet E_T>100 GeV

no lepton with E_T>20 GeV

MissET> max(100, 0.2 M_{eff})

M_{eff} = MET + $\sum_{i=1,..,4} E_T^i$

Transverse sphericity > 0.2



SM Backgrounds

Missing energy ⇒ νs ⇒ W/Z production

"Irreducible": individual events cannot be distinguished from the signal

Z+4jets, Z→ νν

"Reducible": individual events feature properties which distinguish them from the signal, but these can only be exploited with limited efficiency

W+3jets, W→ τν, τ→hadrons (jet)

τ jet has low multiplicity, and originates from a displaced vertex, because of τs lifetime

W+4jets, W→ e/μ ν, lepton undetected

e/μ can be detected, but cannot be vetoed with 100% efficiency, else the signal would be killed as well (e /μ may come from π conversions or decays)

tt → W+jets, with W→ leptons as above

In addition to the above, top decays have b's, but these cannot be detected and vetoed with 100% efficiency

“Instrumental”: individual events resemble the signal because of instrumental “effects” (namely instrumental deficiencies)

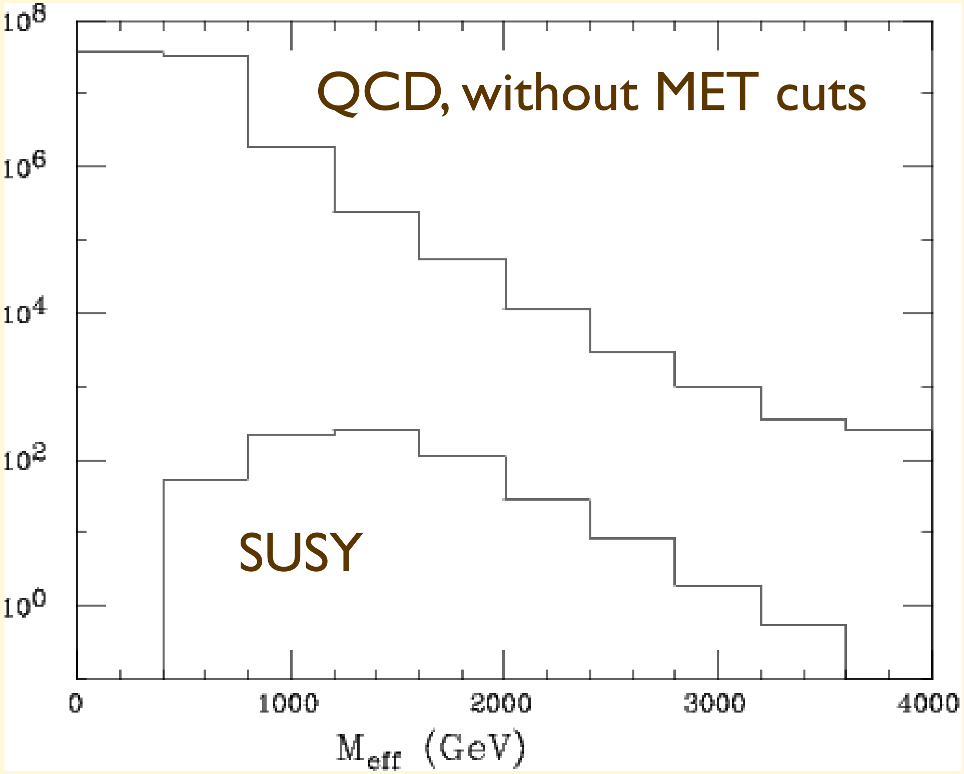
Multijets

The missing ET may originate from several sources:

Mismeasurement of the energy of individual jets

Incomplete coverage in rapidity (forward jets undetected)

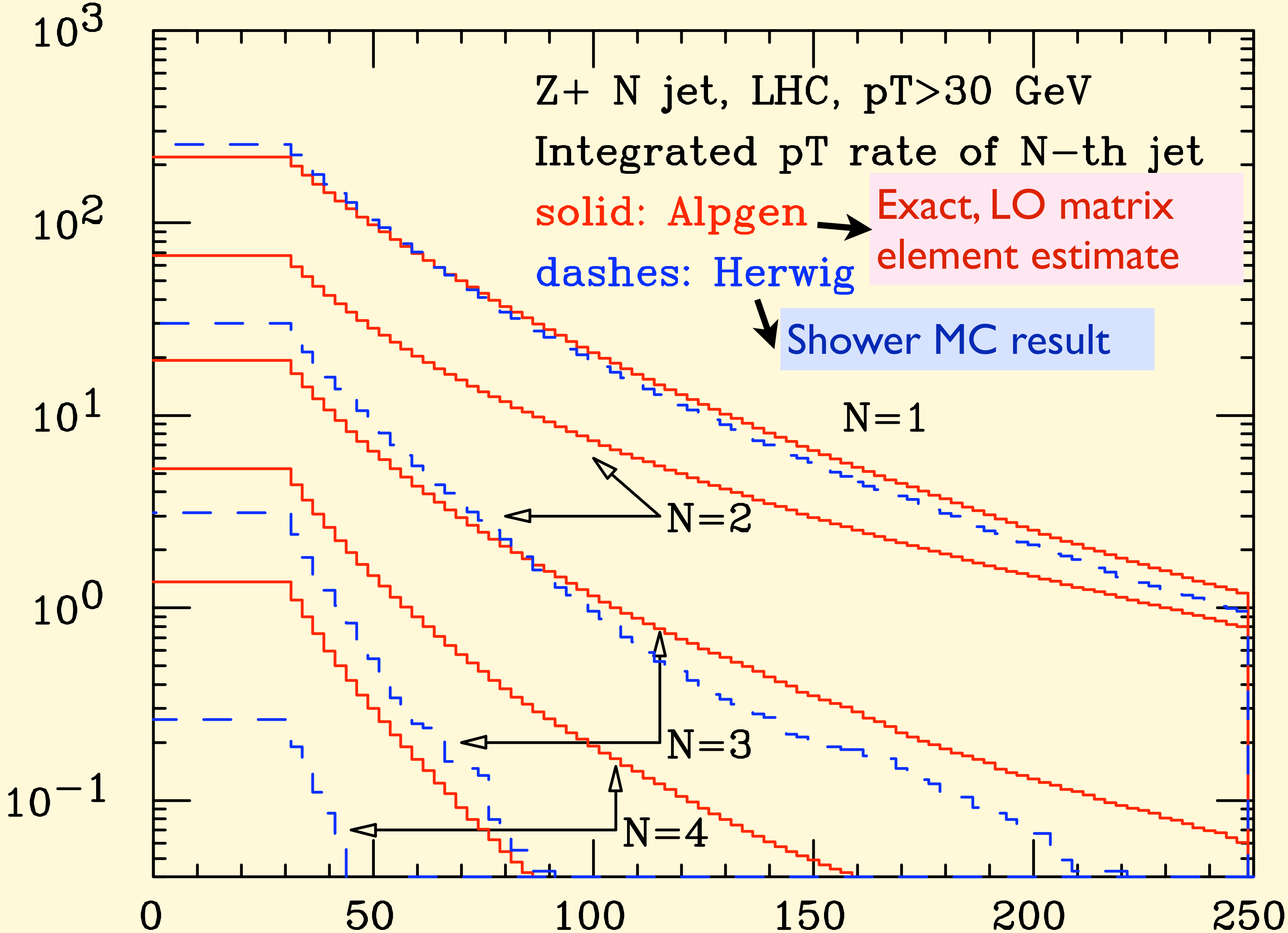
Accidental extra deposits of energy (cosmic rays on time, beam backgrounds, , electronic noise, etc.etc.etc.)



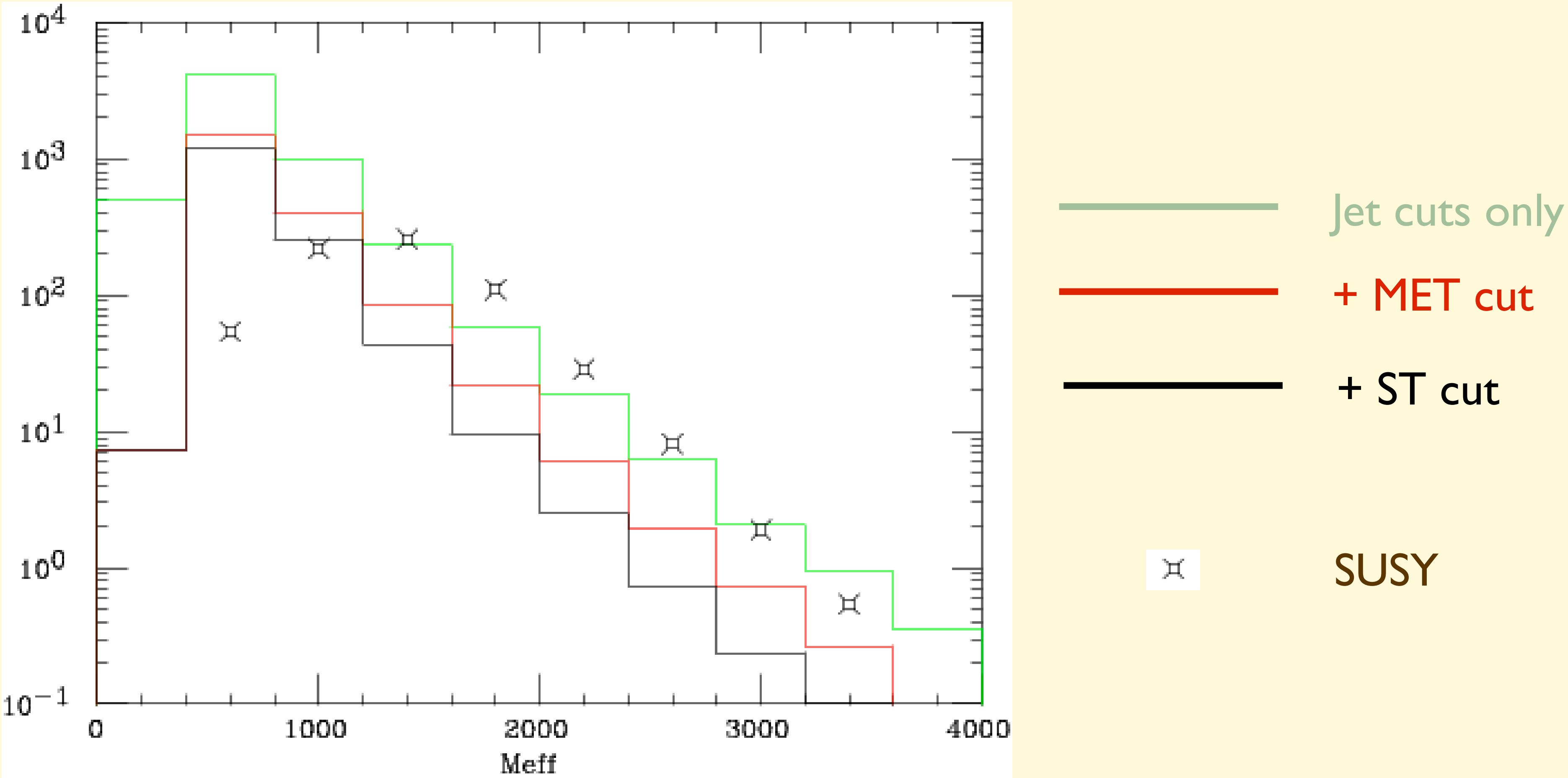
It is sufficient that these effects leave a permille fraction of the QCD rate for the signal to be washed away!

Z(→νν) + jets

1. Shower MC vs Matrix element results



Z(→νν) + jets



10^4
10^3
10^2
10^1
10^0
10^{-1}
0
1000
2000
3000
4000
Meff
Jet cuts only
+ MET cut
+ ST cut
SUSY

Normalizing the bg rate using data ...

Use Z->ee + multijets, apply same cuts as MET analysis but replace MET with ET(e^+e^-)

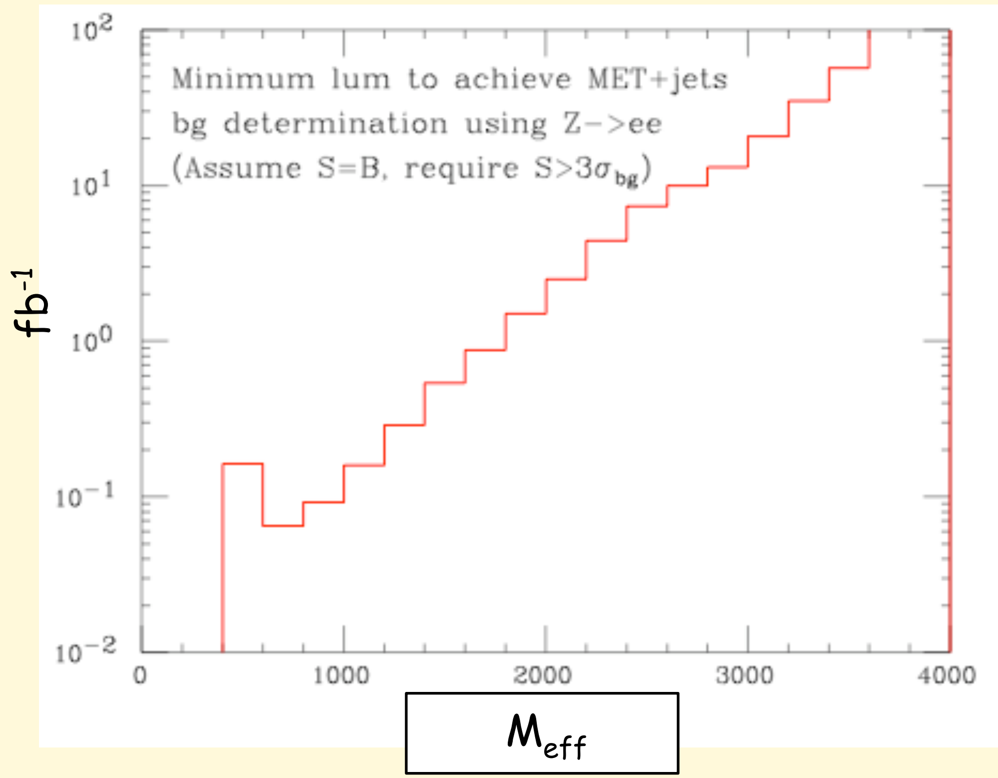
Extract $Z\to\nu\nu$ bg using, bin-by-bin:
$(Z\to\nu\nu) = (Z\to ee)\ B(Z\to\nu\nu)/B(Z\to ee)$

Assume that the SUSY signal is of the same size as the bg, and evaluate the luminosity required to determine the Z->nunu bg with an accuracy such that:

$N_{susy} > 3$ sigma

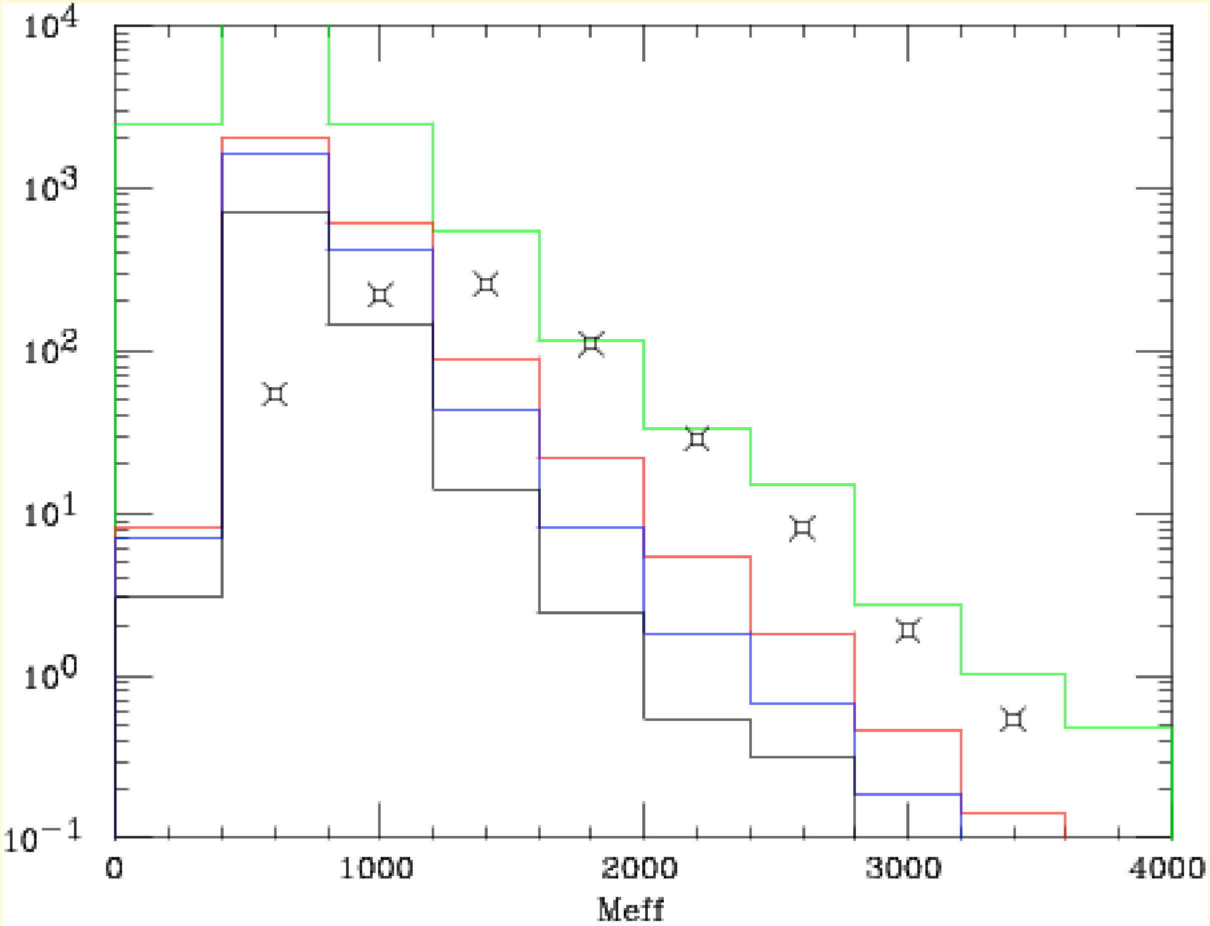
where

sigma=sqrt[N($Z\to ee$)] * $B(Z\to\nu\nu)/B(Z\to ee)$



=> several hundred pb^{-1} are required. They are sufficient if we believe in the MC shape (and only need to fix the overall normalization). Much ore is needed if we want to keep the search completely MC independent

W(→lν) +4 jets



Jet cuts only

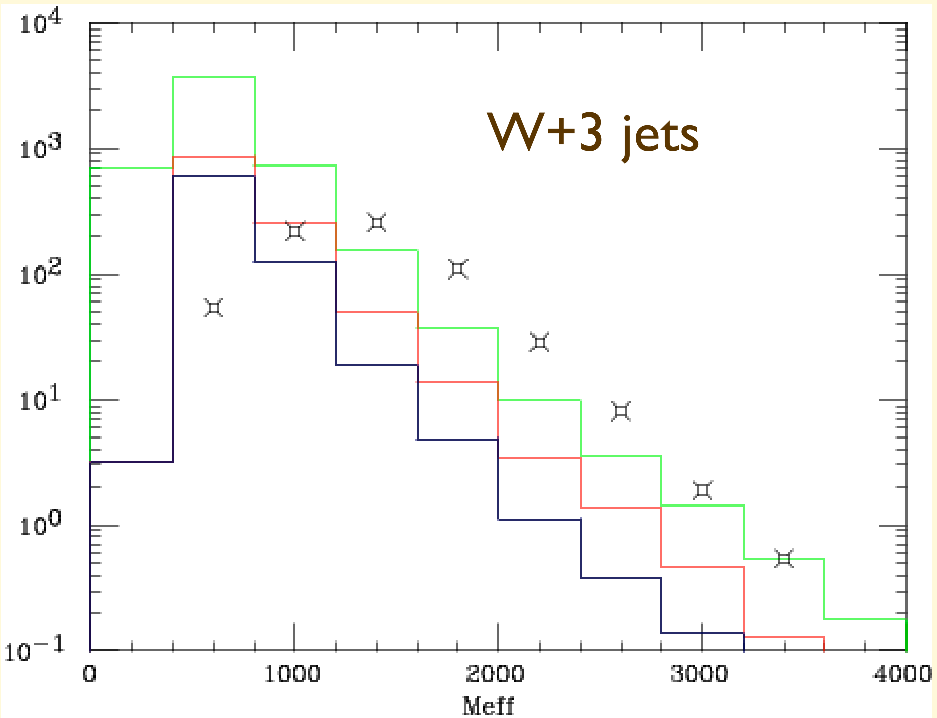
+ MET cut

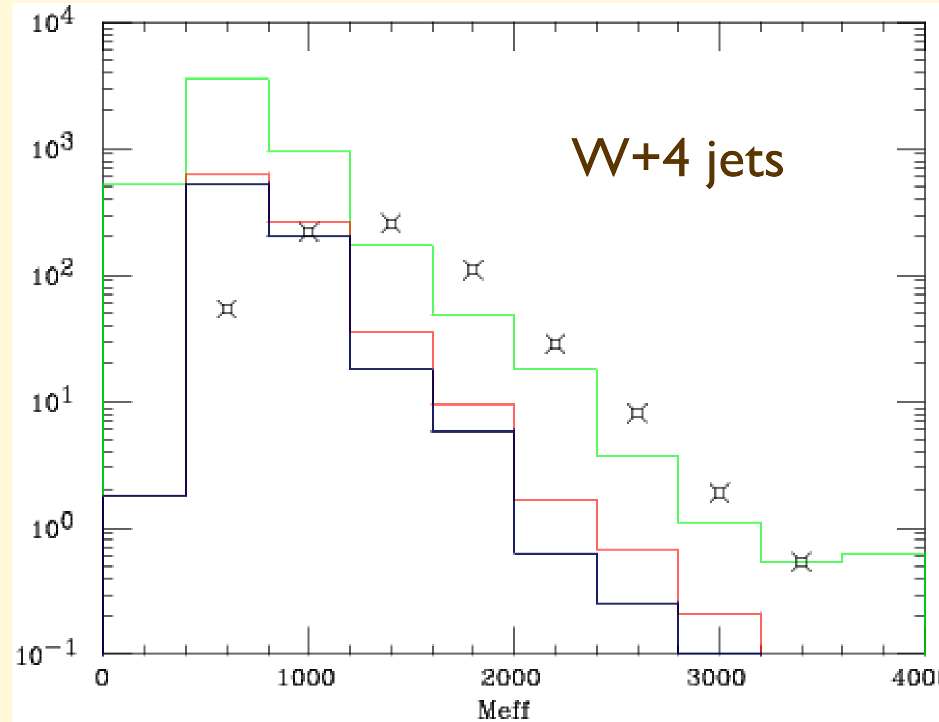
+ ST cut

+ pt_{lept}<20

SUSY

W(→tau-jet ν) + jets





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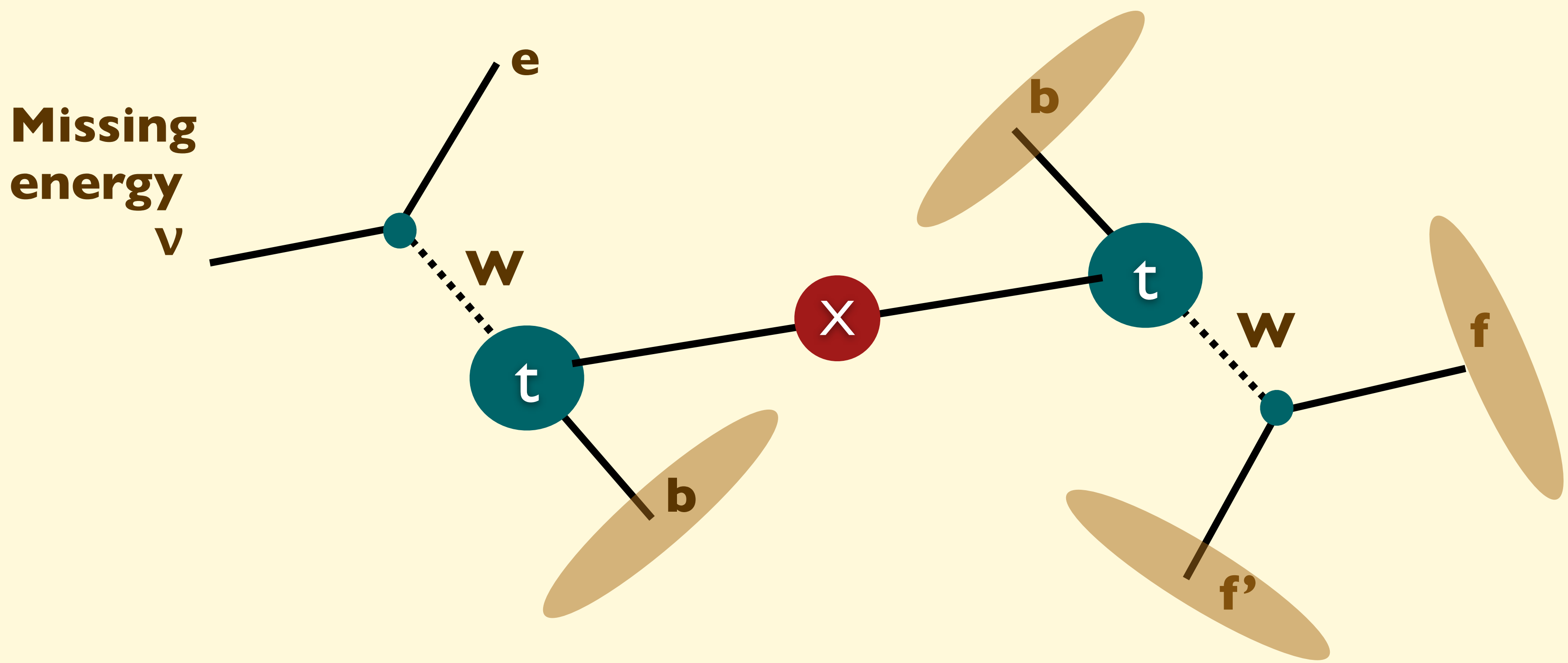
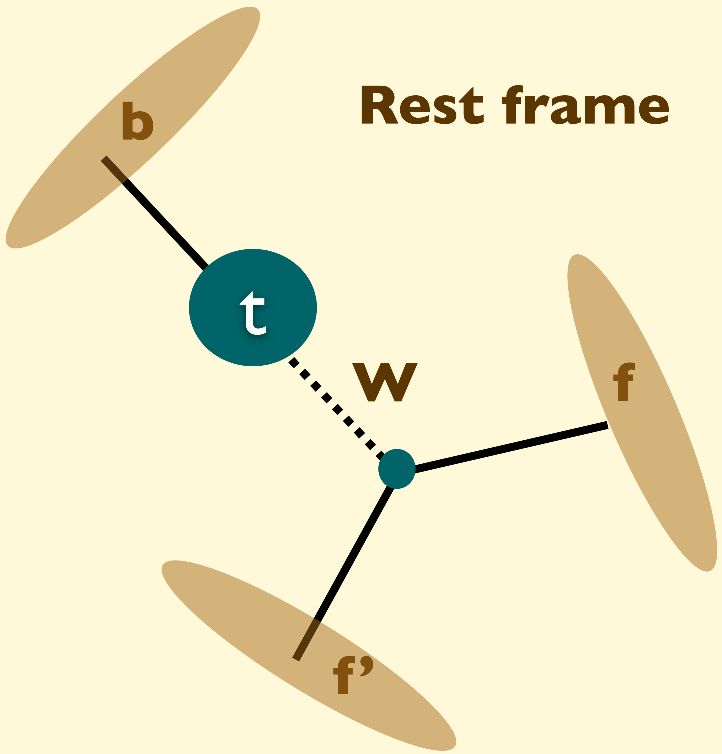
SUSY

Top final states



Top final states

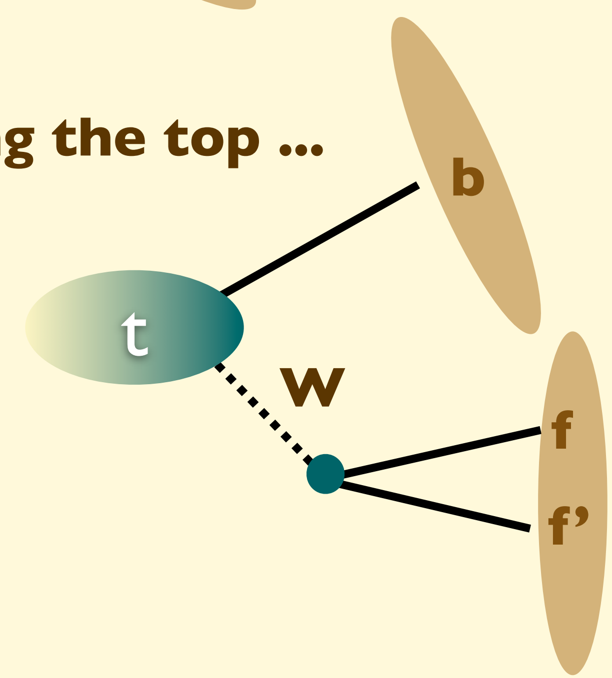
Rest frame

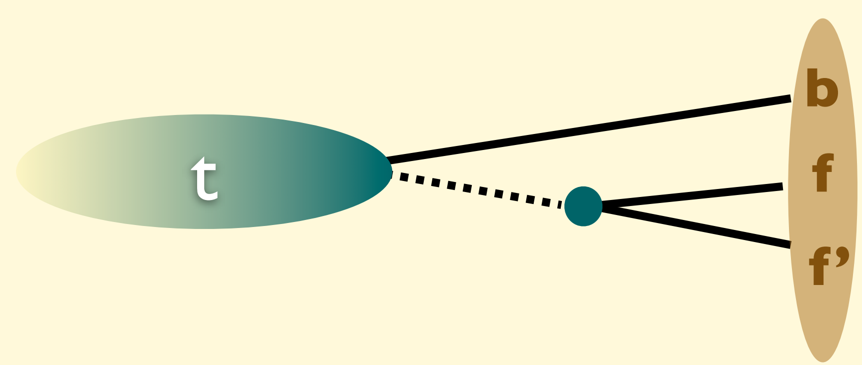


$$p_b = \frac{m_{top}^2 - m_W^2}{2\, m_{top}}$$

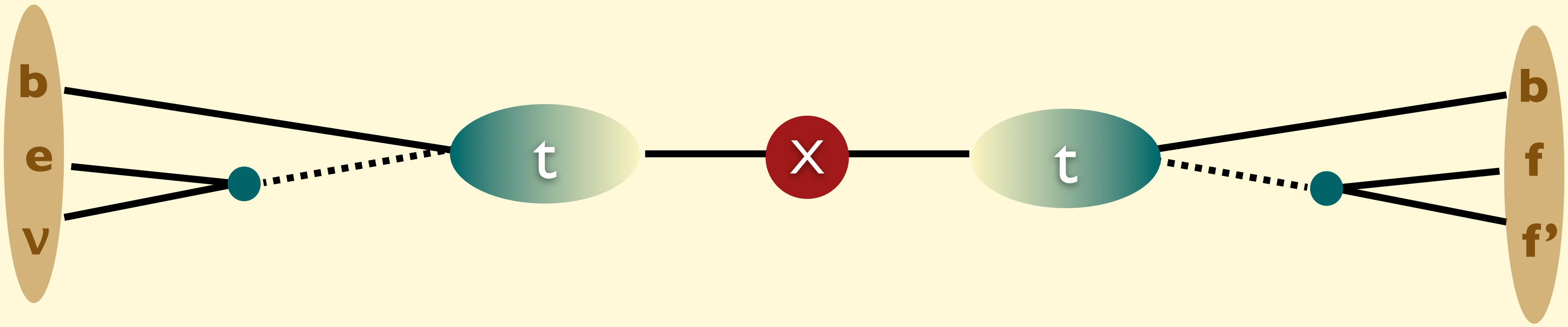
$$p_f^{max} = \frac{m_W}{2}\,\frac{m_{top}^2 + m_W^2}{2\, m_{top}\, m_W}$$

Boosting the top ...



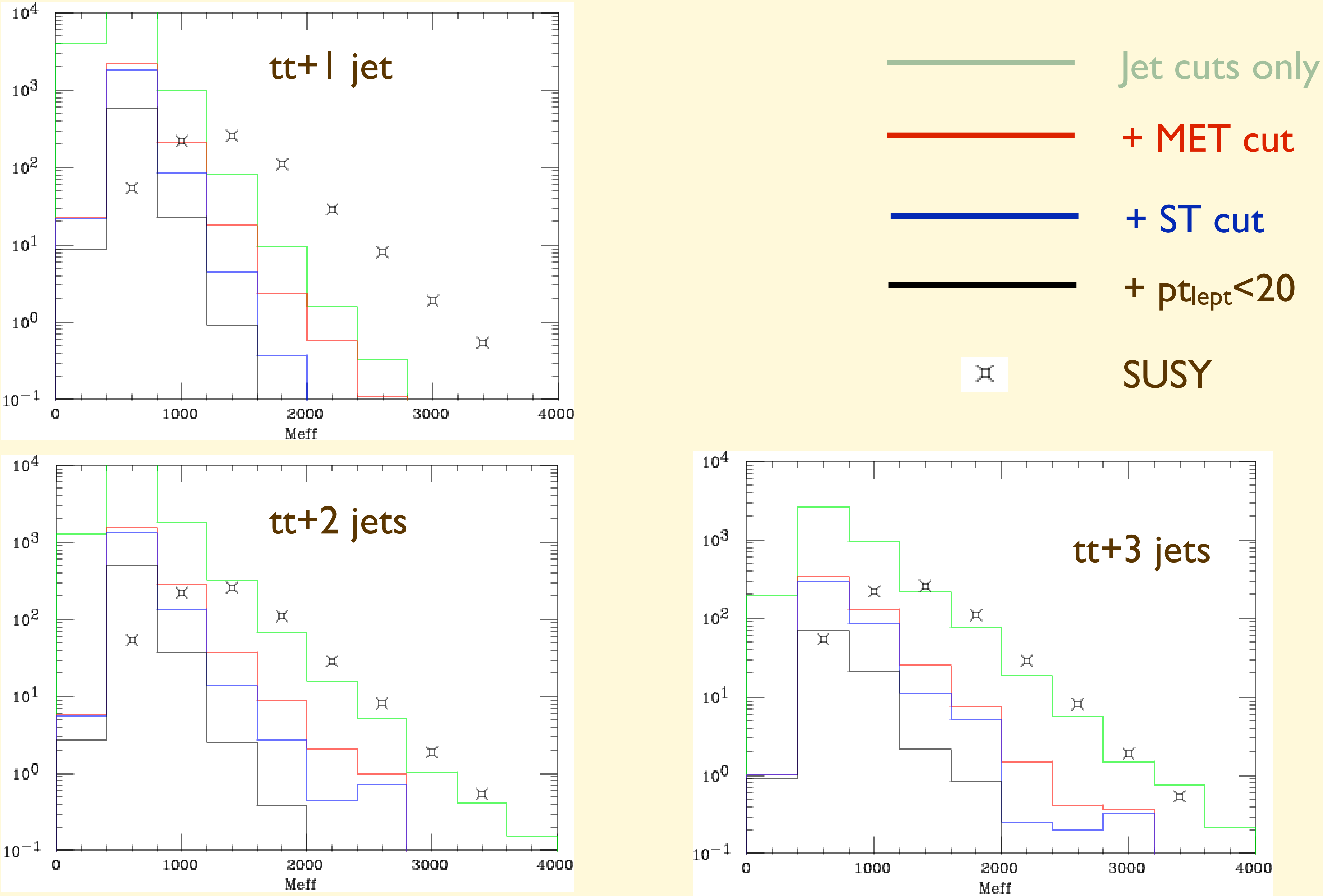


Top final states



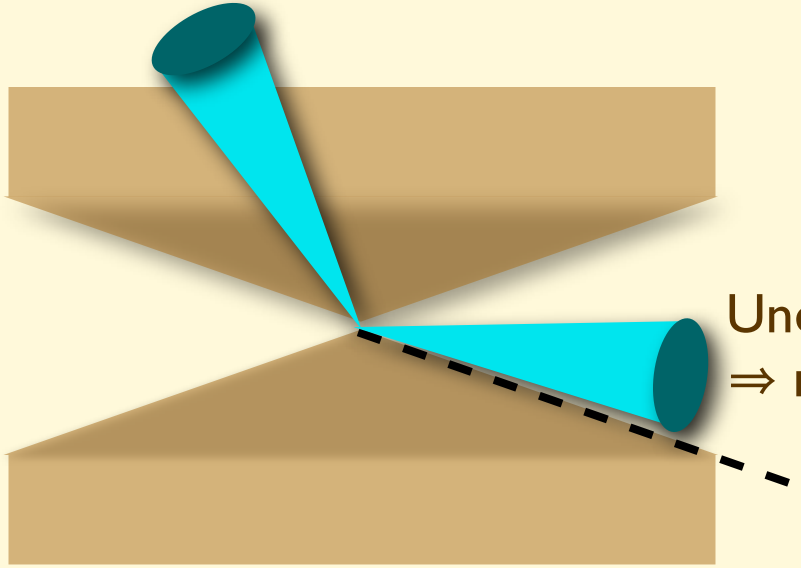
Large Meff leads to highly collimated final states

Sphericity and multi-jet cuts very effective against the leading-order t-tbar contribution!

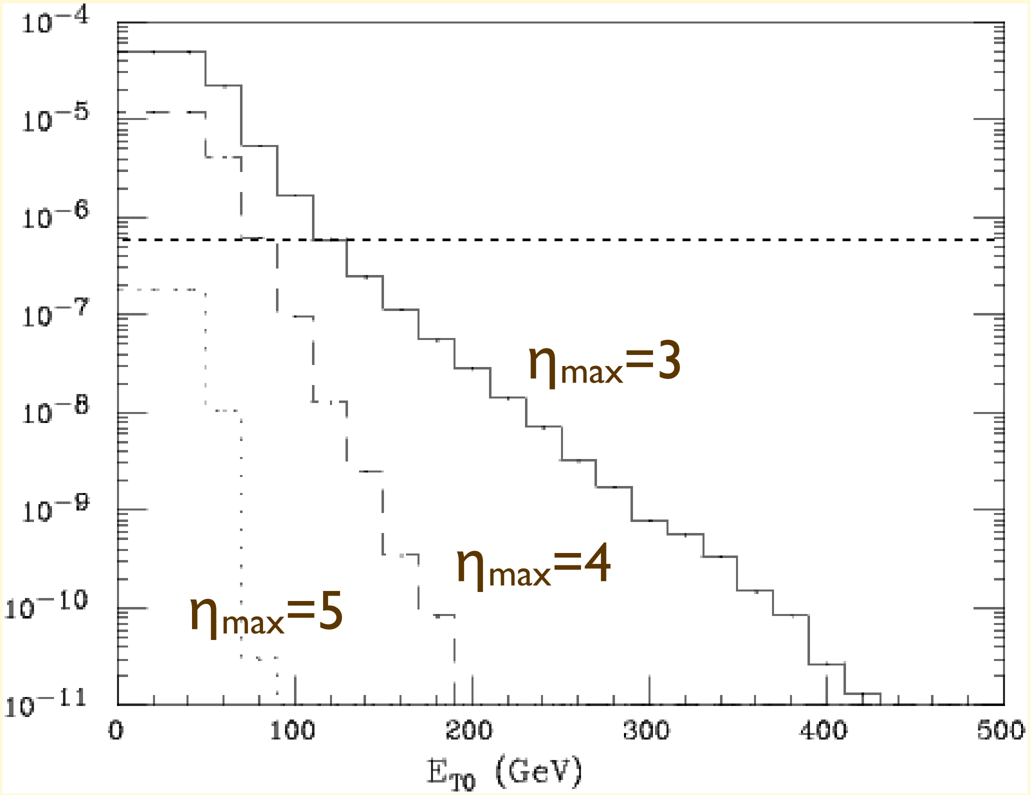


All jet multiplicities contribute at approximately the same level!!

Instrumental sources of missET: incomplete calorimeter η coverage



σ(jet-jet with MET> E_{To}) / σ(pp→X)

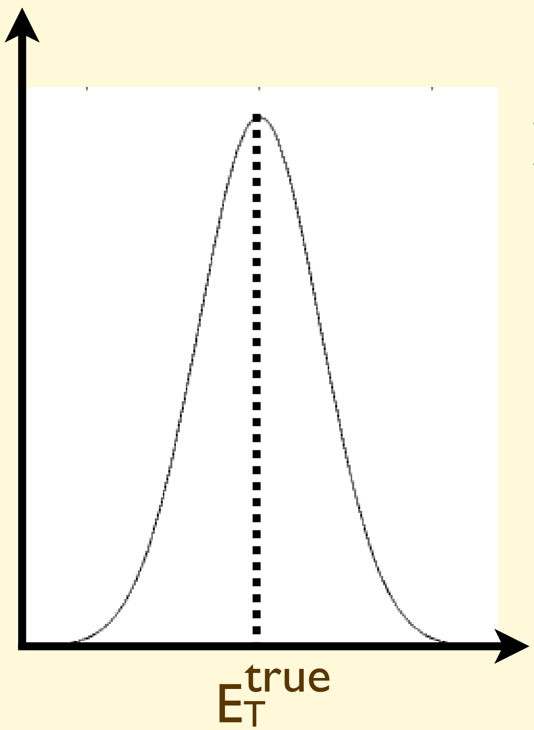


cfr:
$\sigma(W\to l\nu)/\sigma(pp\to X) \approx 6 \times 10^{-7}$

NB:
At $L=10^{34}\ cm^{-2}\ s^{-1}$,
⟨N(pp collisions)⟩ ≈ 20

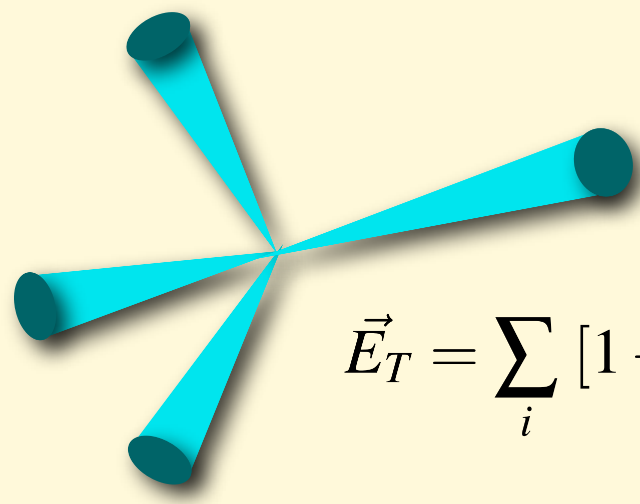
⇒ **probability 20x larger**

Instrumental sources of missET: jet energy resolution



$$\text{Prob}[p_T] \propto \exp -\frac{(p_T - p_T^{\text{true}})^2}{\sigma^2}$$

$$\sigma = C\sqrt{E_T^{\text{true}}/\text{GeV}}, \quad C = O(1)$$

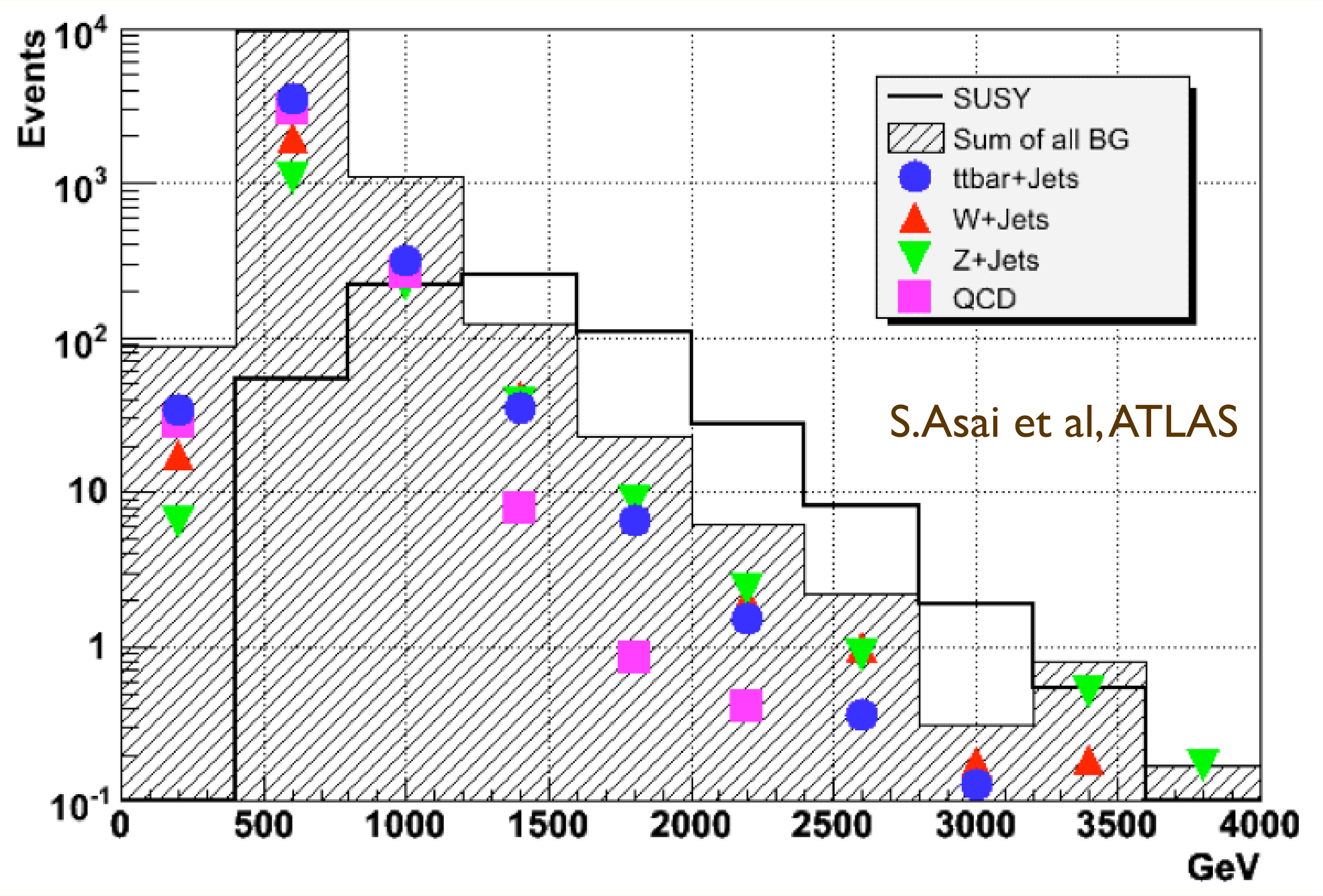


$$\vec{E}_T = \sum_i \left[1 + \delta_i\right] \vec{p}_{T,i}^{\,\text{true}} = \sum_i \delta_i \, \vec{p}_{T,i}^{\,\text{true}}$$

$$\langle |\vec{E}_T|^2 \rangle = \sum_{i,j} \langle \delta_i \delta_j \rangle \vec{p}_{T,i} \cdot \vec{p}_{T,j} \qquad \langle \delta_i \delta_j \rangle = \frac{C^2}{p_{T,i}} \delta_{ij}$$

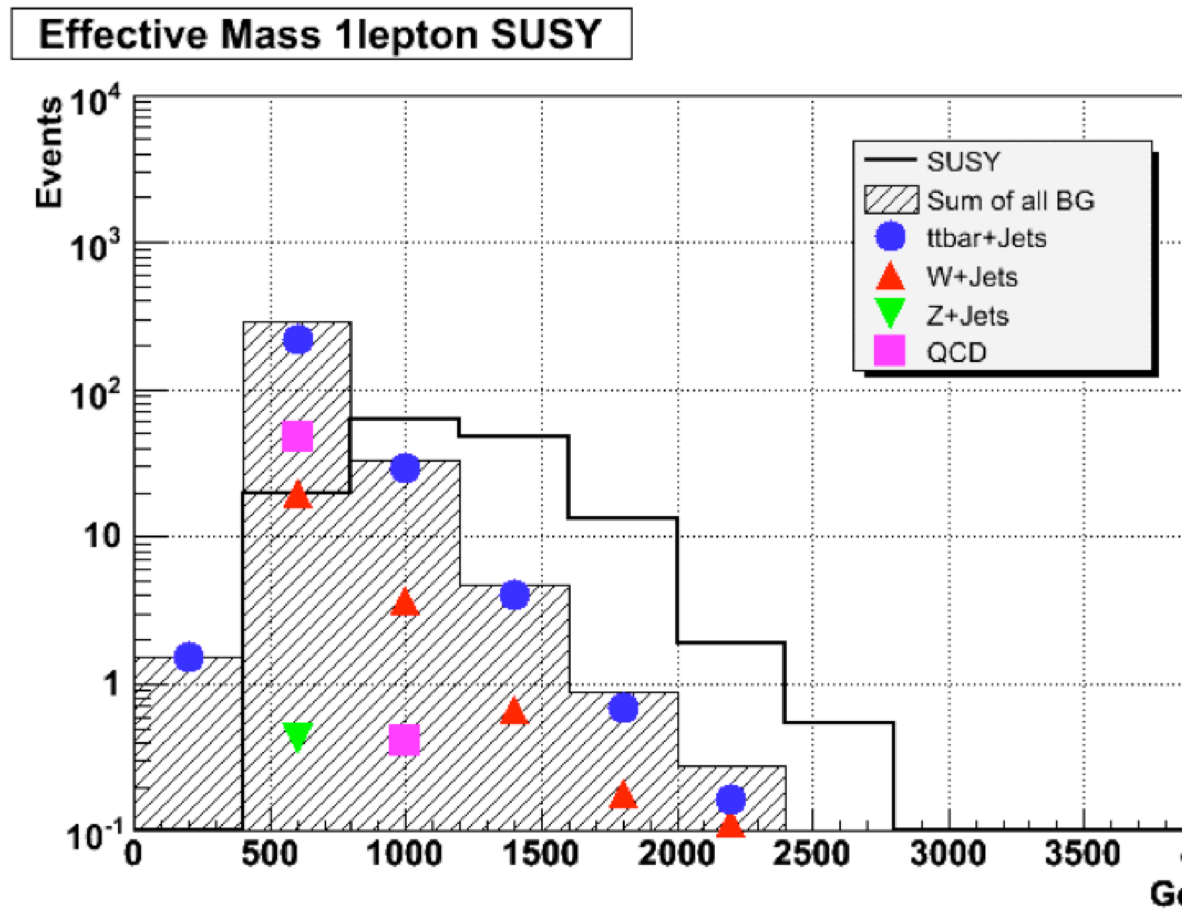
$$\langle \text{MET} \rangle = C \sqrt{\sum_i p_{T,i}}$$

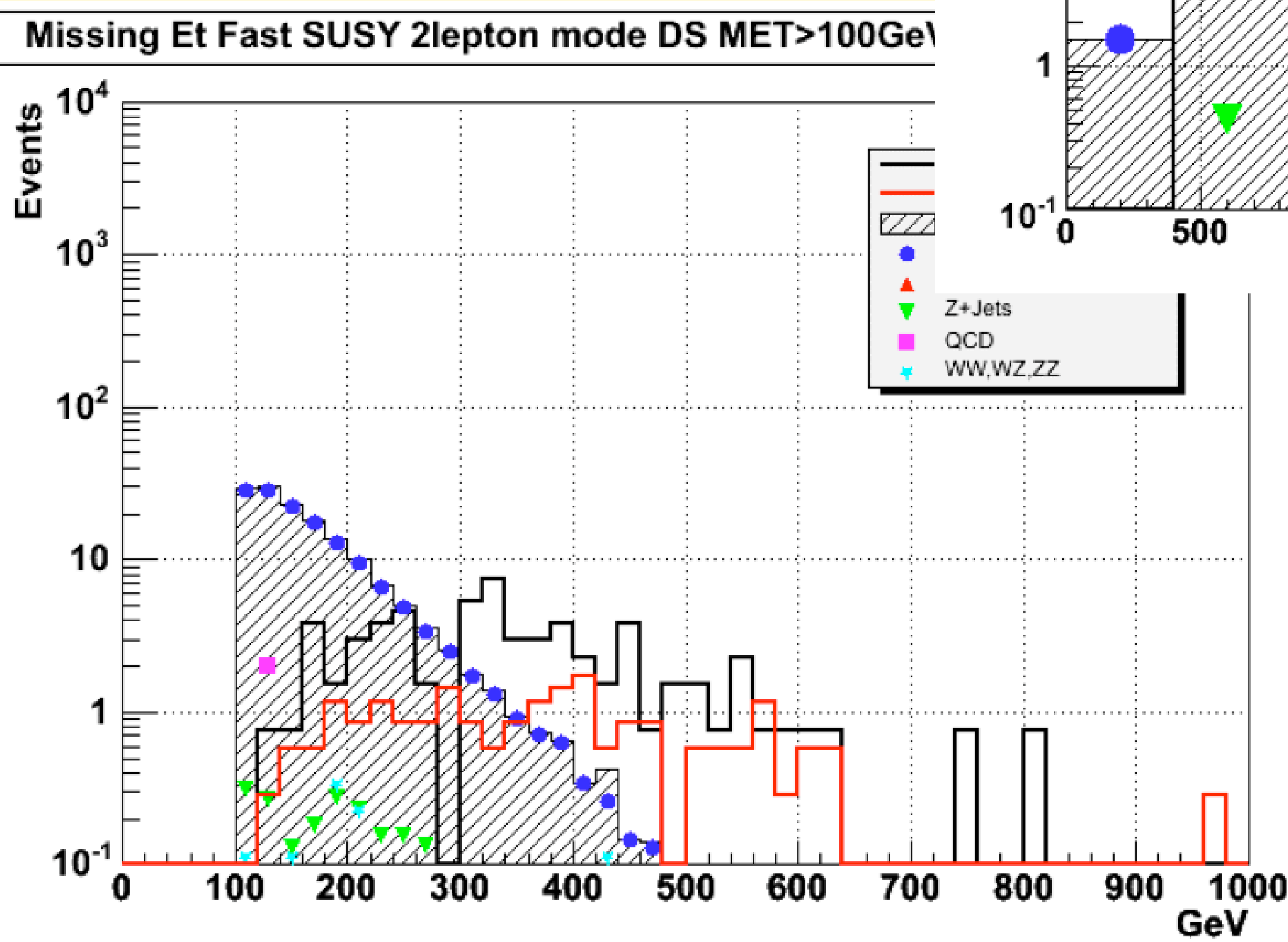
Overall result, after the complete detector simulation, etc....



Adding leptons ...

S.Asai et al, ATLAS

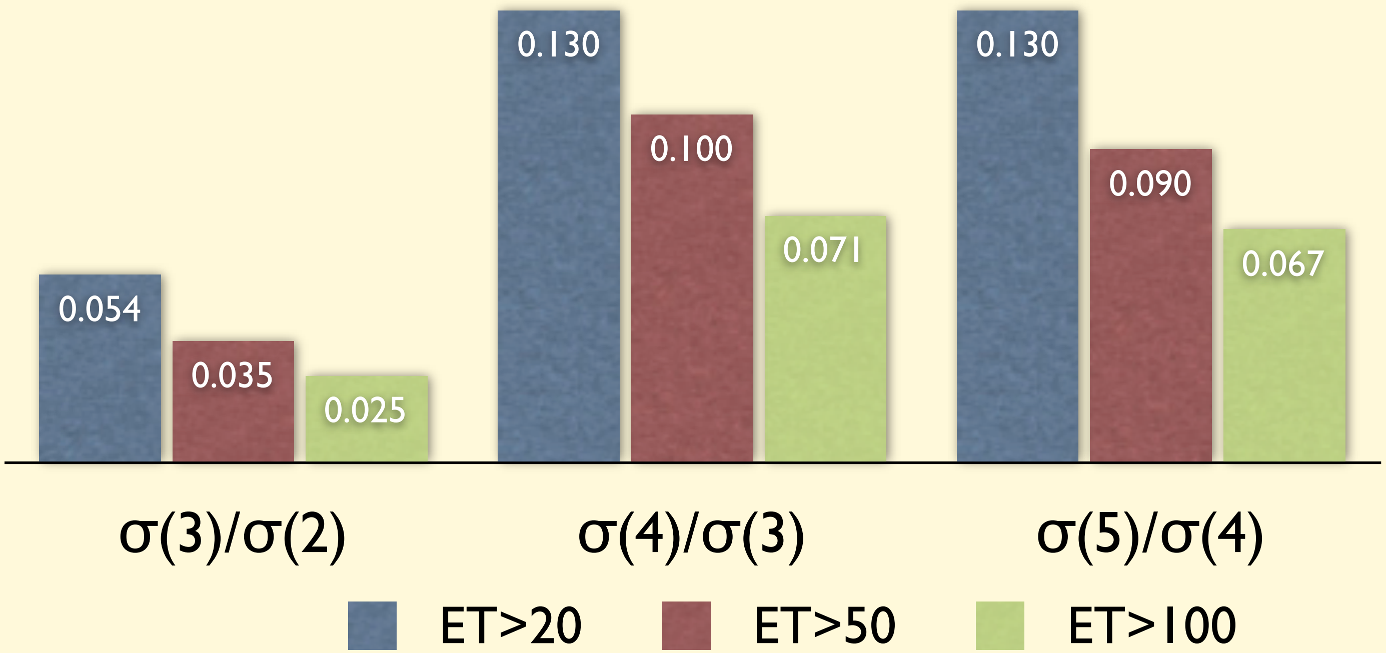




Some properties of rates for multijet final states

Multijet rates

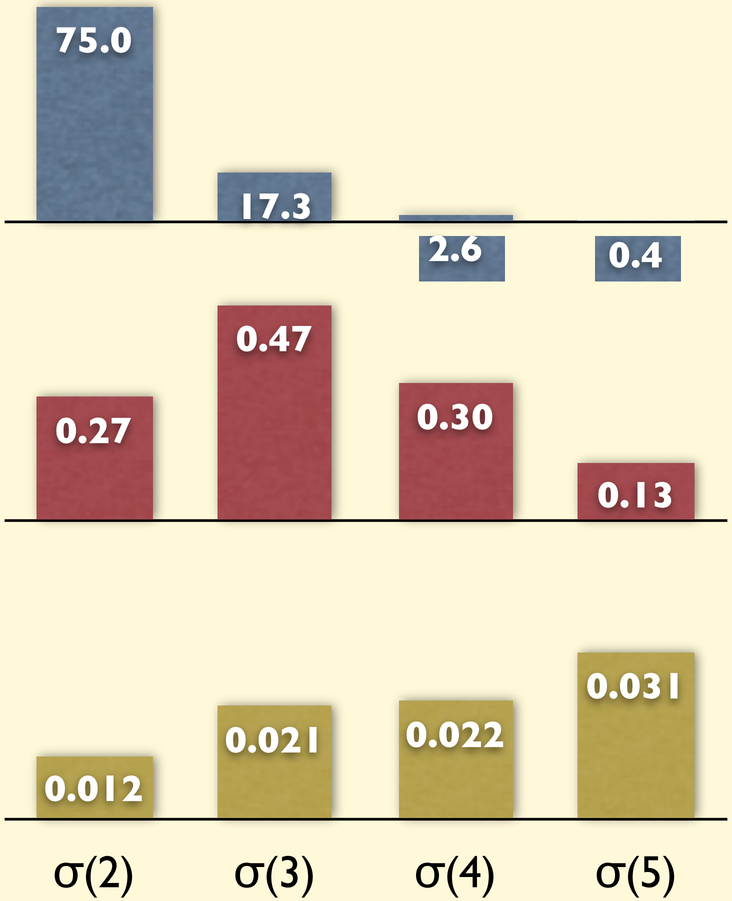
σ [μb]	N jet=2	N jet=3	N jet=4	N jet=5
E_T^{jet} >20 GeV	350	19	2.6	0.35
E_T^{jet} >50 GeV	12.7	0.45	0.045	0.004
E_T^{jet} >100 GeV	0.85	0.021	0.0015	0.0001



- The higher the jet E_T threshold, the harder to emit an extra jet
- When several jets are already present, however, emission of an additional one is less suppressed

Multijet rates, vs $\sqrt{s}$, with $E_T^{jet} > 20$ GeV

σ [μb]	N jet=2	N jet=3	N jet=4	N jet=5
√s > 100 GeV	75	17.3	2.6	0.37
√s > 500 GeV	0.27	0.47	0.30	0.13
√s > 1000 GeV	0.012	0.021	0.022	0.031

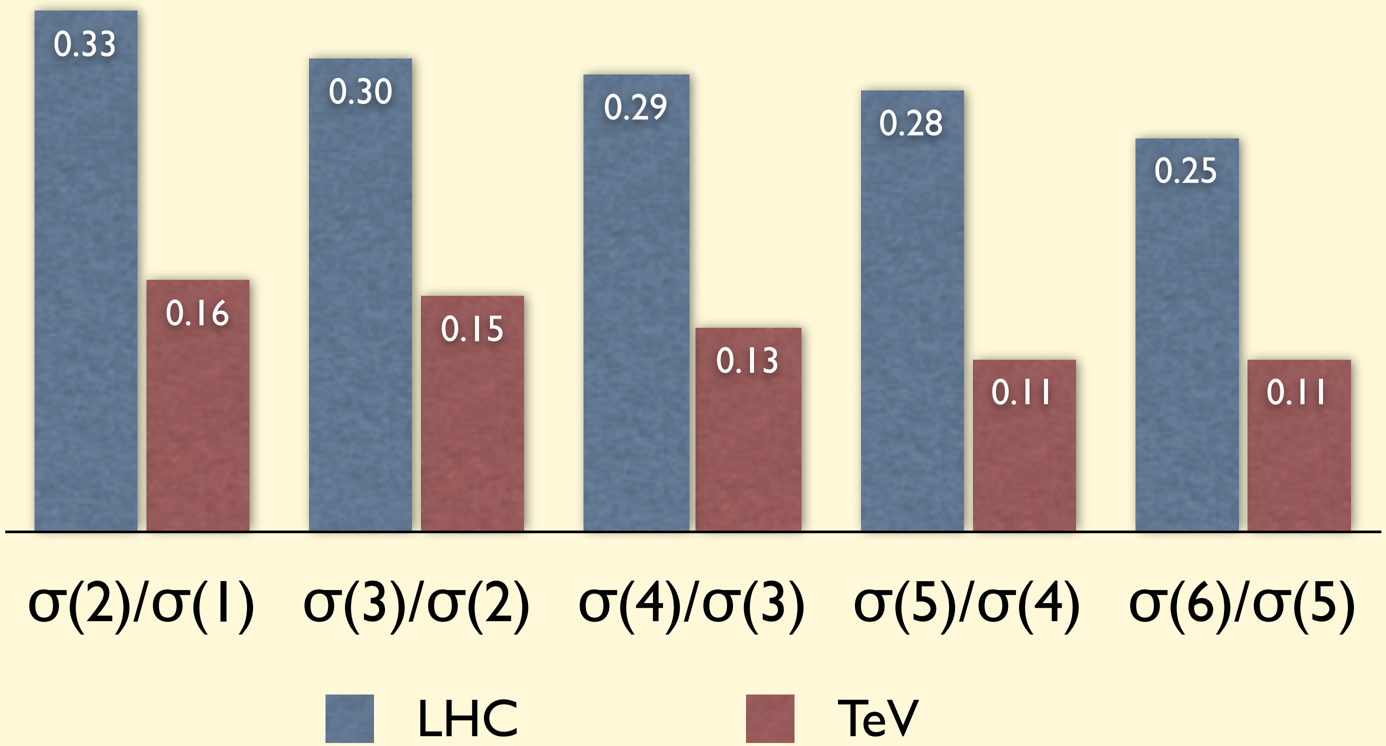


High mass final states are dominated by multijet configurations

W+Multijet rates

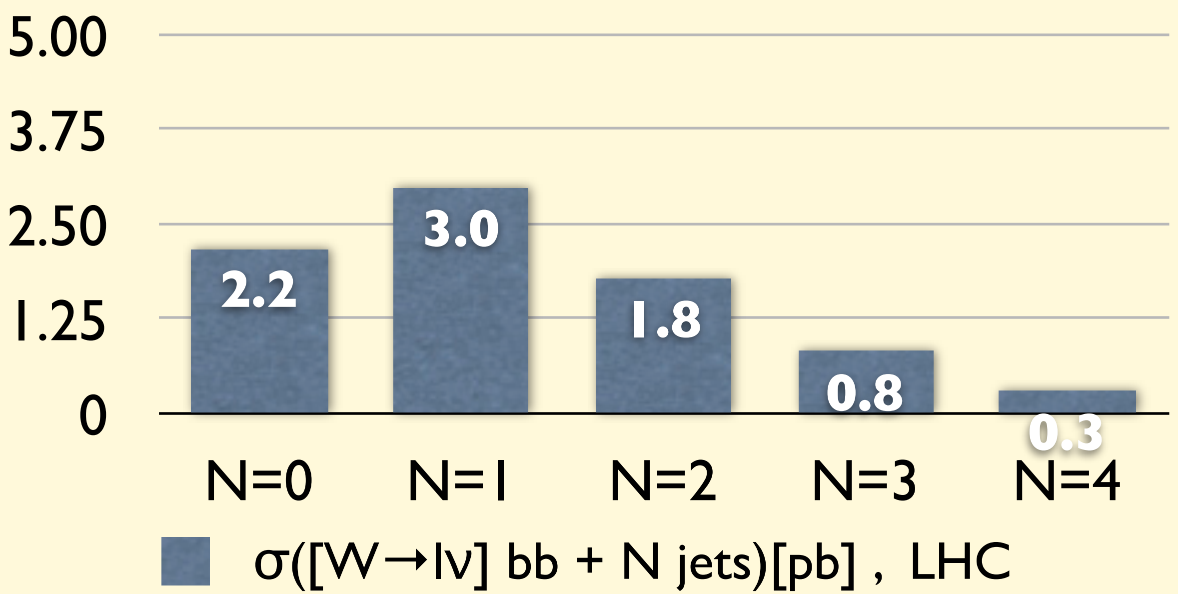
$\sigma \times B(W \to e\nu)$[pb]	N jet=1	N jet=2	N jet=3	N jet=4	N jet=5	N jet=6
LHC	3400	1130	340	100	28	7
Tevatron	230	37	5.7	0.75	0.08	0.009

E_T(jets) > 20 GeV , $|\eta|<2.5$, $\Delta R>0.7$



- Ratios almost constant over a large range of multiplicities
- $O(\alpha_s)$ at Tevatron, but much bigger at LHC

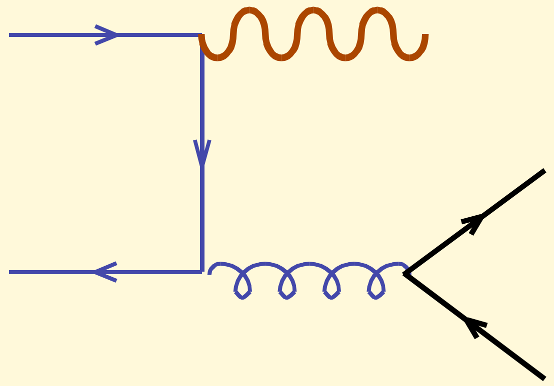
Wbb+jets rates

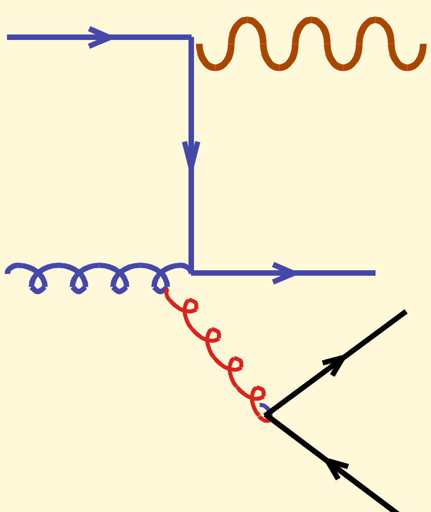


Pattern of multiplicity distribution very different than in W+jets!

In pp collisions (contrary to the Tevatron, p-pbar) :

$N_{jet}=0 \propto \alpha_s^2$ x Lum(q qbar) $\approx$ $N_{jet}=1 \propto \alpha_s^3$ x Lum(q g)





Beware of naive α_s power counting!!

Leptons

Experimentally, electrons, muons and taus are entirely different objects. Their identification requires different components of the detector, different techniques, and is subject to different backgrounds.

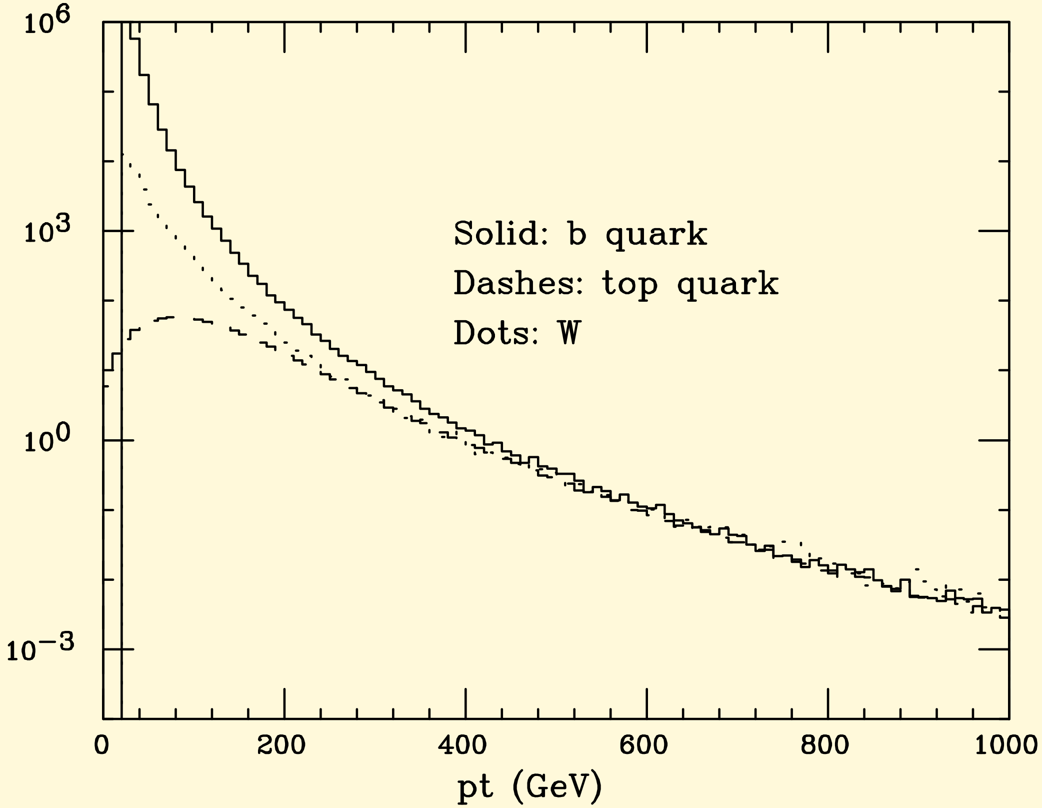
As seen from a theorists, all leptons are produced the same. Nevertheless there is a large variety of possible production mechanisms, each one of them leading to different overall properties of the final state. When considering leptons as a signal for new physics, it is important to have a clear picture of their irreducible SM sources

Single lepton

Sources of single high-pt leptons:

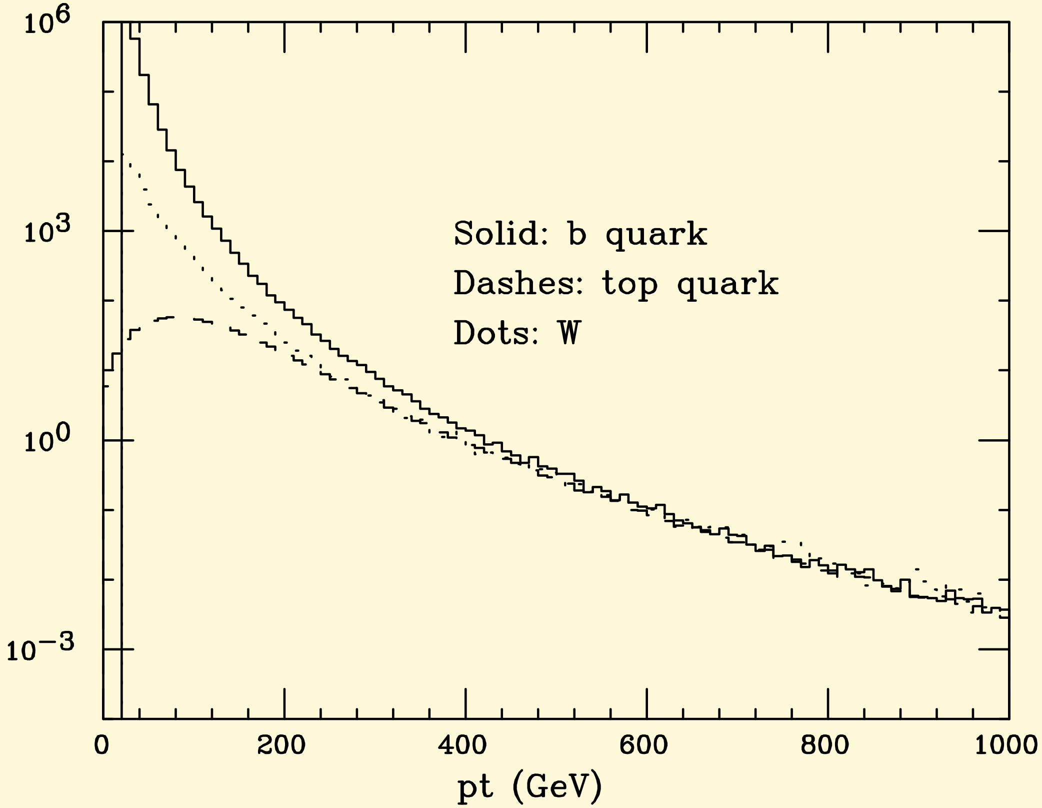
- $W \to e/\mu + \nu$
- $Z \to \tau\tau \to e/\mu + X$
- $b \to e/\mu + X$
- $t \to Wb \to e/\mu + \nu + b$

Differential Rates



- At large pt b and t production ~ equal !
- At large pt, W and heavy quark production ~ equal!

Differential Rates



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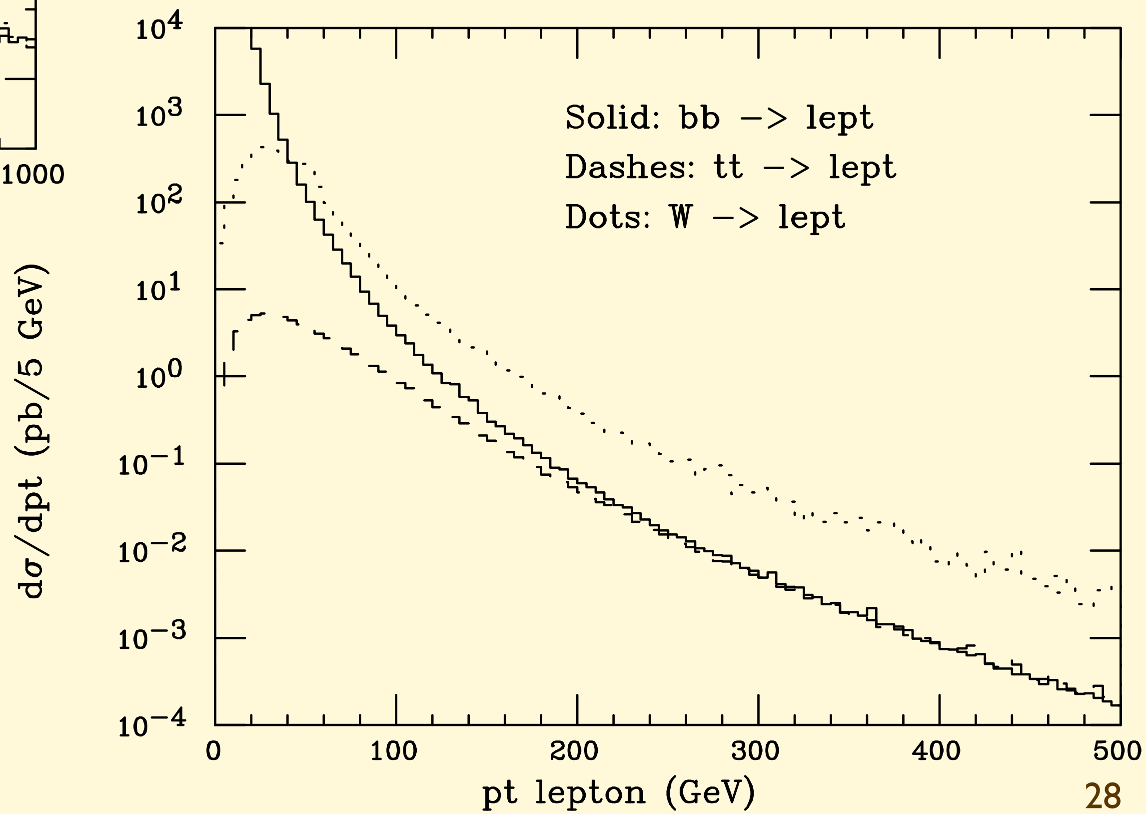

*W → lepton is a 2-body decay, b/t → lepton is 3-body: lepton takes a larger fraction of momentum in W decay => harder spectrum, larger rate at higher pt in W production

*The global features of the event accompanying the lepton will clearly be very different in each case. Which of the three processes will dominate in a given analisys, will therefore depend on the details

How come Q and W spectra are comparable at large Et?

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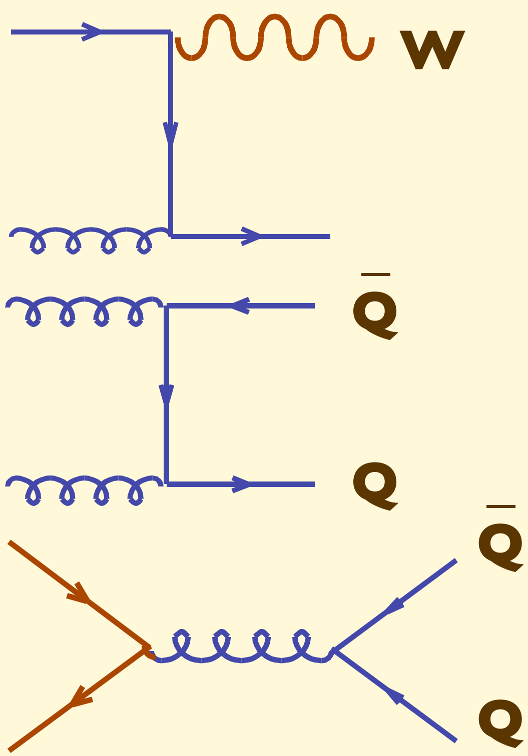
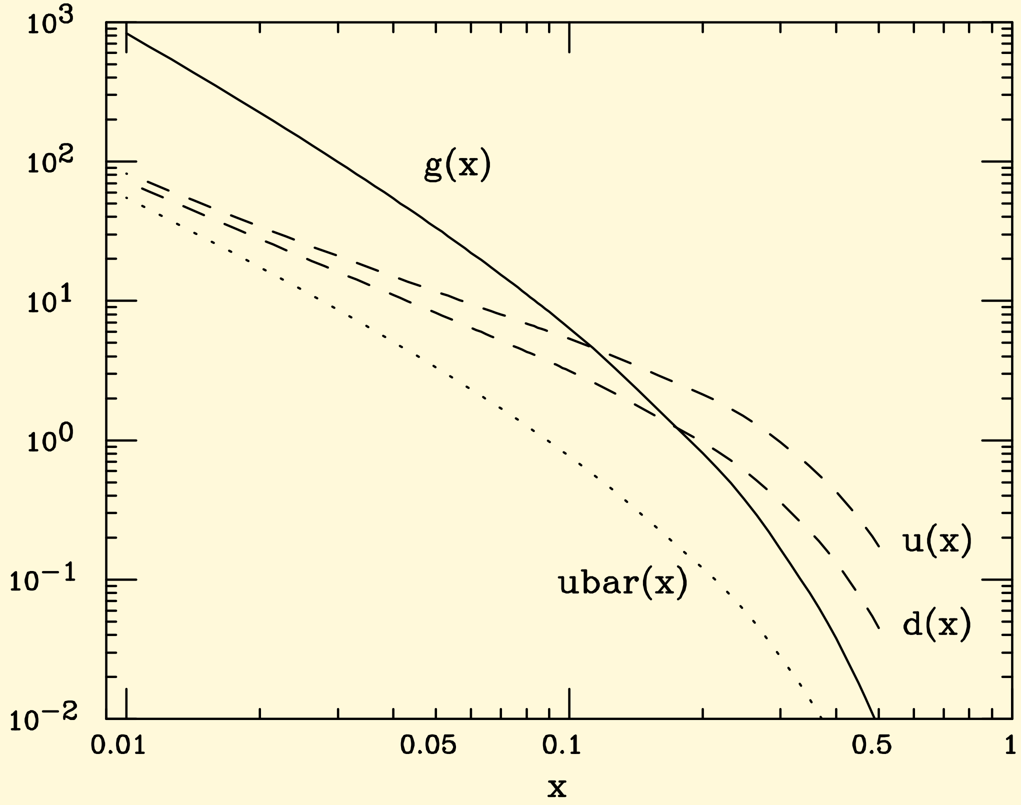
The LO processes for QQ production are weighted by the gg or qqbar luminsity, which drops at large mass much more rapidly than L(qg)





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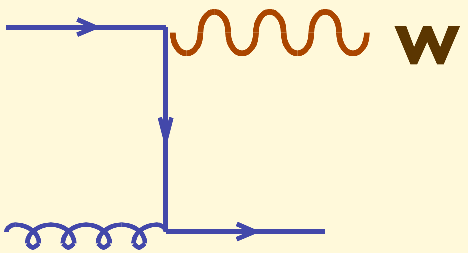

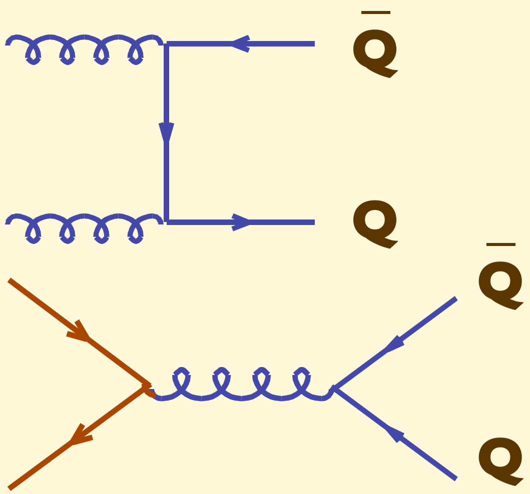

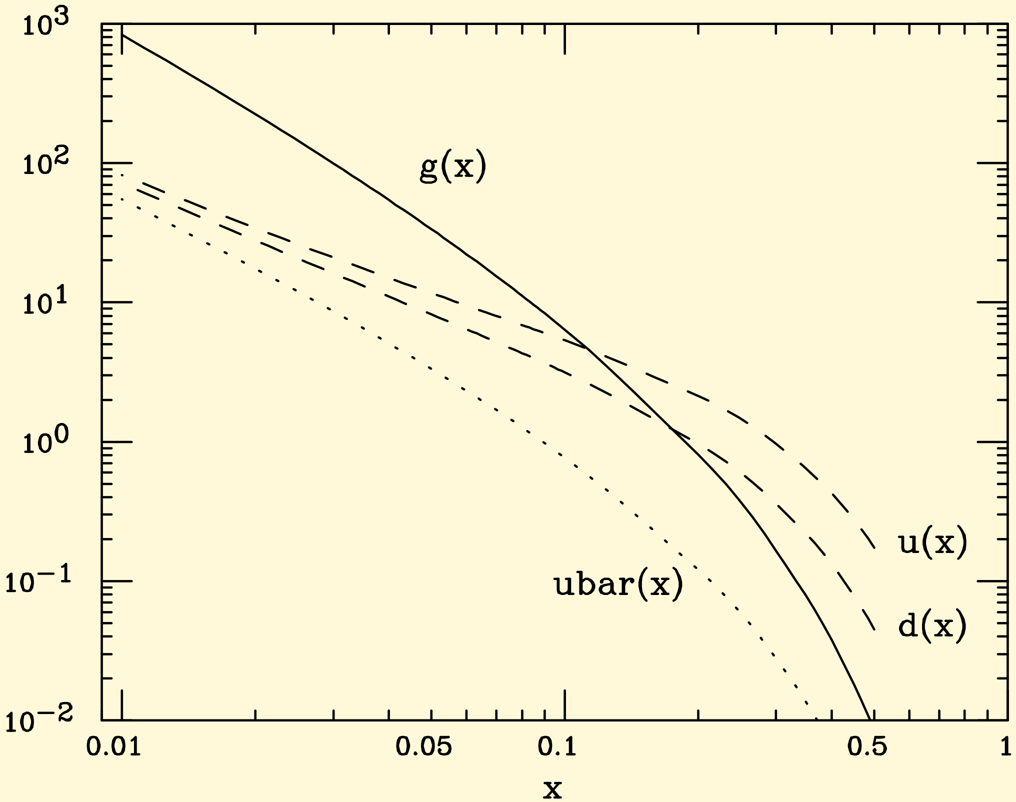

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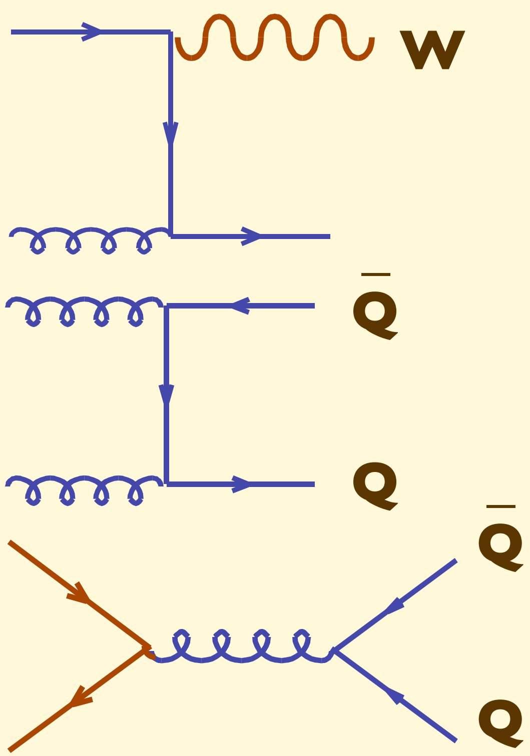
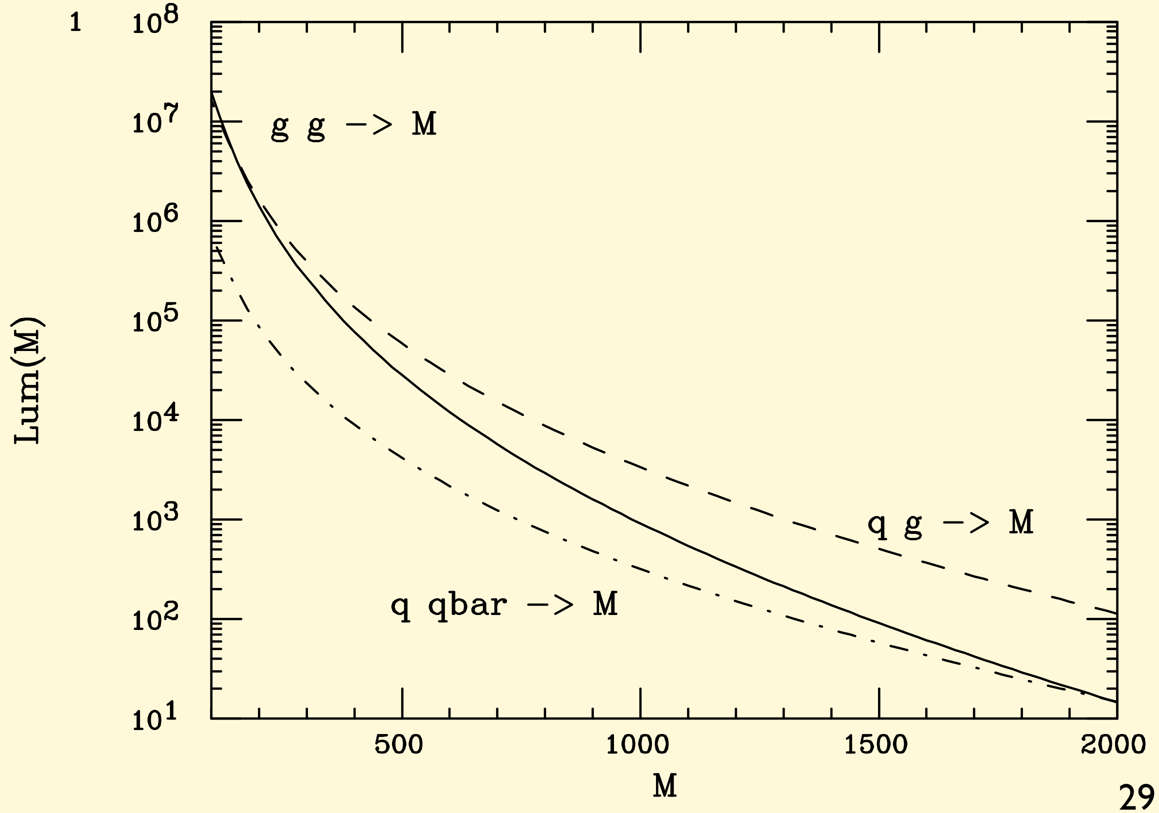
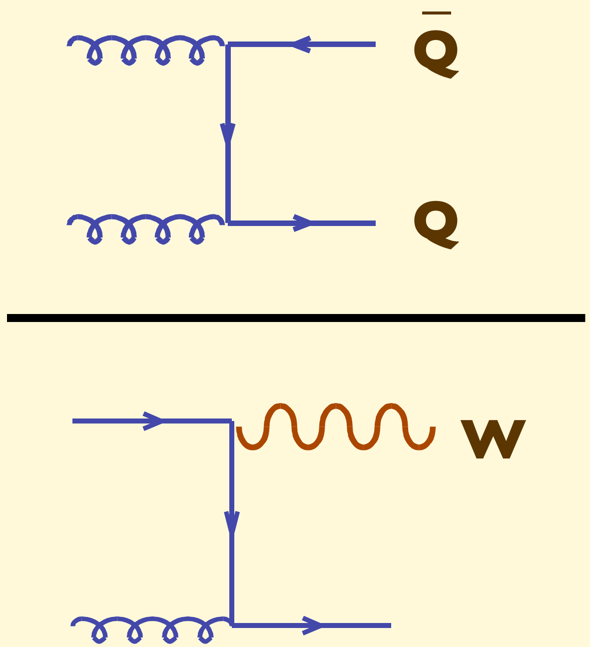






Quark colour charge

Initial state colour averages

~1/3 at 90°

$$= \frac{C_F \alpha_s}{1/2 \times \alpha_w} \times \left(\frac{N}{N^2 - 1}\right) \times \frac{1}{1/2} \times F(s \leftrightarrow u)$$

Quark weak charge

V-A, only L-handed quarks

$$\approx \quad \frac{\alpha_s}{\alpha_W} \sim 3$$

Dileptons

One lepton W: 160 nb

WW	tt	Z
75pb	500pb	50nb
2l+MET, no jets	2l+MET, jets, b's	2l, m(ll)=mZ, no MET, no jets

Dilepton production dominated by top pairs!

Trileptons

WWW	ttW	ZW
130fb	500fb	28pb

ttW ~ 10^{-3} tt => trilepton contribution from tt, with 3rd lepton form b→l decay, important => require isolation!

Quadrileptons

WWWW	tttt	ZWW
0.6fb	12fb	100fb

ZWWW=0.7fb

Ratios

W/Z	WW / WZ	WWW / WWZ	WWWW / WWWZ
3	2.5	1.3	1

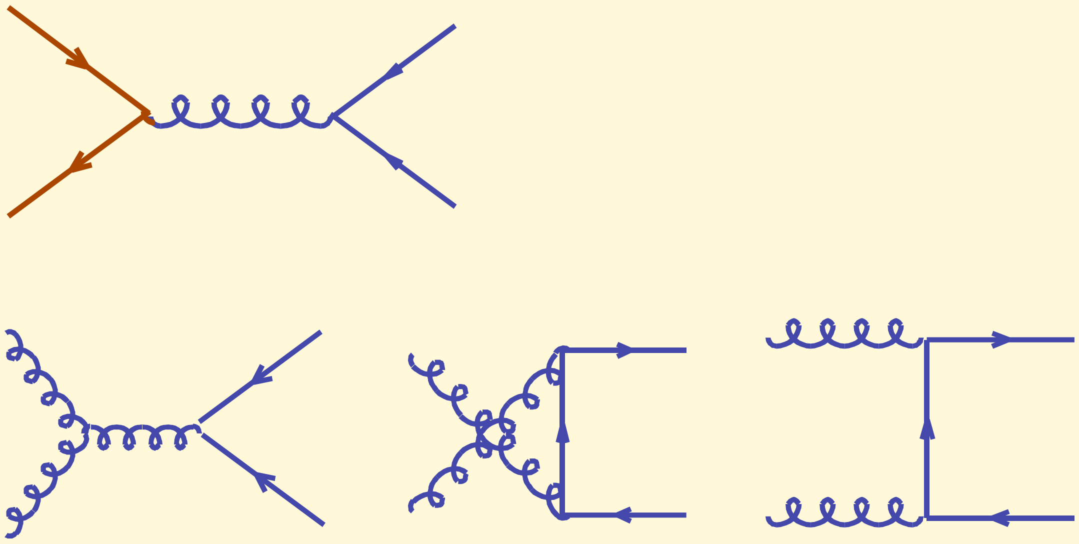
Ratio determined by couplings to quarks, u/d asymmetry of proton → Ratio determined by couplings among W/Z, SU(2) invariance

WW/W	WWW / WW	WWWW / WWW
5.0E-04	2E-03	5E-03

ZW / W	ZWW / WW	ZWWW / WWW
5.0E-04	4E-03	7E-03

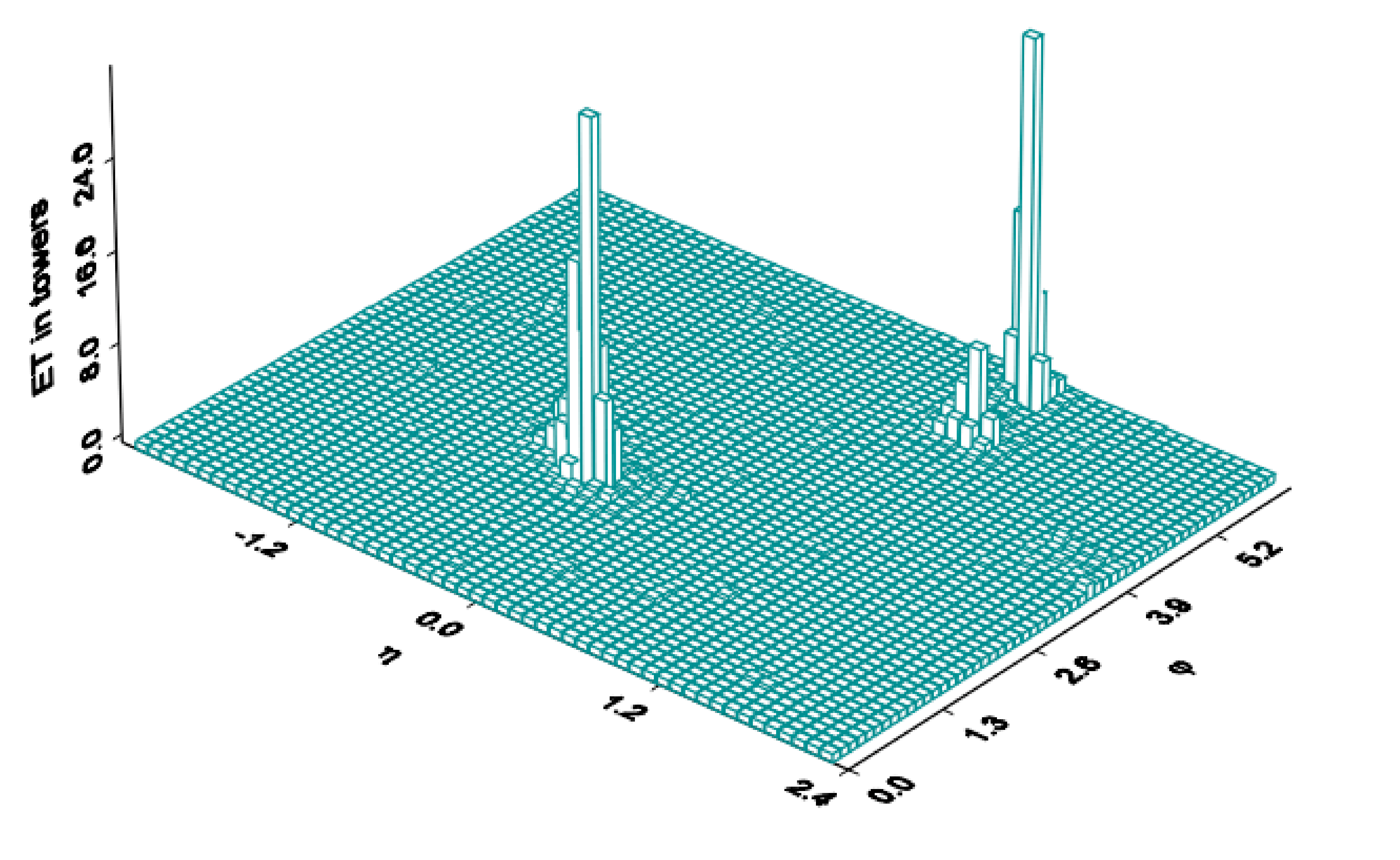
1W $\sim 10^{-3}$

Top production and bgs



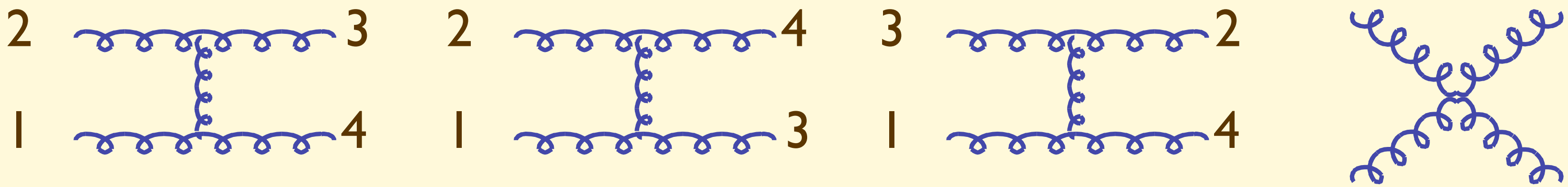
	σ(tt) [pb]	σ(W+X)	σ(W+bbX) [ptb>20 GeV]	σ(W+bbjj X) [ptb,ptj >20 GeV]
Tevatron	6	20 x 10^3	3	0.16
LHC	800	160 x 10^3	20	16
Increase	x 100	x 10	x 10	x 100

Jets in hadronic collisions

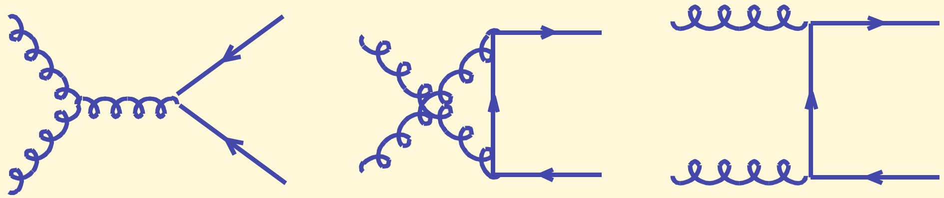


- Inclusive production of jets is the largest component of high-Q phenomena in hadronic collisions
- QCD predictions are known up to NLO accuracy
- Intrinsic theoretical uncertainty (at NLO) is approximately 10%
- Uncertainty due to knowledge of parton densities varies from 5-10% (at low transverse momentum, p_T to 100% (at very high p_T, corresponding to high-x gluons)
- Jet are used as probes of the quark structure (possible substructure implies departures from point-like behaviour of cross-section), or as probes of new particles (peaks in the invariant mass distribution of jet pairs)

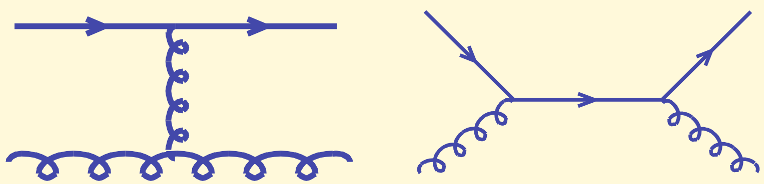
$gg \to gg$



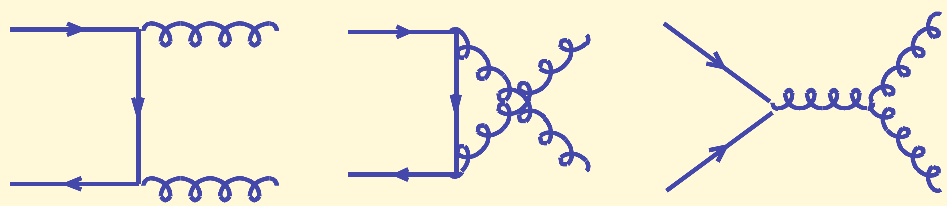
$gg \to q\bar{q}$



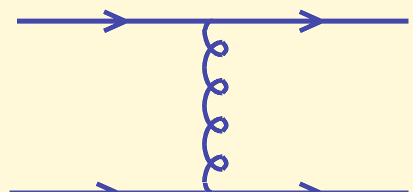
$qg \to qg$



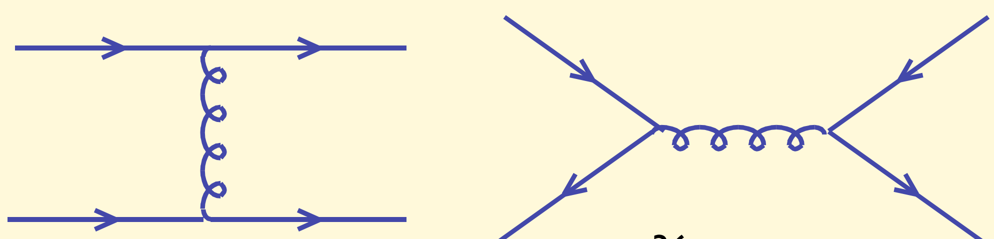
$q\bar{q} \to gg$



$qq' \to qq'$



$q\bar{q} \to q\bar{q}$



Phase space and cross-section for LO jet production

$$d[PS] = \frac{d^3p_1}{(2\pi)^2 2p_1^0} \frac{d^3p_2}{(2\pi)^2 2p_2^0} (2\pi)^4 \delta^4(P_{in} - P_{out})\, dx_1\, dx_2$$

(a) $$\delta(E_{in} - E_{out})\, \delta(P^z_{in} - P^z_{out})\, dx_1\, dx_2 = \frac{1}{2E^2_{beam}}$$

(b) $$\frac{dp^z}{p^0} = dy \equiv d\eta$$

$$\Rightarrow \quad d[PS] = \frac{1}{4\pi S} p_T\, dp_T\, d\eta_1\, d\eta_2$$

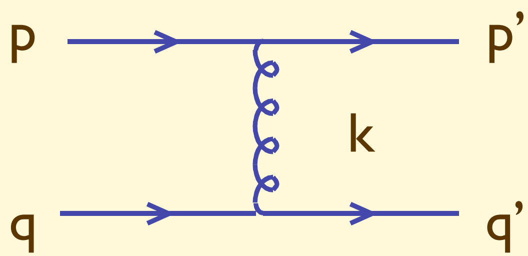
$$\Rightarrow \quad \frac{d^3\sigma}{dp_T d\eta_1 d\eta_2} = \frac{p_T}{4\pi S} \sum_{i,j} f_i(x_1) f_j(x_2) \frac{1}{2\hat{s}} \overline{\sum_{kl}} |M(ij \to kl)|^2$$

The measurement of pT and rapidities for a dijet final state uniquely determines the parton momenta x_1 and x_2. Knowledge of the partonic cross-section allows therefore the determination of partonic densities f(x)

Small-angle jet production, a useful approximation for the determination of the matrix elements and of the cross-section

At small scattering angle, $t = (p_1 - p_3)^2 \sim (1 - \cos\theta) \to 0$

and the $1/t^2$ propagators associated with t-channel gluon exchange dominate the matrix elements for all processes. In this limit it is easy to evaluate the matrix elements. For example:



$$\sim (\lambda^a)_{ij} (\lambda^a)_{kl} (2p_\mu) \frac{1}{t} (2q_\mu) = \frac{2s}{t} (\lambda^a)_{ij} (\lambda^a)_{kl}$$

where we used the fact that, for k=p-p'<<p (small angle scattering),

$$\bar{u}(p') \gamma_\mu u(p) \sim \bar{u}(p) \gamma_\mu u(p) = 2p_\mu$$

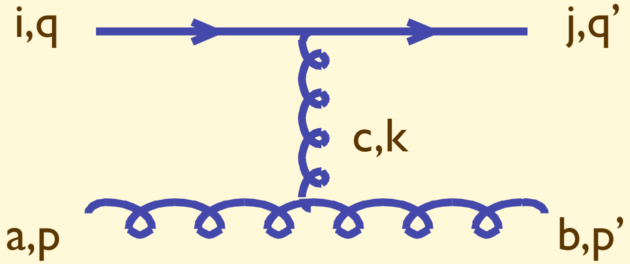
Using our colour algebra results, we then get: $\overline{\sum_{col,spin}} |M|^2 = \frac{1}{N_c^2} \frac{N_c^2 - 1}{4} \frac{4s^2}{t^2}$

Noting that the result must be symmetric under s↔u exchange, and setting Nc=3, we finally obtain: $\overline{\sum_{col,spin}} |M|^2 = \frac{4}{9} \frac{s^2 + u^2}{t^2}$

which turns out to be the exact result!

Quark-gluon and gluon-gluon scattering

We repeat the exercise in the more complex case of qg scattering, assuming the dominance of the t-channel gluon-exchange diagram:



$$\sim f^{abc}\lambda^c_{ij} 2p_\mu \frac{1}{t} 2q_\mu = 2\frac{s}{t} f^{abc}\lambda^c_{ij}$$

Using the colour algebra results, and enforcing the s↔u symmetry, we get:

$$\overline{\sum_{col,spin}} |M|^2 = \frac{s^2+u^2}{t^2}$$

which differs by only 20% from the exact result

$$\overline{\sum_{col,spin}} |M|^2 = \frac{s^2+u^2}{t^2} - \frac{4}{9}\frac{s^2+u^2}{us}$$

even in the large-angle region, at 90°

In a similar way we obtain for gg scattering (using the t↔u symmetry):

$$\overline{\sum_{col,spin}} |M(gg\to gg)|^2 = \frac{9}{2}\left(\frac{s^2}{t^2}+\frac{s^2}{u^2}\right)$$

compared to the exact result

$$\overline{\sum_{col,spin}} |M(gg\to gg)|^2 = \frac{9}{2}\left(3-\frac{ut}{s^2}-\frac{us}{t^2}-\frac{st}{u^2}\right)$$

with a 20% difference at 90°

Note that in the leading 1/t approximation we get the following result:

$$\hat{\sigma}_{gg} : \hat{\sigma}_{qg} : \hat{\sigma}_{qq} = \frac{9}{4} : 1 : \frac{4}{9}$$

where 4/9 = C_F / C_A = [(N^2-1)/2N] / N is the ratio of the squared colour charges of quarks and gluons

and therefore

$$d\sigma_{jet} = \int dx_1 dx_2 \sum_{ij} f_i(x_1) f_j(x_2) d\hat{\sigma}_{ij} = \int dx_1 dx_2 \sum_{ij} F(x_1) F(x_2) d\hat{\sigma}_{gg}$$

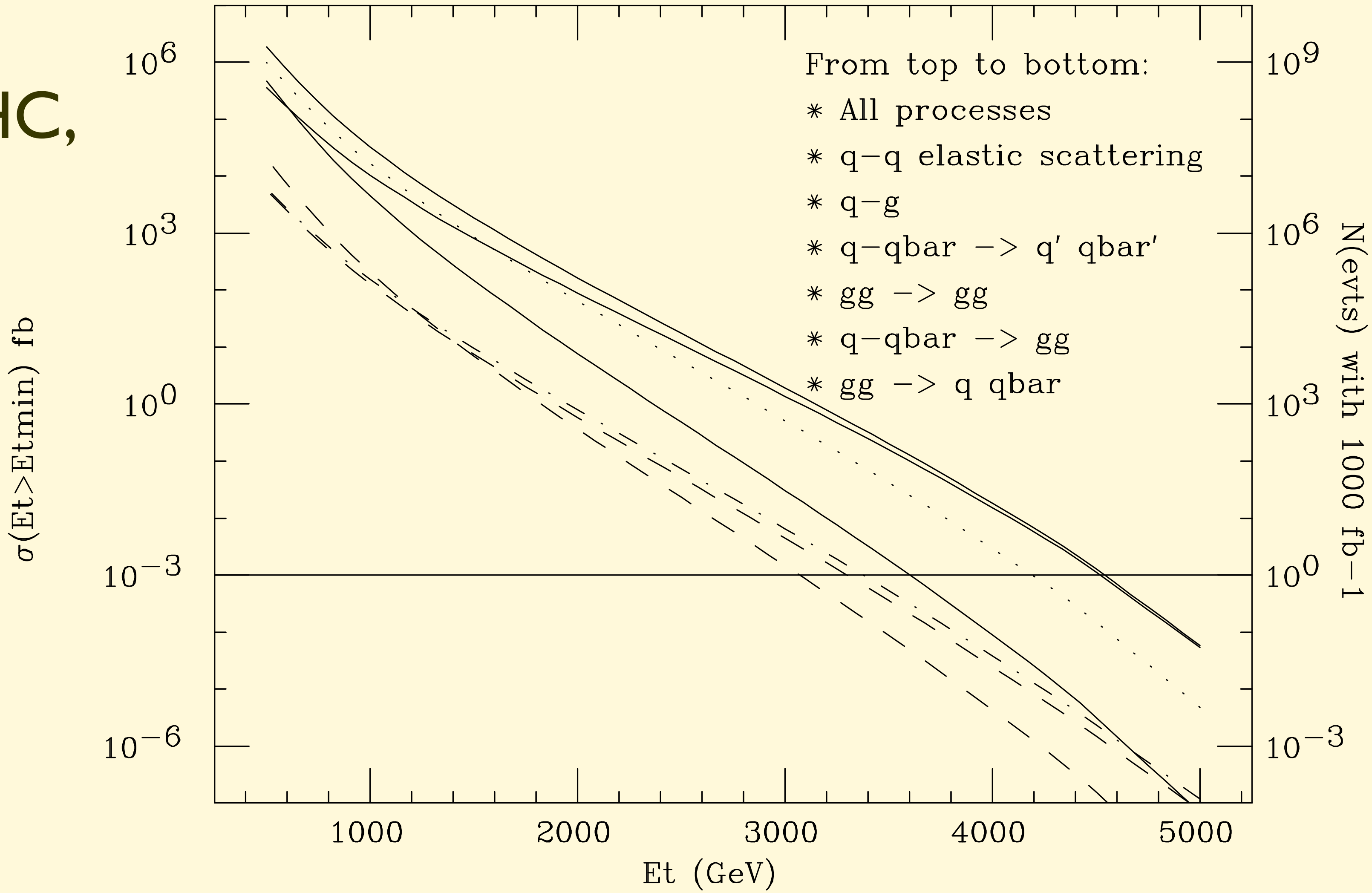
where we defined the `effective parton density' F(x):

$$F(x) = g(x) + \frac{4}{9} \sum_i [q_i(x) + \bar{q}_i(x)]$$

As a result jet data cannot be used to extract separately gluon and quark densities. On the other hand, assuming an accurate knowledge of the quark densities (say from HERA), jet data can help in the determination of the gluon density

Process	$\frac{d\hat{\sigma}}{d\Phi_2}$	at 90°
$qq' \to qq'$	$\frac{4}{9}\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}$	2.22
$qq \to qq$	$\left[\frac{4}{9}\left(\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}+\frac{\hat{s}^2+\hat{t}^2}{\hat{u}^2}\right)-\frac{8}{27}\frac{\hat{s}^2}{\hat{u}\hat{t}}\right]$	3.26
$q\bar{q} \to q'\bar{q}'$	$\frac{4}{9}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}$	0.22
$q\bar{q} \to q\bar{q}$	$\left[\frac{4}{9}\left(\frac{\hat{s}^2+\hat{u}^2}{\hat{t}^2}+\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}\right)-\frac{8}{27}\frac{\hat{u}^2}{\hat{s}\hat{t}}\right]$	2.59
$q\bar{q} \to gg$	$\left[\frac{32}{27}\frac{\hat{t}^2+\hat{u}^2}{\hat{t}\hat{u}}-\frac{8}{3}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}\right]$	1.04
$gg \to q\bar{q}$	$\left[\frac{1}{6}\frac{\hat{t}^2+\hat{u}^2}{\hat{t}\hat{u}}-\frac{3}{8}\frac{\hat{t}^2+\hat{u}^2}{\hat{s}^2}\right]$	0.15
$gq \to gq$	$\left[-\frac{4}{9}\frac{\hat{s}^2+\hat{u}^2}{\hat{s}\hat{u}}+\frac{\hat{u}^2+\hat{s}^2}{\hat{t}^2}\right]$	6.11
$gg \to gg$	$\frac{9}{2}\left(3-\frac{\hat{t}\hat{u}}{\hat{s}^2}-\frac{\hat{s}\hat{u}}{\hat{t}^2}-\frac{\hat{s}\hat{t}}{\hat{u}^2}\right)$	30.4

Jet production rates at the LHC, subprocess composition



The presence of a quark substructure would manifest itself via contact interactions (as in Fermi's theory of weak interactions). On one side these new interactions would lead to an increase in cross-section, on the other they would affect the jets' angular distributions. In the dijet CMF, QCD implies Rutherford law, and extra point-like interactions can then be isolated using a fit. With the anticipated statistics of 300 fb-1, limits on the scale of the new interactions in excess of 40 TeV should be reached (to increase to 60 TeV with 3000 fb-1)

Some more kinematics

Prove as an **exercise** that

$$x_{1,2} = \frac{p_T}{E_{beam}} \cosh y^* \, e^{\pm y_b}$$

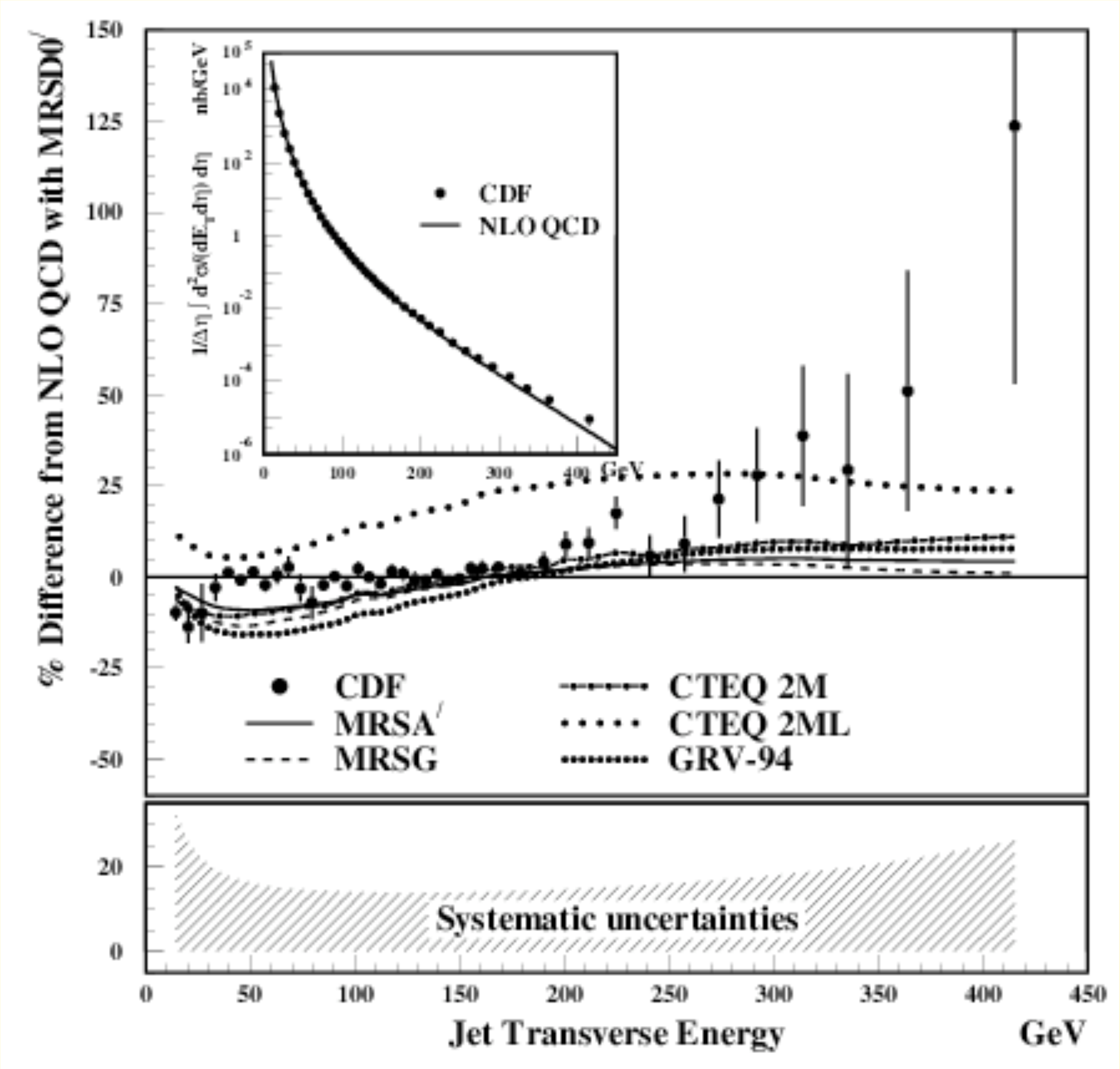
where

$$y^* = \frac{\eta_1 - \eta_2}{2}, \quad y_b = \frac{\eta_1 + \eta_2}{2}$$

We can therefore reach large values of x either by selecting large invariant mass events:

$$\frac{p_T}{E_{beam}} \cosh y^* \equiv \sqrt{\tau} \to 1$$

or by selecting low-mass events, but with large boosts (y_b large) in either positive of negative directions. In this case, we probe large-x with events where possible new physics is absent, thus setting consistent constraints on the behaviour of the cross-section in the high-mass region, which could hide new phenomena.



% Difference from NLO QCD with MRSD0′
150
125
100
75
50
25
0
-25
-50
20
0
CDF
NLO QCD
nb/GeV
1/Δη ∫ d²σ/(dE_T dη) dη
GeV
CDF
MRSA′
MRSG
CTEQ 2M
CTEQ 2ML
GRV-94
Systematic uncertainties
0 50 100 150 200 250 300 350 400 450
Jet Transverse Energy
GeV

Example, at the Tevatron