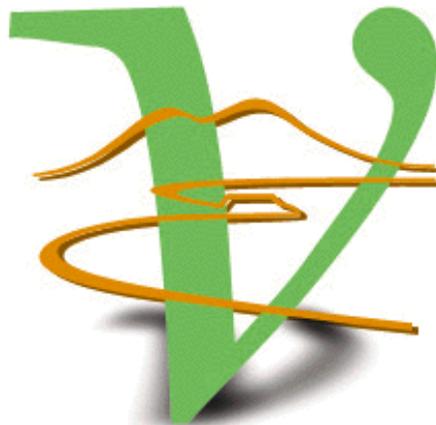


Neutrino Physics: experimental status, theoretical aspects and perspectives

Pasquale Migliozzi
INFN - Napoli



... neutrinos induce courage in
theoreticians and perseverance in
experimenters
Maurice Goldhaber, 1974

Charm physics with neutrinos

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Abstract

High energy neutrino interactions induce charmed hadron production at the level of a few percent and therefore they constitute a powerful tool to study charm physics. After 30 years of investigations with different neutrino beams and different detection techniques, important results have been achieved while other topics still need to be clarified. Recently, relevant results have been reported by several collaborations. We review the composite scenario of charm physics as it emerges from 30 years of investigations, including the latest results and pointing out possible future developments in this field.

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The potential for neutrino physics at muon colliders and dedicated high current muon storage rings

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Abstract

Conceptual design studies are underway for muon colliders and other high-current muon storage rings that have the potential to become the first true "neutrino factories". Muon decays in long straight sections of the storage rings would produce precisely characterized beams of electron and muon type neutrinos of unprecedented intensity. This article reviews the prospects for these facilities to greatly extend our capabilities for neutrino experiments, largely emphasizing the physics of neutrino interactions. © 2002 Elsevier Science B.V. All rights reserved.

PACS: 13.15.+g

Keywords: Muon colliders; Muon storage rings; Neutrino factories

PHYSICS AT THE FRONT-END OF A NEUTRINO FACTORY: A QUANTITATIVE APPRAISAL

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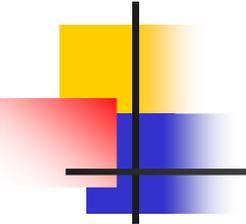
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Abstract

We present a quantitative appraisal of the physics potential for neutrino experiments at the front-end of a muon storage ring. We estimate the foreseeable accuracy in the determination of several interesting observables, and explore the consequences of these measurements. We discuss the extraction of individual quark and antiquark densities from polarized and unpolarized deep-inelastic scattering. In particular we study the implications for the understanding of the nucleon spin structure. We assess the determination of α_S from scaling violation of structure functions, and from sum rules, and the determination of $\sin^2 \theta_W$ from elastic νe and deep-inelastic νp scattering. We then consider the production of charmed hadrons, and the measurement of their absolute branching ratios. We study the polarization of Λ baryons produced in the current and target fragmentation regions. Finally, we discuss the sensitivity to physics beyond the Standard Model.

Neutrino Physics is not ONLY oscillations, but MUCH more. For a review of past results, present status and future activities, we refer to the abovementioned reviews.



Outline

- Neutrino oscillation formalism
- Experimental status: solar, atmospheric, reactor and accelerator data
- Interpretation of all available data
- Key measurements with neutrinos
- Future projects
- Conclusion

The PMNS leptonic mixing matrix

If neutrinos have mass: $|\nu_l\rangle = \sum U_{li} |\nu_i\rangle$

Flavour e.s.

Mass e.s.

For three neutrinos

$$U_{li} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

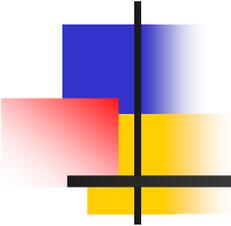
where $c_{ij} = \cos \theta_{ij}$, and $s_{ij} = \sin \theta_{ij}$

Oscillation probability

$$P_{\nu_\alpha \rightarrow \nu_\beta}(L) = \sum_{k=1}^3 |U_{\alpha k}|^2 |U_{\beta k}|^2 + 2 \operatorname{Re} \sum_{\substack{k=1 \\ k>j}}^3 U_{\alpha k}^* U_{\beta k} U_{\alpha j} U_{\beta j}^* \exp\left(-i \frac{\Delta m_k^2 L}{2E}\right)$$

Fixed by nature

Tuned by experiments



The MSW effect

MSW is such a beautiful phenomenon that Nature would be well advised to use it. After all, it may eventually give us the unambiguous, incontrovertible, uncontestable, clear and definitive evidence we so eagerly seek that the neutrino has mass

S.P. Rose, 1986

Mikheyev-Smirnov-Wolfenstein (MSW)

2-ν vacuum oscillation

$$P(\nu_e \rightarrow \nu_e) = 1 - \sin^2 2\theta \sin^2 \left(1.27 \Delta m^2_{[\text{eV}^2]} \frac{L[\text{m}]}{E[\text{MeV}]} \right) \text{ where } \Delta m^2 = m_2^2 - m_1^2$$

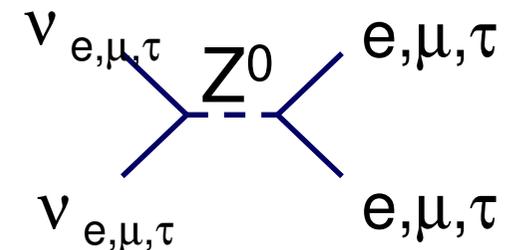
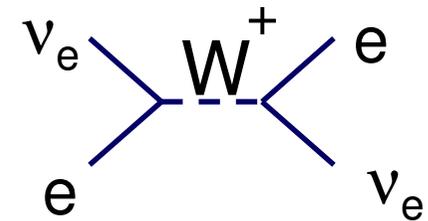
- In matter, ν_e and ν_μ/ν_τ have different "effective masses", flavor conversion can be a resonant effect \Rightarrow

Mikheyev-Smirnov-Wolfenstein (MSW)

Mixing in matter

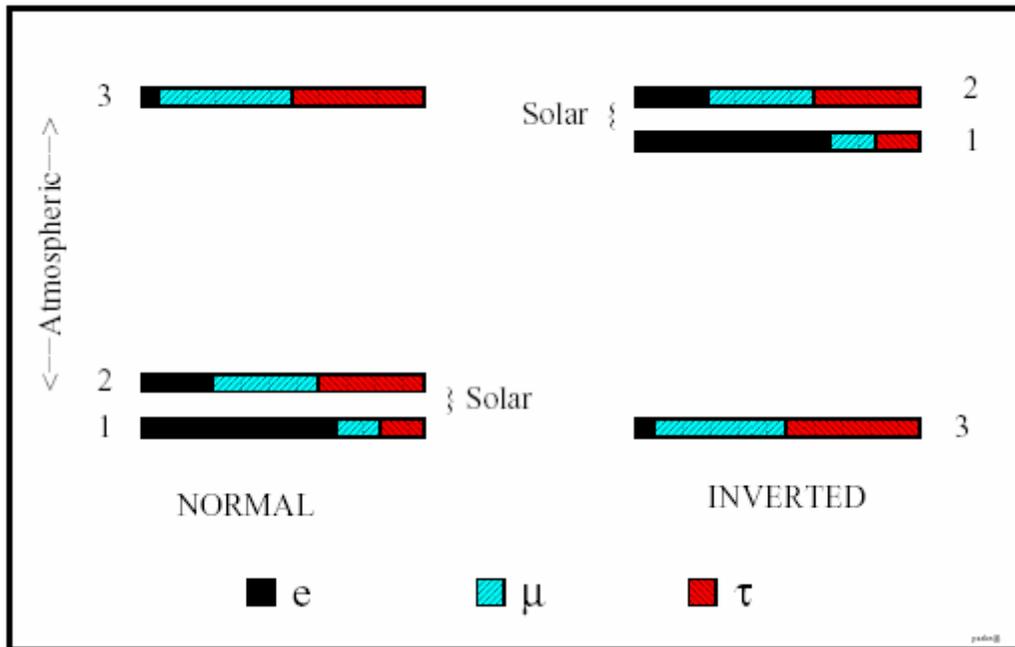
Effective Mixing Angle in Matter: $\tan 2\theta_M = \frac{\tan 2\theta}{1 - \frac{A_{CC}}{\Delta m^2 \cos 2\theta}}$

Resonance: $A_{CC}^R = \Delta m^2 \cos 2\theta \Rightarrow N_e^R = \frac{\Delta m^2 \cos 2\theta}{2\sqrt{2}EG_F}$



Notation

- Mixing parameters: $U = U(\theta_{12}, \theta_{13}, \theta_{23}, \delta)$ as for CKM matrix
- Mass-gap parameters: $M^2 = \Delta m_{12}^2, \pm \Delta m_{23}^2$



The absolute mass scale should be set by other measurements:

- β -decay
- $0\nu 2\beta$ -decay
- anisotropies in cosmic background radiation

Appearance channels: $\nu_\mu \rightarrow \nu_e$

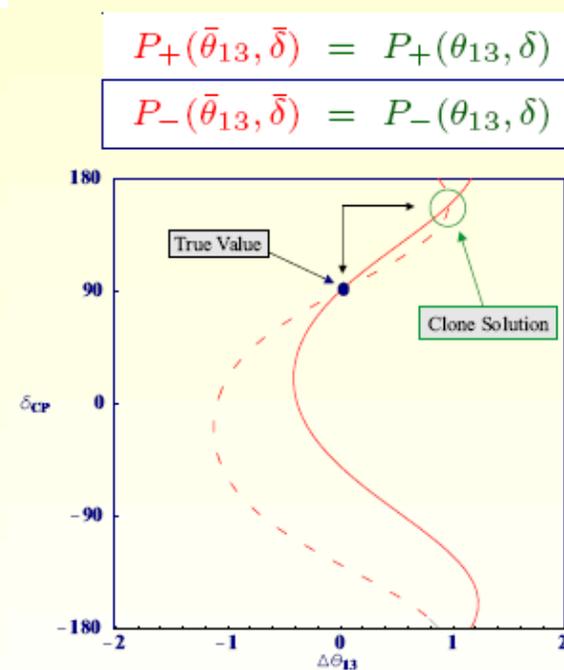
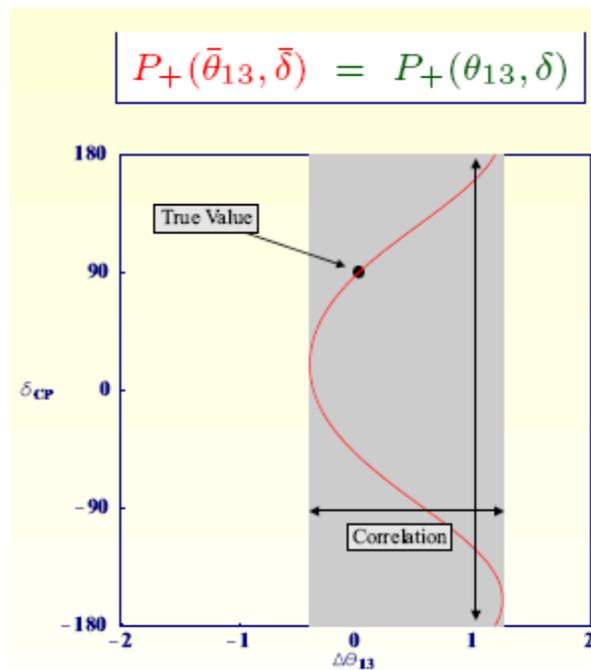
$$\begin{aligned}
 P_{\text{app}} &\simeq \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2[(1 - \hat{A})\Delta]}{(1 - \hat{A})^2} \\
 &\pm \alpha \sin 2\theta_{13} \xi \sin \delta_{\text{CP}} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin[(1 - \hat{A})\Delta]}{\hat{A} (1 - \hat{A})} \\
 &+ \alpha \sin 2\theta_{13} \xi \cos \delta_{\text{CP}} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin[(1 - \hat{A})\Delta]}{\hat{A} (1 - \hat{A})} \\
 &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2},
 \end{aligned}$$

$$\alpha \equiv \frac{\Delta m_{21}^2}{\Delta m_{31}^2} \simeq \pm 0.03, \quad \Delta \equiv \frac{\Delta m_{31}^2 L}{4E}, \quad \xi \equiv \sin 2\theta_{12} \sin 2\theta_{23}, \quad \hat{A} \equiv \pm \frac{2\sqrt{2}G_F n_e E}{\Delta m_{31}^2}$$

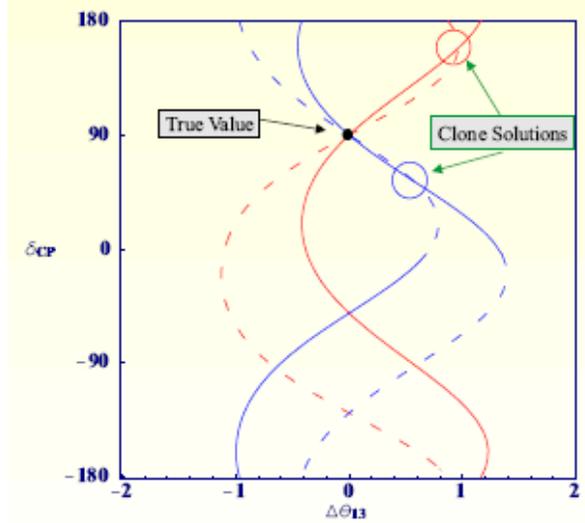
- Complicated, but all interesting information there: θ_{13} , δ_{CP} , mass hierarchy (via A)

The intrinsic degeneracy

There is a strong correlation between θ_{13} and δ



Measuring the $\nu_e \rightarrow \nu_\mu$ **GOLDEN** Appearance Channel;
 Measuring the $\nu_e \rightarrow \nu_\tau$ **SILVER** Appearance Channel;



There are infinite solutions!
 Infinite degeneracies

By using neutrinos and anti-neutrinos there are two solutions: the true and the clone

No clone solutions!

- Degeneracies in (θ_{13}, δ) Measure: **EIGHTFOLD DEGENERACY**

Besides θ_{13} and δ other two (discrete) quantities will be unknown in 5-10 years at the time of next generation neutrino experiments:

- The SIGN of the ATM mass difference $s_{atm} = \text{sign}(\Delta m_{23}^2)$
- The OCTANT of the ATM angle $s_{oct} = \text{sign}(\tan 2\theta_{23})$

Consequently, for taking into account **ALL OUR IGNORANCE** on the neutrino **masses** and **mixings** one has to make a simultaneous fit to these 4 parameters, i.e. to solve the following equation:

$$N_i^\pm \underbrace{(\bar{\theta}_{13}, \bar{\delta}; \bar{s}_{atm}, \bar{s}_{oct})}_{\text{“true parameters”}} = N_i^\pm \underbrace{(\theta_{13}, \delta; s_{atm}, s_{oct})}_{\text{“guessed parameters”}}$$

One has to solve ALL the following FOUR systems of equations, each of them having in general two distinct solutions:

intrinsic degeneracy (Burguet01)

$$N_i^\pm(\bar{\theta}_{13}, \bar{\delta}; \bar{s}_{atm}, \bar{s}_{oct}) = N_i^\pm(\theta_{13}, \delta; s_{atm} = \bar{s}_{atm}, s_{oct} = \bar{s}_{oct})$$

sign degeneracy (Minakata01)

$$N_i^\pm(\bar{\theta}_{13}, \bar{\delta}; \bar{s}_{atm}, \bar{s}_{oct}) = N_i^\pm(\theta_{13}, \delta; s_{atm} = -\bar{s}_{atm}, s_{oct} = \bar{s}_{oct})$$

octant degeneracy (Fogli96, Barger01)

$$N_i^\pm(\bar{\theta}_{13}, \bar{\delta}; \bar{s}_{atm}, \bar{s}_{oct}) = N_i^\pm(\theta_{13}, \delta; s_{atm} = \bar{s}_{atm}, s_{oct} = -\bar{s}_{oct})$$

mixed degeneracy (Barger01)

$$N_i^\pm(\bar{\theta}_{13}, \bar{\delta}; \bar{s}_{atm}, \bar{s}_{oct}) = N_i^\pm(\theta_{13}, \delta; s_{atm} = -\bar{s}_{atm}, s_{oct} = -\bar{s}_{oct})$$

The Eightfold Degeneracy

The disappearance channels

$$1 - P_{ee} \cong \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta m_{\text{atm}}^2 L}{4E} \right) + \underbrace{\sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta m_{\text{solar}}^2 L}{4E} \right)}_{\text{negligible on the atm. peak}} + \dots$$

\uparrow
 Survival probability

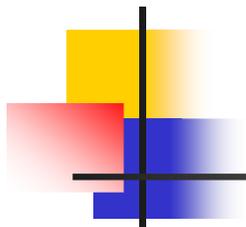
$(\Delta m_{\text{atm}}^2 \gg \Delta m_{\text{solar}}^2)$

- No sensitivity to CP phase, to θ_{23} and to the sign of Δm_{23}^2 : the eightfold degeneracy is not an issue
- Drawback: only one parameter can be measured
- Difficult measurement: extremely sensitive to the knowledge of the flux, of the signal and of the background

$$P_{\nu_\mu \nu_\mu}^\pm \approx 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{23}^2 L}{4E} \right) + \mathcal{O}(\theta_{13}^2 \sin^2 (\Delta m_{23}^2 L/4E))$$

$$+ \mathcal{O}(\cos \delta_{CP} \cdot \theta_{13} \cdot \Delta_{12} \cdot \sin(\Delta m_{23}^2 L/4E)) + \mathcal{O}(\Delta_{12}^2)$$

- Extremely sensitive measurement of θ_{23} and Δm_{23}^2
- No sensitivity at all to θ_{13} and to CP phase
- Difficult measurement: extremely sensitive to the knowledge of the flux, of the signal and of the background



SOLAR NEUTRINOS

Data taking: 1970 – 1995 108 solar runs.

Results:

ν interaction rate on ^{37}Cl : **2.56 +/- 0.16 (stat) +/- 0.16 (sys) SNU**

R (exp/SSM) = 0.34 +/- 0.03 (exp) +/- 0.05 (theo)

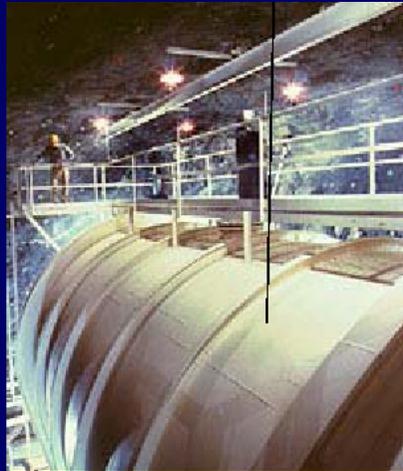
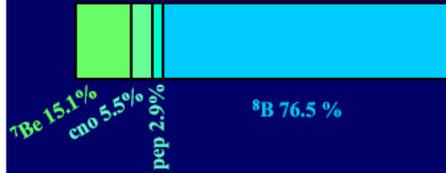
FIRST EVIDENCE FOR NEUTRINOS COMING FROM THE SUN

Neutrino interaction: $^{37}\text{Cl} (\nu_e, e) ^{37}\text{Ar}$
Threshold : 814 keV

Detection Technique: radiochemical

^{37}Ar is extracted from the tank by He purging every two months, and then counted inside a gas proportional counter

Signal composition:



Results: (Runs 1-108) (May 91- Jan 02)
70.8 +/- 4.5 (stat) +/- 3.8 (sys) SNU
R (exp/SSM) = 0.55 +/- 0.05 (exp) +/- 0.03 (theo)

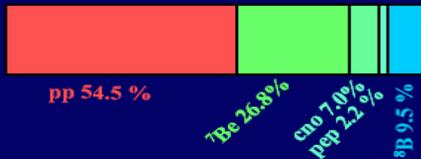
Results: (Runs 1 – 104) (Jan 90 – Jan 03)
70.5 +/- 4.8 (stat) +/- 3.5 (sys) SNU
R (exp/SSM) = 0.55 +/- 0.05 (exp) +/- 0.03 (theo)

FIRST DETECTION OF PP SOLAR NEUTRINOS

SUPPRESSION OF INTERMEDIATE ENERGY SOLAR NEUTRINO FLUX

Neutrino interaction: $^{71}\text{Ga} (\nu_e, e) ^{71}\text{Ge}$ Threshold : 233 keV

Signal composition:



GALLEX/GNO

Location:
Laboratori Nazionali del Gran Sasso
Abruzzo (Italy)

Target: 101 tons of GaCl_3 acidic solution
(30 tons of nat. Ga)

SAGE

Location:
Baksan Neutrino Observatory
Caucasus (Russia)

Target: 55 tons of metallic Ga

❖ **Location:** Kamioka mine, Japan, 2700 mwe depth

❖ **Target:** 4500 tons of pure water
2150 tons (fid. Vol.)

❖ **Neutrino interaction:** $e + \nu_x \rightarrow e + \nu_x$ (CC+NC)

❖ **Detection Technique:** Cerenkov light of scattered electron

❖ **Signal composition:** 100% from $^8\text{B} \nu$
Direction of recoil electron
Energy spectrum of recoil electron
Threshold : 7 MeV

❖ **Data taking:** 1987 → 1995
2079 days of live-time.

Results : 597 ν events observed in 2079 days

$\Phi(^8\text{B}) = 2.80 +/- 0.19$ (stat) +/- 0.33 (sys) $10^6 \text{ cm}^{-2} \text{ s}^{-1}$

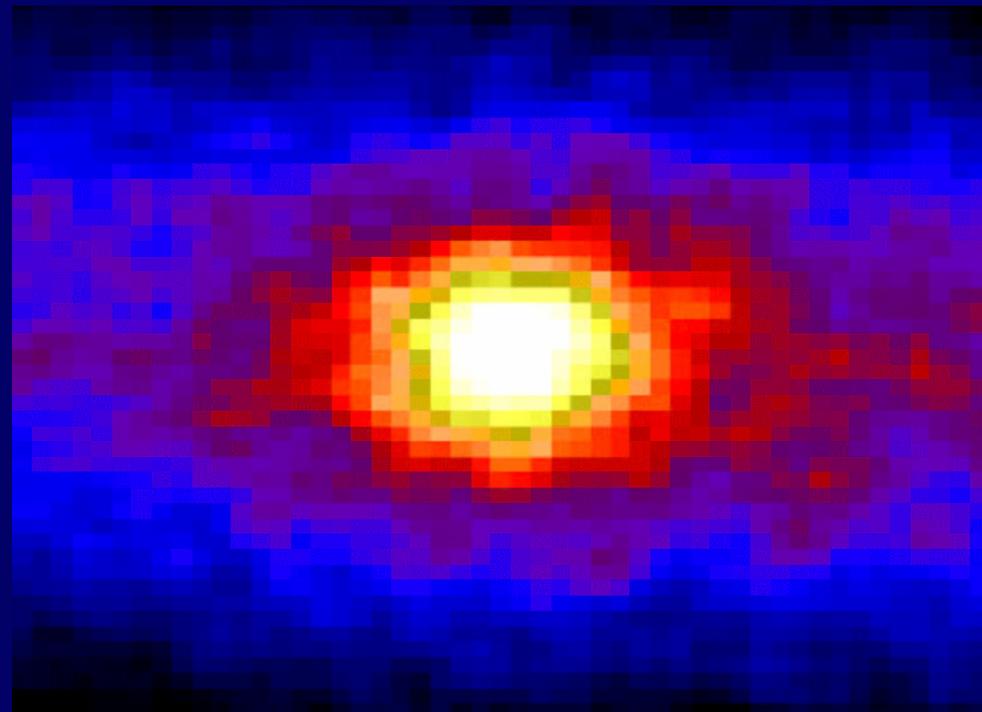
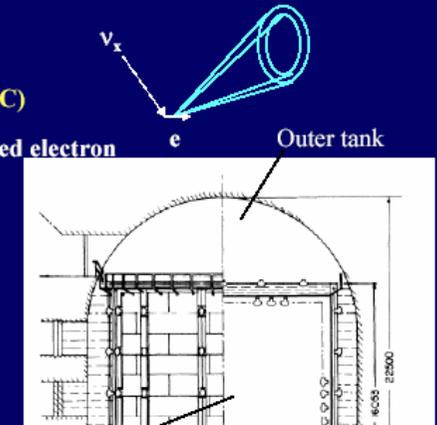
R (exp/SSM) = 0.37

Confirmation of the evidence for a “solar neutrino problem”

First evidence that neutrino signals are correlated with the Sun direction

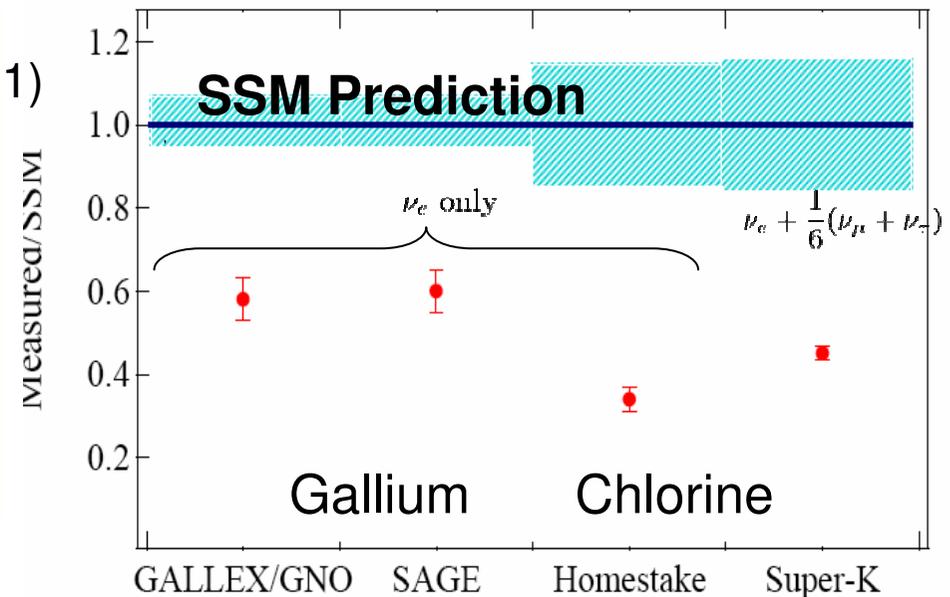
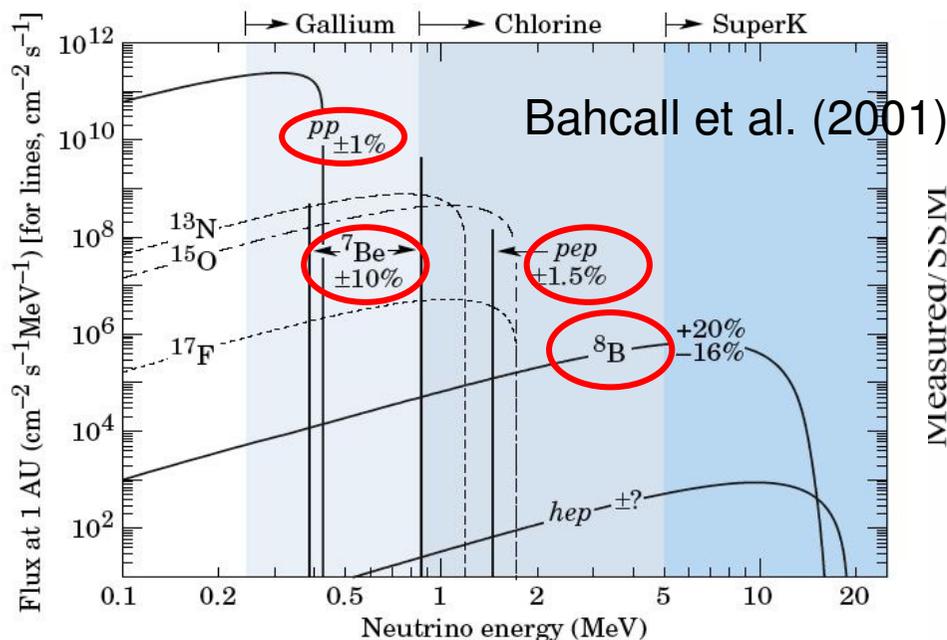
First test of the ^8B neutrino energy spectrum

First measurement of solar neutrinos in real time



The old standing solar neutrino problem (before year 2000)

Consistent predictions of ν_e flux from a number of Standard Solar Models (e.g. Bahcall *et al.*, and Turck-Chieze *et al.*)

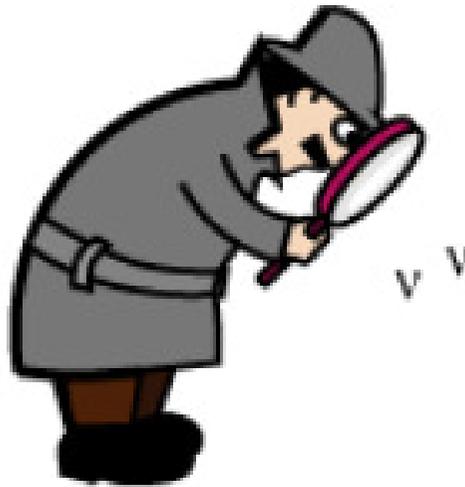


There is a strong deficit in the measured flux as well as an "energy" dependence of the deficit!!!

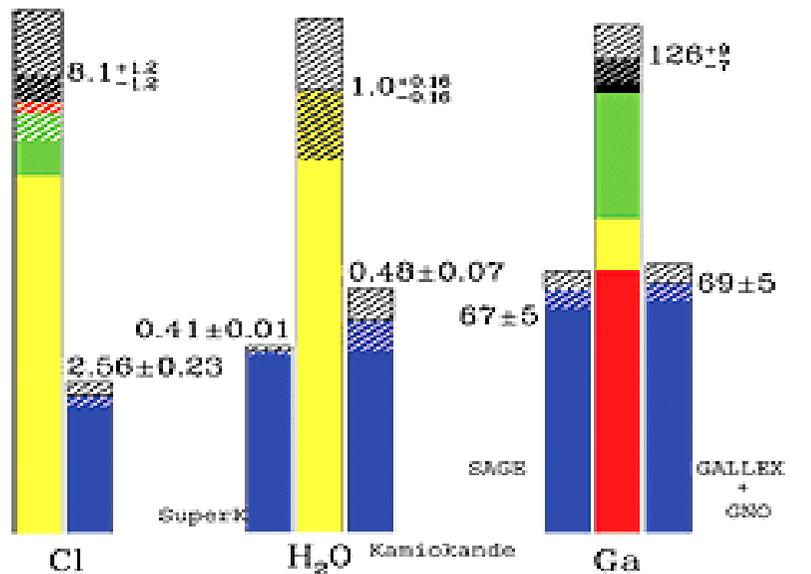
Although with a lower threshold Cl exp sees less ν than SK!!!

The Solar Neutrino Problem

(we don't get enough neutrinos)



Total Rates: Standard Model
Bahcall-Serenelli 2005 [1]



Theory ■ ⁷Be ■ P-P, pep

■ ⁸B ■ CNO

Experiments ■

Uncertainties

Data are consistent with:

- Full ν_e flux from $p + p \rightarrow e^+ + \nu_e + d$
- ~50% of the ν_e flux from $B^8 \rightarrow Be^8 + e^+ + \nu_e$
- Very strong (almost complete) suppression of the ν_e flux from $e^- + Be^7 \rightarrow \nu_e + Li^7$

The real solar neutrino puzzle:

There is evidence for B^8 in the Sun (with deficit 50%), but no evidence for Be^7 ; yet Be^7 is needed to make B^8 by the fusion reaction $p + Be^7 \rightarrow \gamma + B^8$

Possible solutions:

- At least one experiment is wrong
- The SSM is totally wrong
- The ν_e from $e^- + Be^7 \rightarrow \nu_e + Li^7$ are no longer ν_e when they reach the Earth and become invisible $\Rightarrow \nu_e$ OSCILLATIONS

ν reactions in SNO

ES



- Both SK, SNO
- Mainly sensitive to ν_e , less to ν_μ and ν_τ
- Strong directional sensitivity

CC



- Good measurement of ν_e energy spectrum
- Weak directional sensitivity $\propto 1 - 1/3 \cos(\theta)$

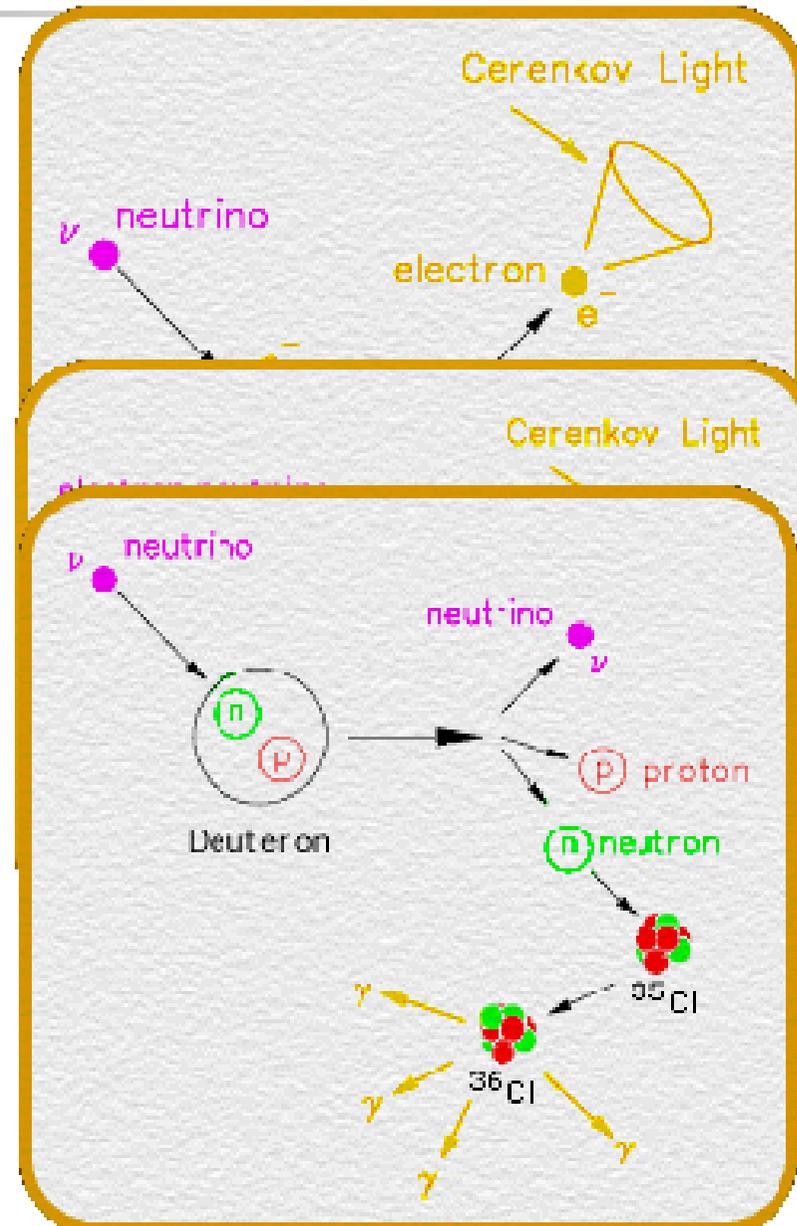
- ν_e ONLY

NC

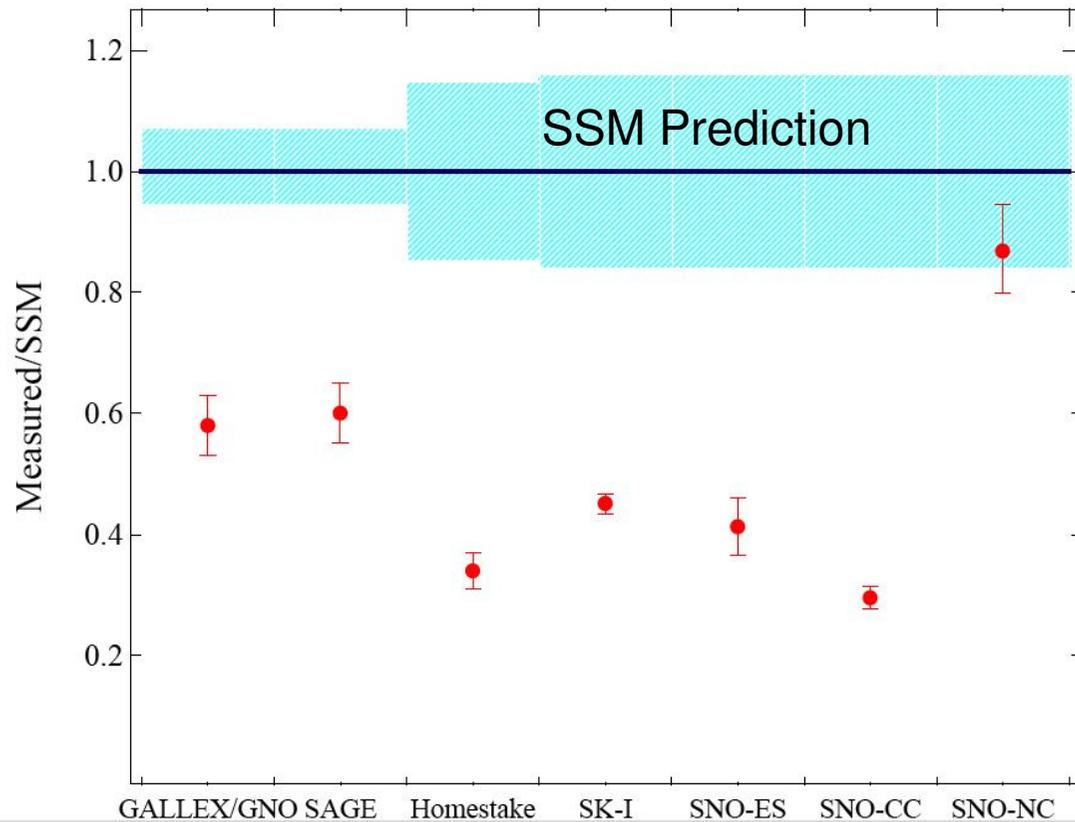


- Measure total ^8B ν flux from the sun.

- Equal cross section for all ν types



Resolution of the solar neutrino problem



SNO-measured “energy-unconstrained (no hypothesis on ^8B)” flux (Phase II) [$\times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$]:

$$\phi_{\text{CC}}^{\text{uncon}} = 1.68_{-0.06}^{+0.06} (\text{stat})_{-0.09}^{+0.08} (\text{syst})$$

$$\phi_{\text{ES}}^{\text{uncon}} = 2.35_{-0.22}^{+0.22} (\text{stat})_{-0.15}^{+0.15} (\text{syst})$$

$$\phi_{\text{NC}}^{\text{uncon}} = 4.94_{-0.21}^{+0.21} (\text{stat})_{-0.34}^{+0.38} (\text{syst})$$

arXiv:nucl-ex/0502021

Standard Solar Model flux
($\times 10^6 \text{ cm}^{-2} \text{ s}^{-1}$):

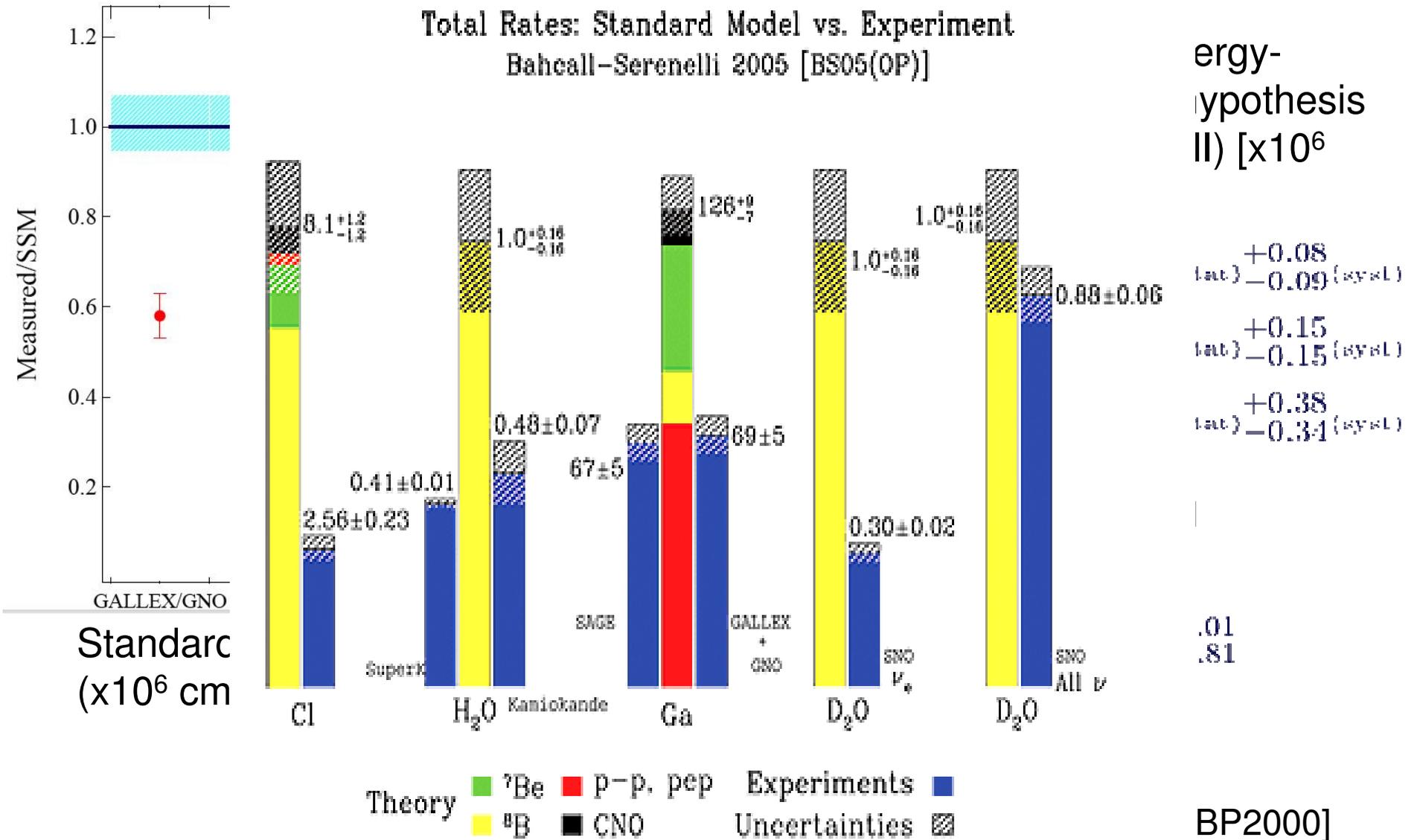
$$\text{Bahcall-Pinsonneault 2000} = 5.05_{-0.81}^{+1.01}$$

$$\text{Bahcall-Serenelli 2005} = 5.69 \pm 0.91$$

$$\text{Turck-Chieze 2004} = 5.31 \pm 0.60$$

$$\frac{\phi_{\text{CC}}^{\text{HICOE1}}}{\phi_{\text{NC}}^{\text{HICOE1}}} = 0.340 \pm 0.023 (\text{stat})_{-0.031}^{+0.029} (\text{syst}) \quad [\text{BP2000}]$$

Resolution of the solar neutrino problem



Indirect evidence for MSW effect

Study of the tolerance of the solar data for variations of standard MSW interaction energy

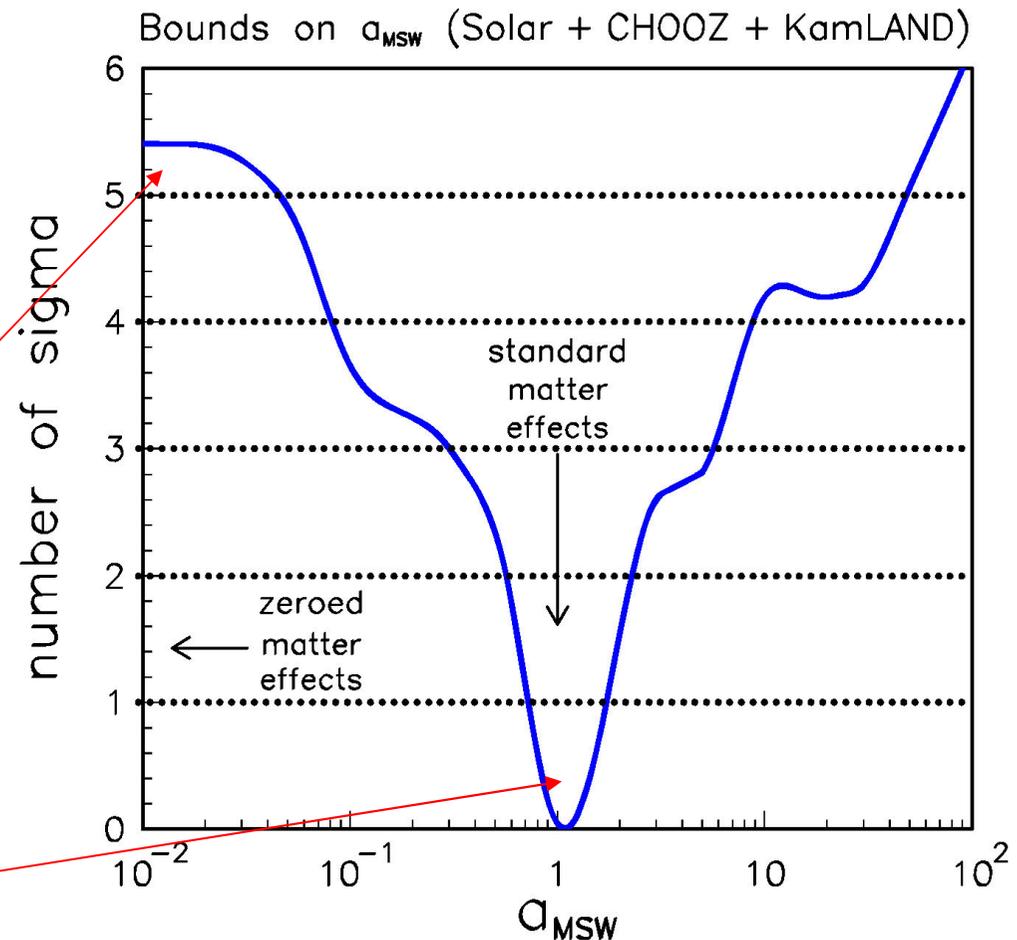
$$V = \sqrt{2} G_F N_e$$

through a shift

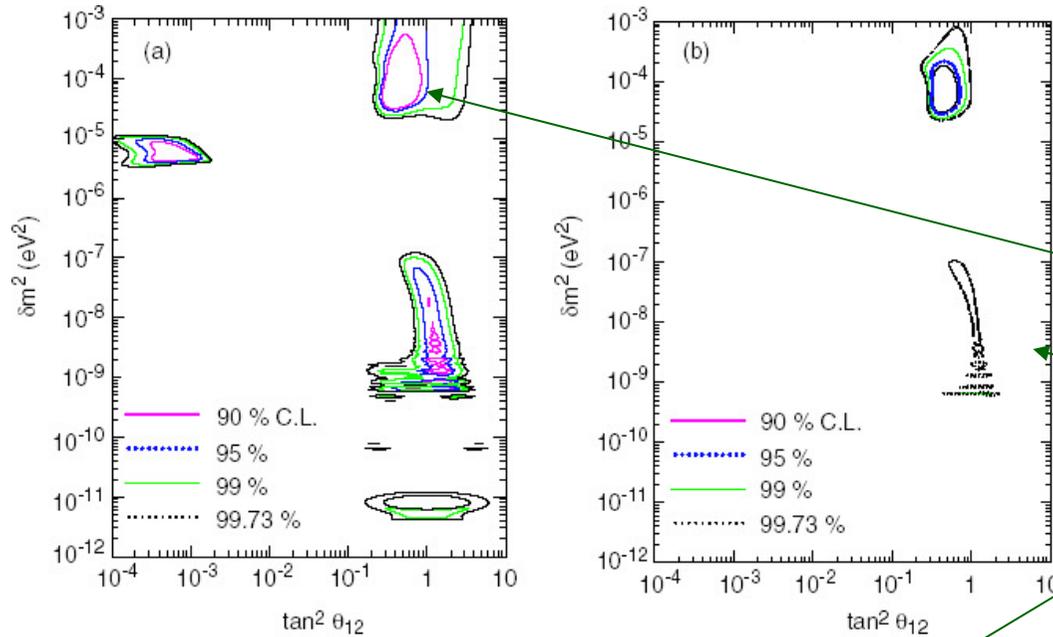
$$V \rightarrow V \cdot a_{MSW}$$

Case of no matter effects in the Sun is ruled out at >5 sigma.

Clear indication in favor of standard matter effects ($a_{MSW} = 1$).



What do we know about ν mixing from solar sector?

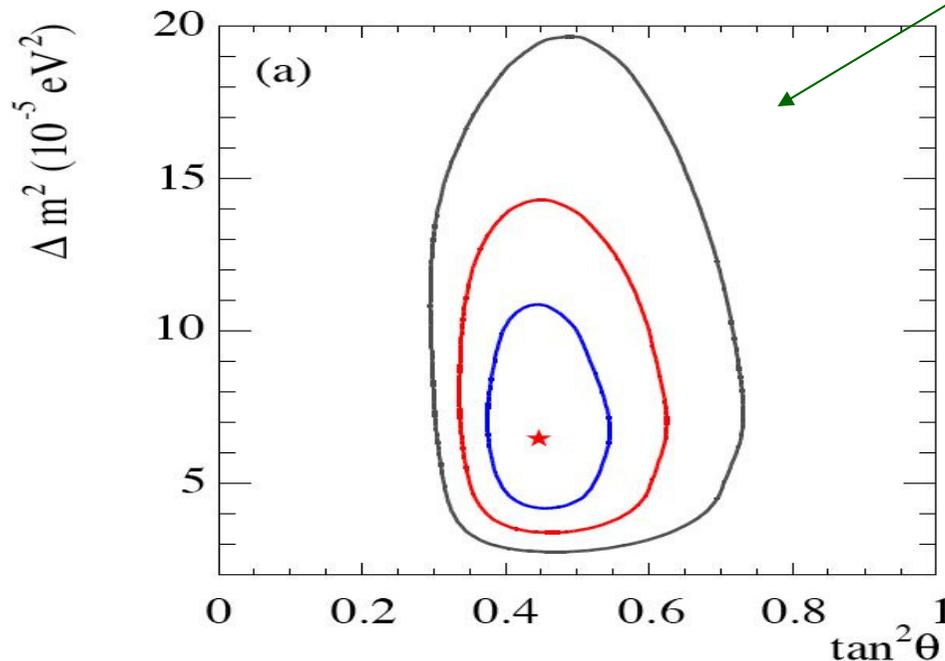


Dramatic reduction of the $(\delta m^2, \theta_{12})$ param. space in **2001-2003**
(note change of scales)

Cl+Ga+SK (2001)

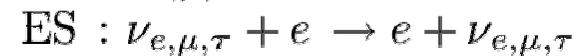
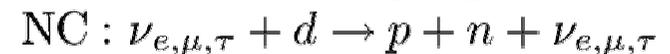
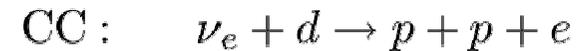
+SNO-I (2001-2002)

+SNO-II (2003)



(+ confirmation of solar model)

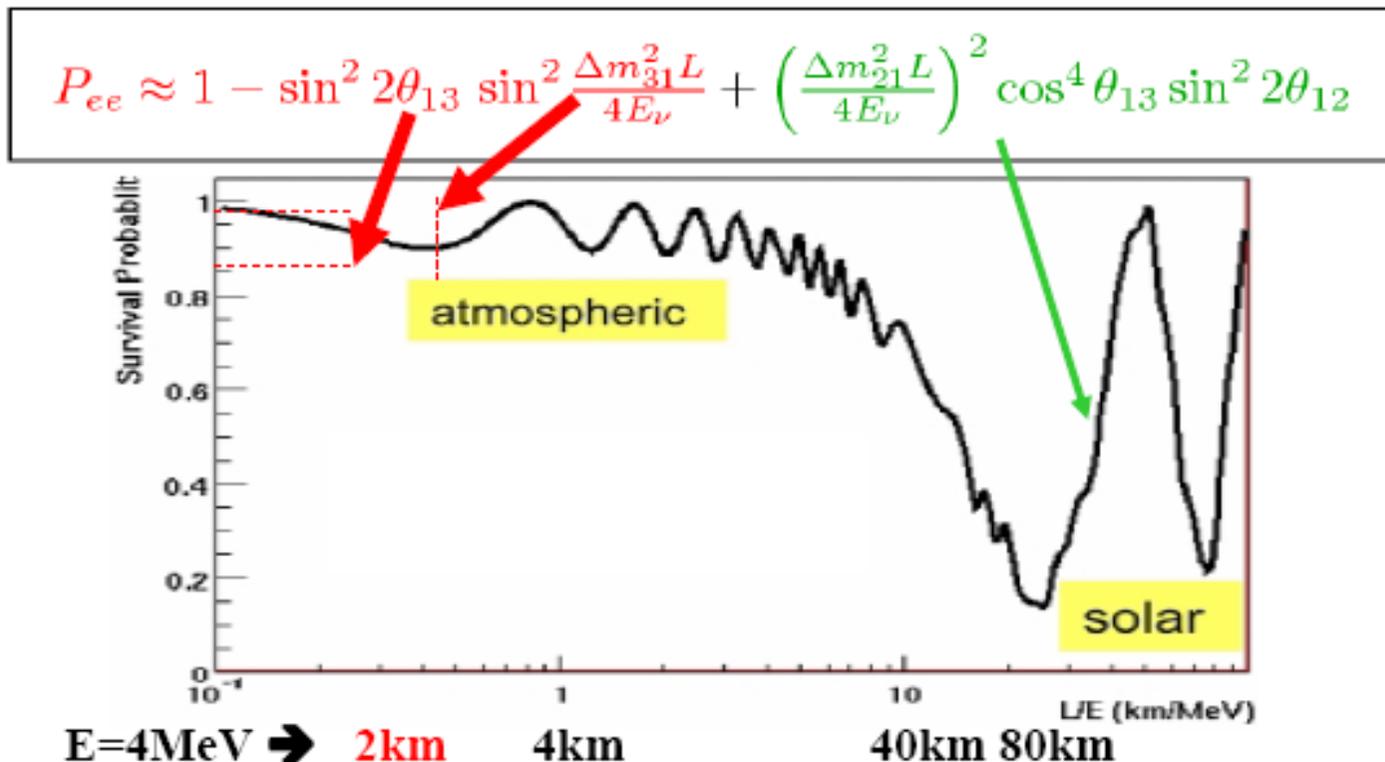
Direct proof of solar $\nu_e \rightarrow \nu_{\mu,\tau}$ in SNO through comparison of



From solar ν to reactor anti- $\bar{\nu}$?

CPT Theorem

Solar ν_e Reactor $\bar{\nu}_e$

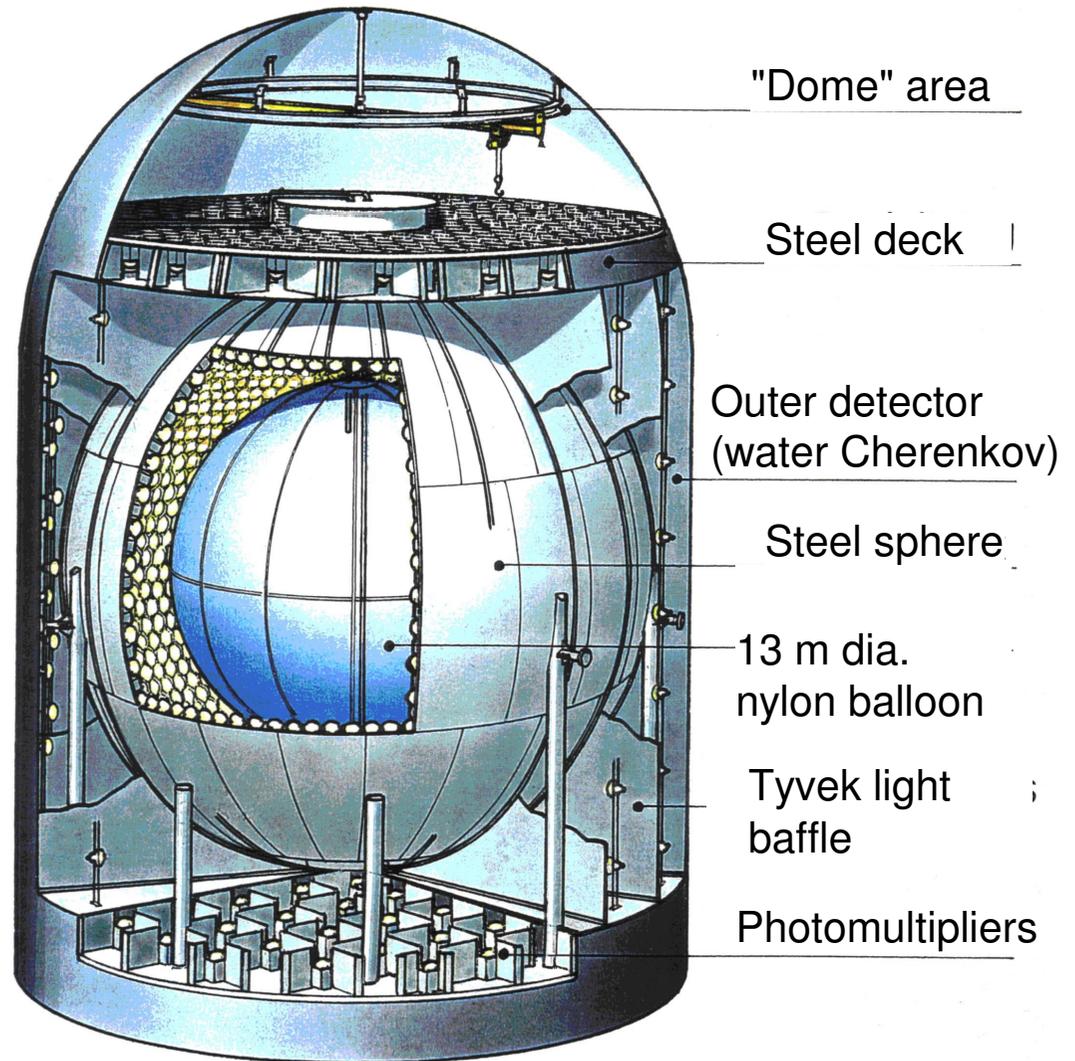
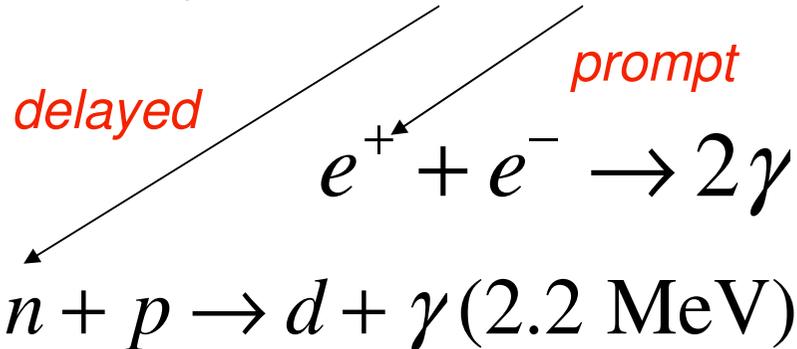
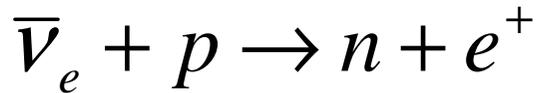


To test the hypothesis of neutrino oscillation as the underlying mechanism for flavor transformation, we need a baseline of $\sim 100\text{-}200$ km for reactor anti-neutrino experiments

The KamLAND Detector

- 70 GW (7% of world total) is generated at 130-220 km distance from Kamioka.
"Effective baseline" ~ 180 km

- Detects electron anti- ν by inverse β decay:

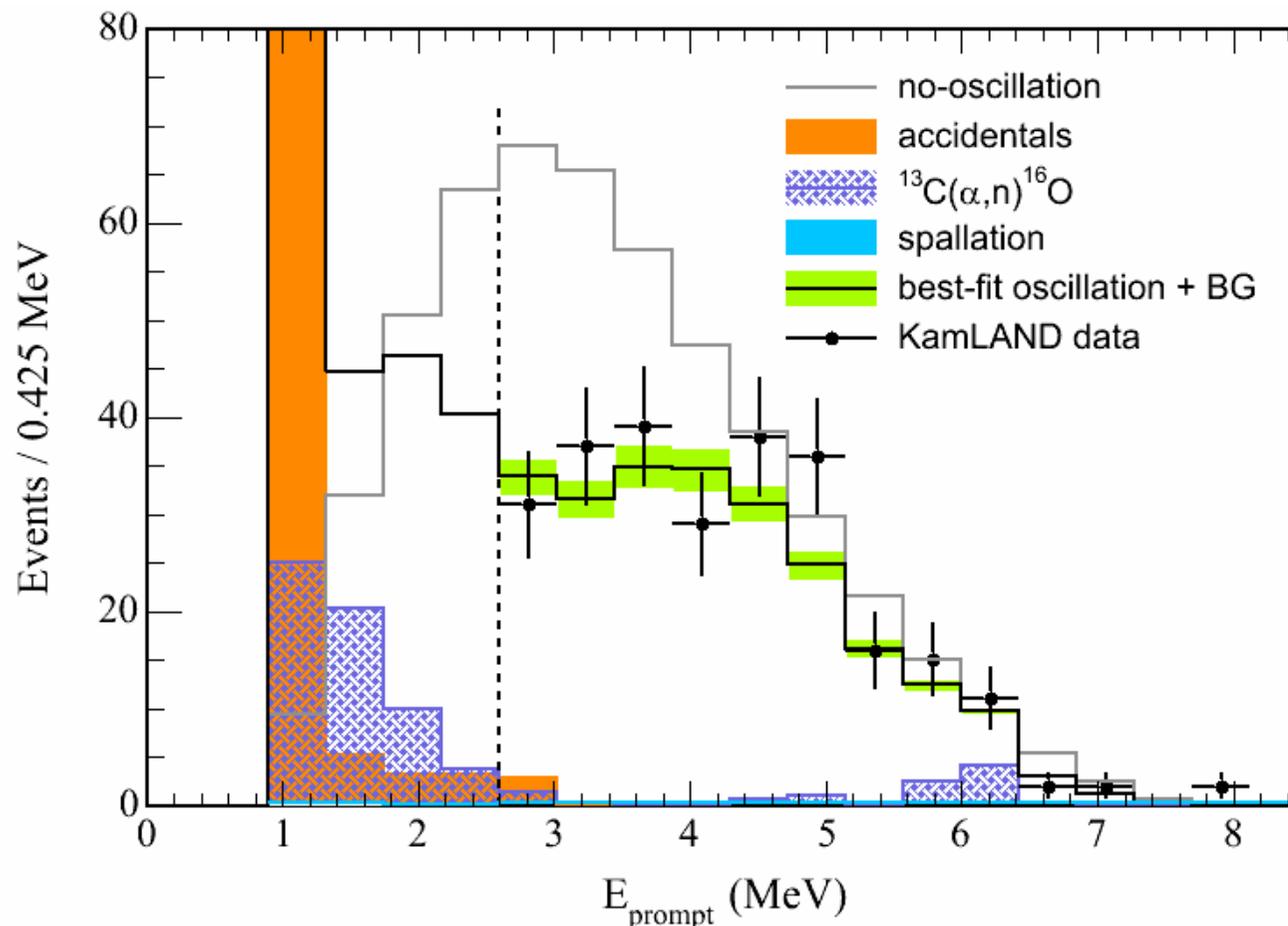


KamLAND: Signatures for Neutrino Oscillations

Energy spectrum Best-fit KamLAND only oscillation:

$$\tan^2 \theta = 0.46$$

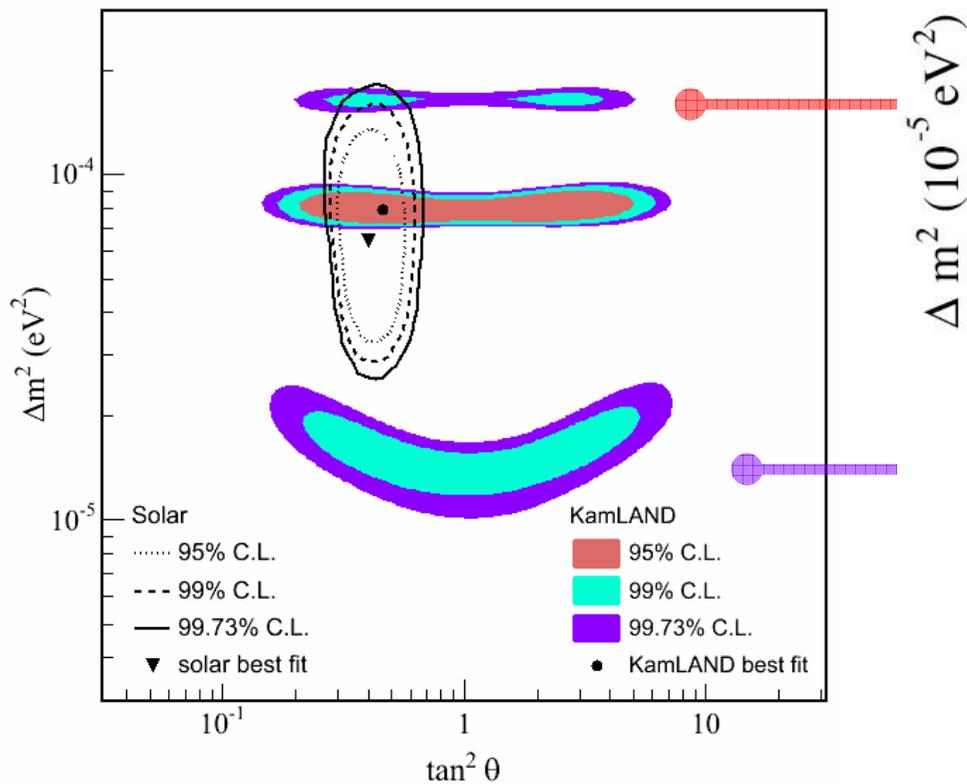
$$\Delta m^2 = 7.9_{-0.5}^{+0.6} \times 10^{-5} eV^2$$



A fit to a simple rescaled reactor spectrum is excluded at 99.6% CL

Solar+Reactor: What do we know about ν mixing now?

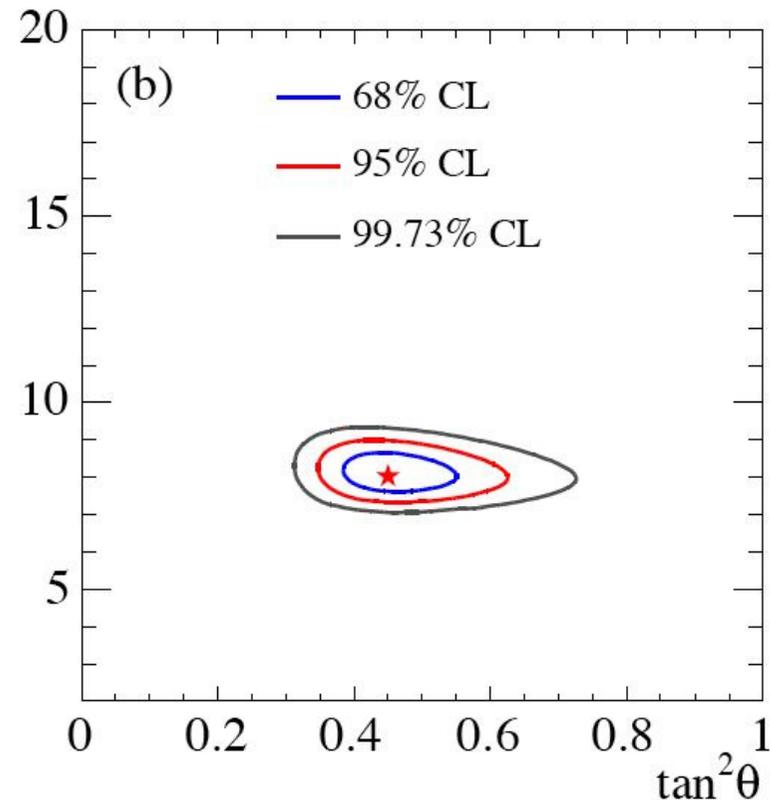
KamLAND-only



Best fit: $\Delta m^2 = 7.9_{-0.5}^{+0.6} \times 10^{-5} eV^2$
 $\tan^2 \theta = 0.46$

KamLAND $\rightarrow \Delta m_{12}^2$

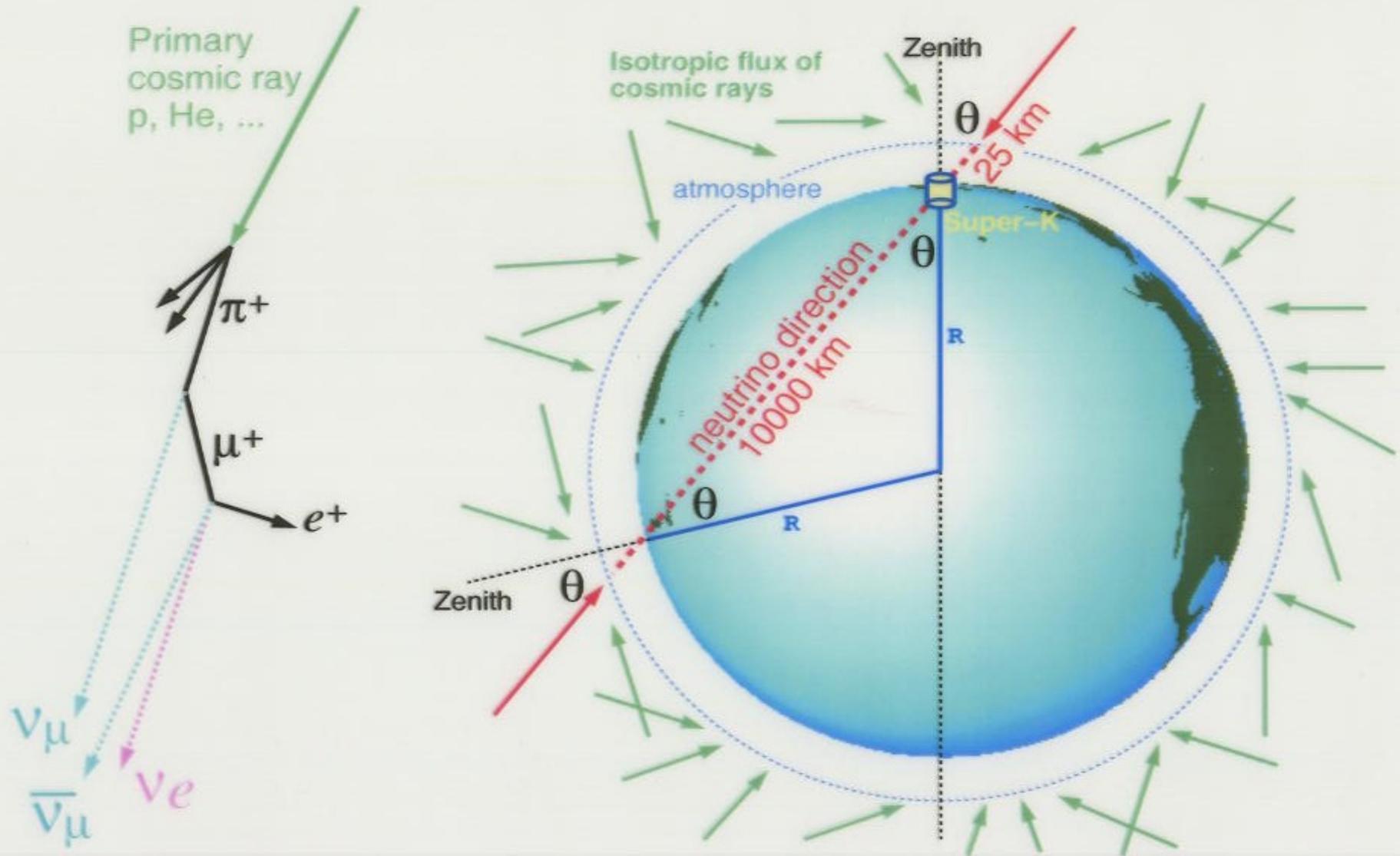
Solar+KamLAND



Best fit: $\Delta m^2 = 8.0_{-0.4}^{+0.6} \times 10^{-5} eV^2$
 $\tan^2 \theta = 0.45_{-0.07}^{+0.09}$

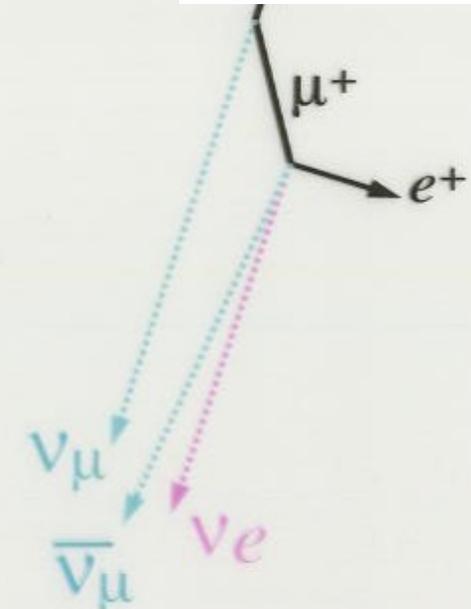
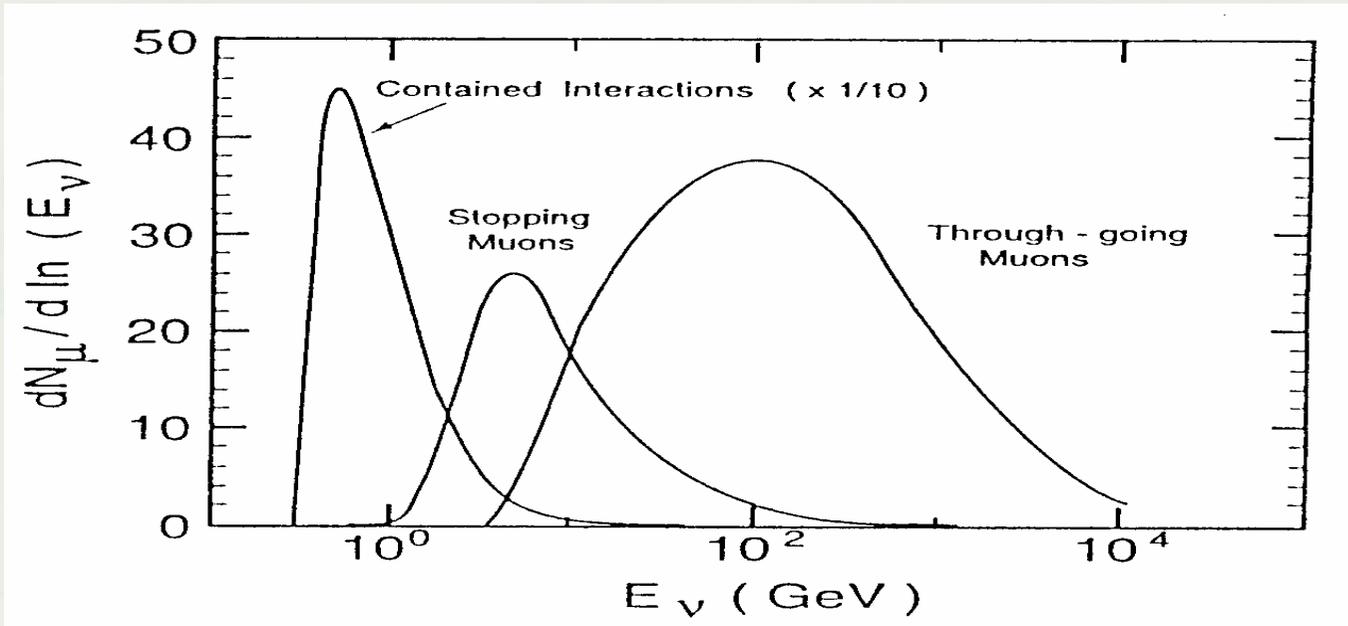
Solar $\rightarrow \theta_{12}$

ATMOSPHERIC NEUTRINOS



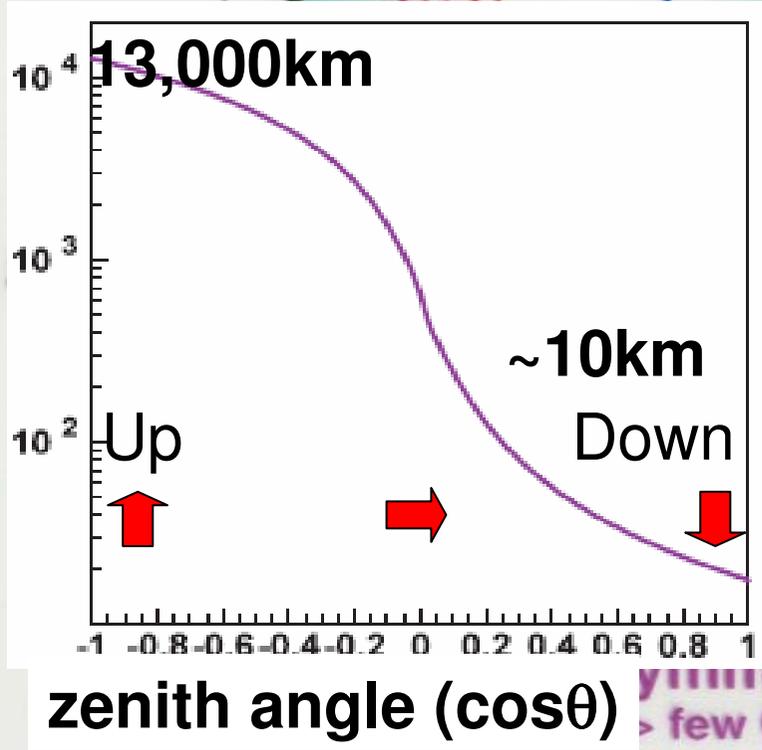
Ratio of $\nu_\mu/\nu_e \sim 2$
(for $E_\nu < \text{few GeV}$)

Up-Down Symmetric Flux
(for $E_\nu > \text{few GeV}$)

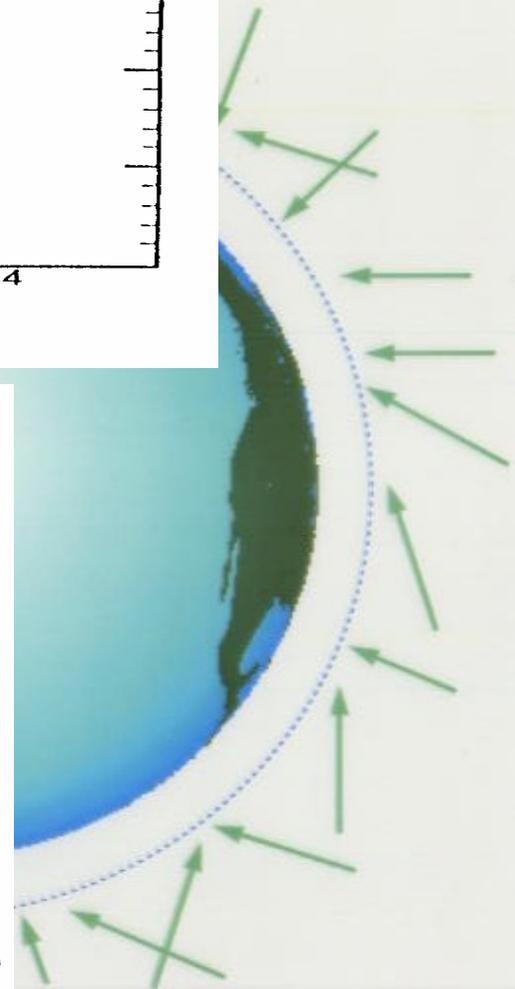


Ratio of $\nu_{\mu}/\bar{\nu}_{\mu}$ (for $E_{\nu} < \text{few GeV}$)

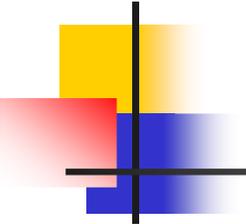
Flight path-length (km)



zenith angle ($\cos\theta$)



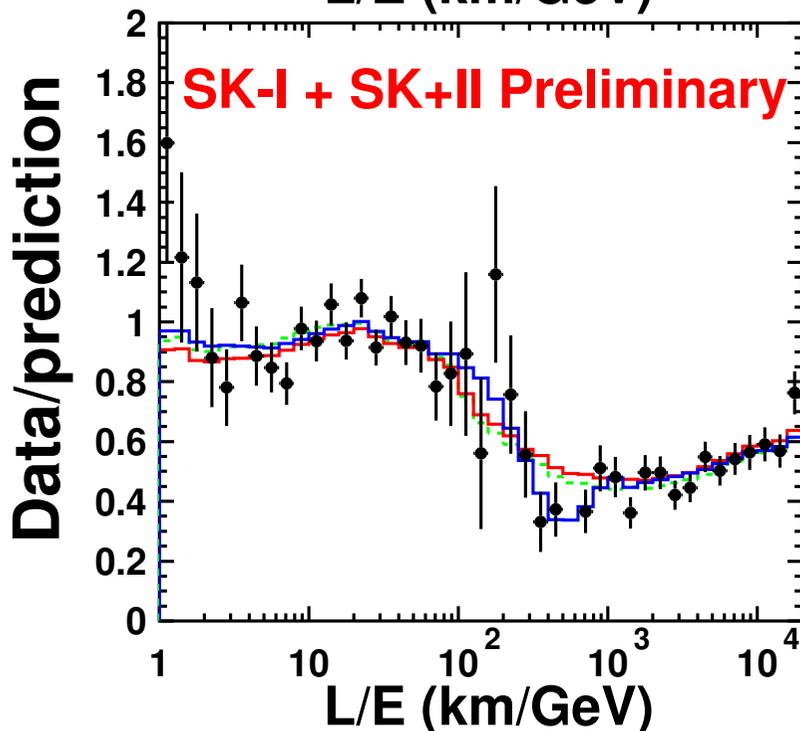
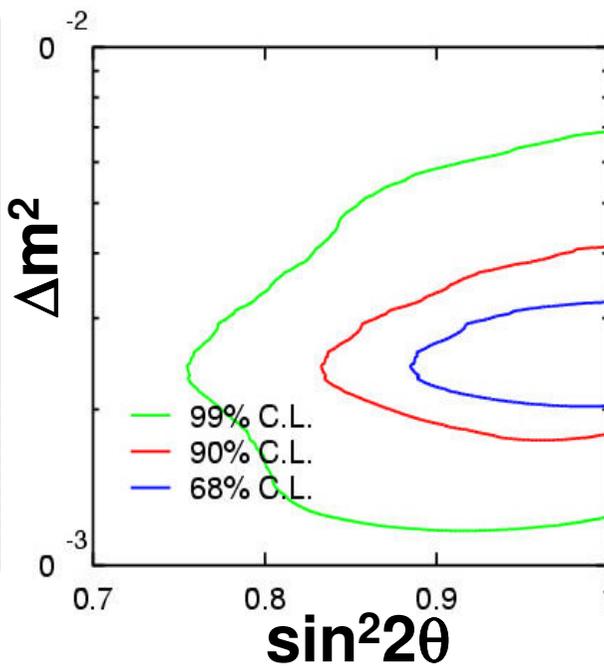
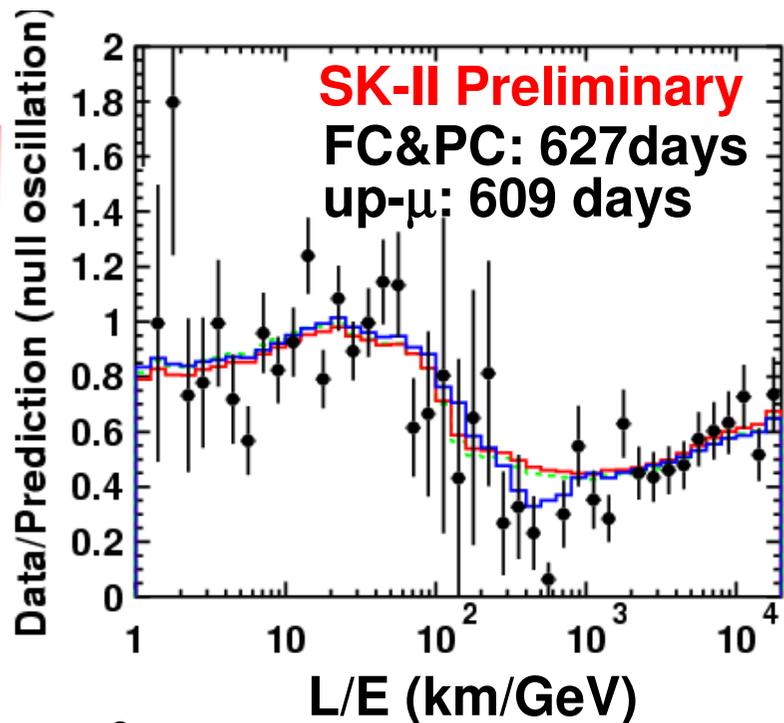
Symmetric Flux > few GeV



A bit of history

- In the '70 the most important problem in particle physics was the proton decay detection
- In this search the atmospheric neutrino interactions constituted the most tricky background: this is the reason why the study of atmospheric neutrinos started!
- The atmospheric neutrinos remained a "simple background" till when an anomalous results was obtained with Cerenkov detectors (Kamiokande, IMB) and lately confirmed with calorimeters (Soudan2)

Since then atmospheric neutrinos became the "high-way" towards new physics beyond the Standard Model



Combined results from SK-I and SK-II will be shown.

Need correct treatment of different systematics of SK-I and II .

Guide Line χ^2

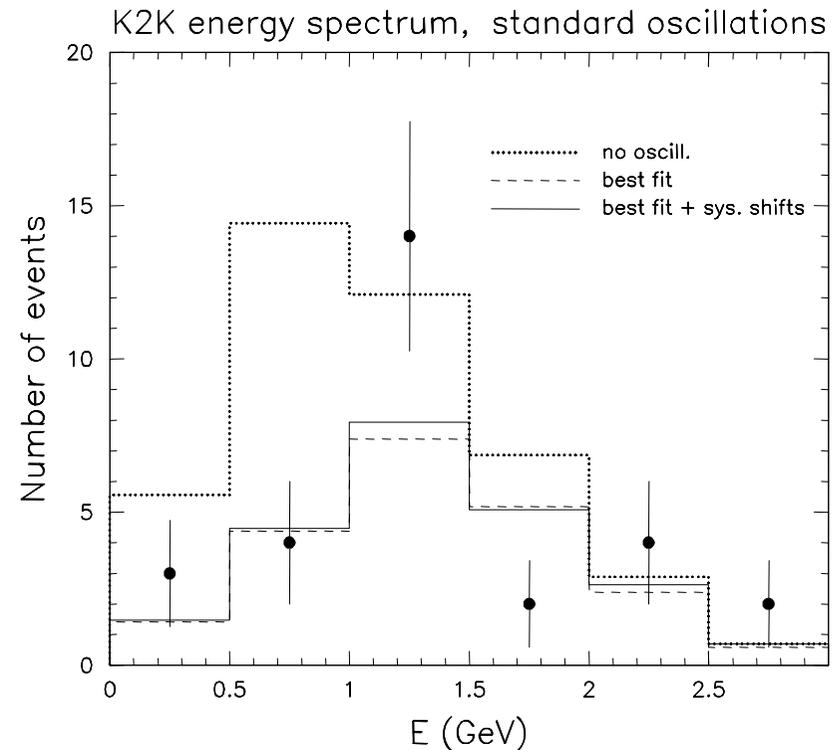
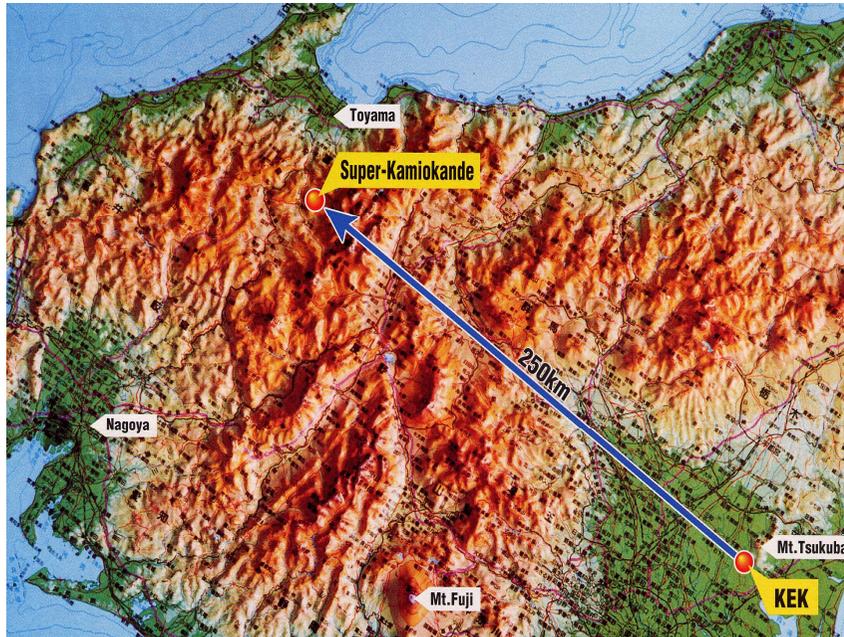
for $\Delta m^2 = 2.4 \times 10^{-3} \text{ eV}^2$, $\sin^2 2\theta = 1.02$

$\chi^2_{\text{osc}} = 42.9/42 \text{ d.o.f. (43\%)}$

neutrino decay $\Delta\chi^2 = 16.5 (4.1\sigma)$

de-coherence $\Delta\chi^2 = 20.9 (4.6\sigma)$

First-generation LBL accelerator experiment: KEK-to-Kamioka (K2K)



Aimed at testing disappearance of
accelerator ν_μ in the same range
probed by atmospheric ν :

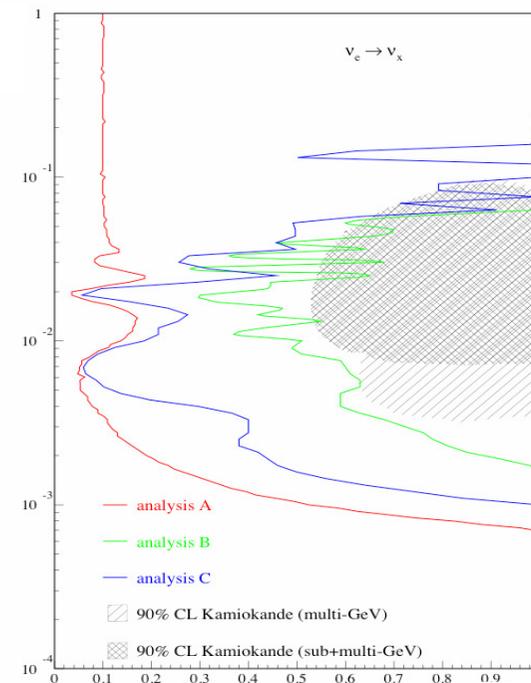
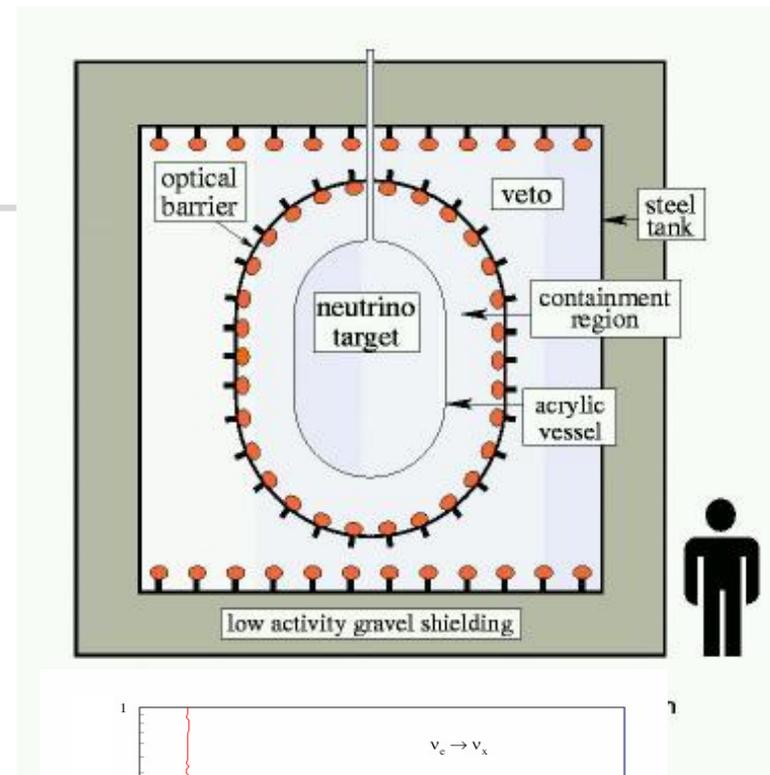
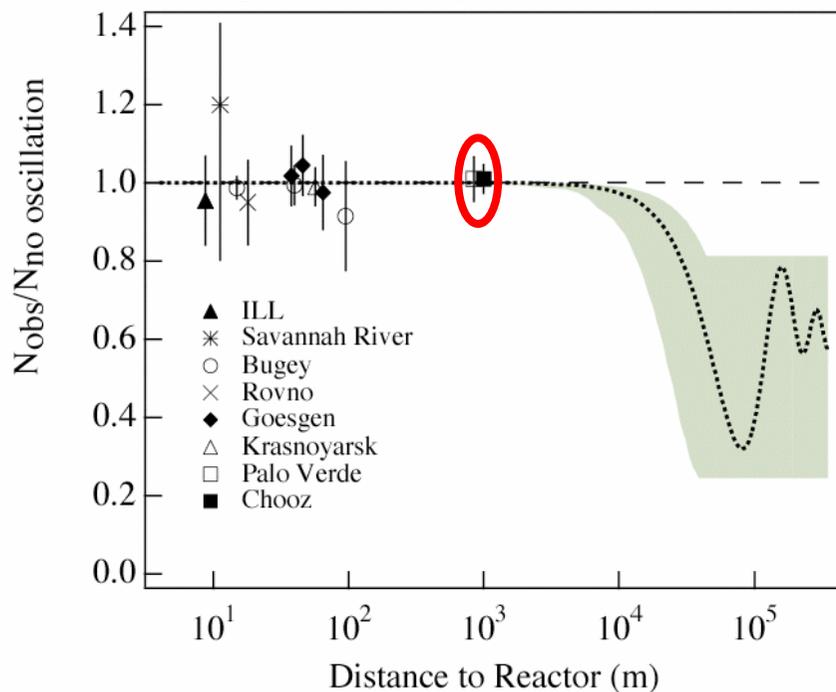
$$(L/E)_{K2K} \sim (250 \text{ km}/1.3 \text{ GeV}) \sim (L/E)_{ATM}$$

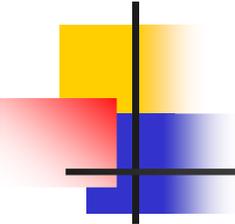
**2002: muon disappearance
observed at >99% C.L.**

No electron appearance.

The CHOOZ reactor experiment and θ_{13}

- Searched for disappearance of reactor ν_e ($E \sim \text{few MeV}$) at distance $L=1 \text{ km}$
- L/E range comparable to atmospheric ν
→ probe the same Δm^2
- No disappearance signal was found (1998)
→ Exclusion plot in $(\Delta m^2, \theta_{13})$ plane
- Results also confirmed by later reactor experiment (Palo Verde)





Numerical $\pm 2\sigma$ ranges (95% CL for 1dof), 2004 data:

$$\delta m^2 \simeq 8.0_{-0.7}^{+0.8} \times 10^{-5} \text{ eV}^2$$

$$\Delta m^2 \simeq 2.4_{-0.6}^{+0.5} \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{12} \simeq 0.29_{-0.04}^{+0.05} \quad (\text{SNO '05 : } 0.29 \rightarrow 0.31)$$

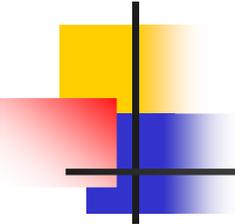
$$\sin^2 \theta_{23} \simeq 0.45_{-0.11}^{+0.18}$$

$$\sin^2 \theta_{13} < \sim 0.035$$

$\text{sign}(\pm \Delta m^2)$: unknown

CP phase δ : unknown

Note: Precise values for θ_{12} and θ_{23} relevant for model building



Probing absolute ν masses through non-oscillation searches

Three main tools (m_β , $m_{\beta\beta}$, Σ):

- 1) **β decay**: $m_i^2 \neq 0$ can affect spectrum endpoint. Sensitive to the "effective electron neutrino mass":

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{\frac{1}{2}}$$

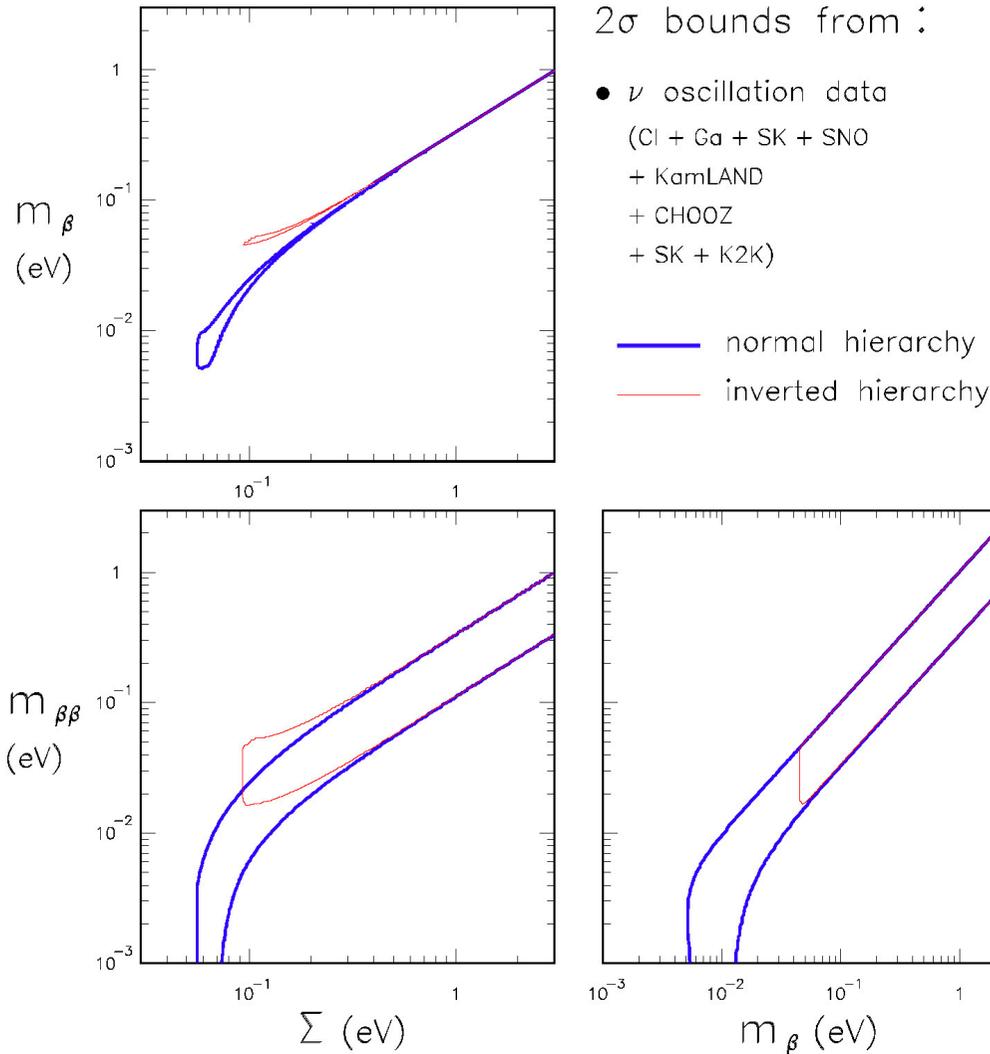
- 2) **$0\nu 2\beta$ decay**: Can occur if $m_i^2 \neq 0$ and $\nu = \bar{\nu}$. Sensitive to the "effective Majorana mass" (and phases):

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

- 3) **Cosmology**: $m_i^2 \neq 0$ can affect large scale structures in (standard) cosmology constrained by CMB+other data. Probes:

$$\Sigma = m_1 + m_2 + m_3$$

Even without non-oscillation data, the $(m_\beta, m_{\beta\beta}, \Sigma)$ parameter space is constrained by previous oscillation results



Partial overlap between the two hierarchies

Large $m_{\beta\beta}$ spread due to unknown Majorana phases

But we do have information from non-oscillation experiments:

1) β decay: no signal so far. Mainz & Troitsk expts: $m_\beta < O(\text{eV})$

2) $0\nu 2\beta$ decay, no signal in all experiment, except in the most sensitive one (*Heidelberg-Moscow*). Rather debated claim.

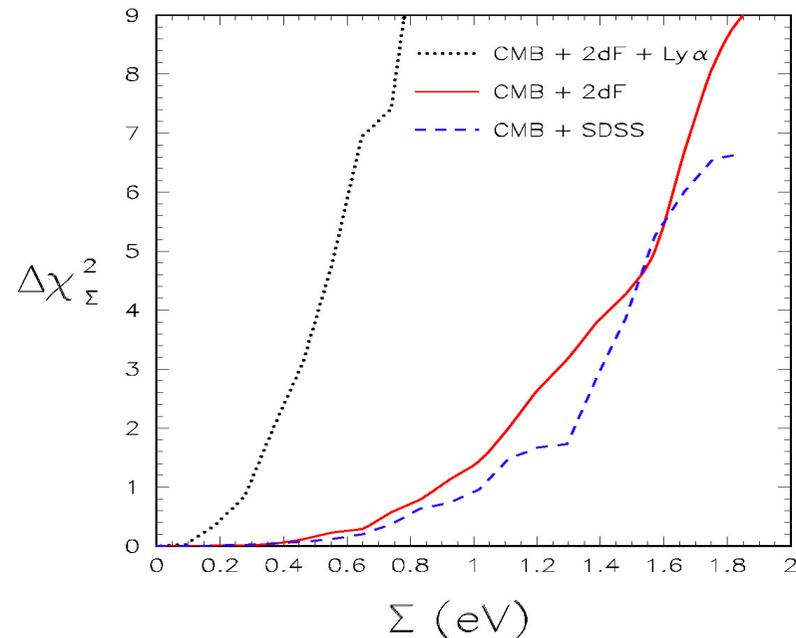
Claim accepted: $m_{\beta\beta}$ in sub-eV range (with large uncertainties)

Claim rejected: $m_{\beta\beta} < O(\text{eV})$.

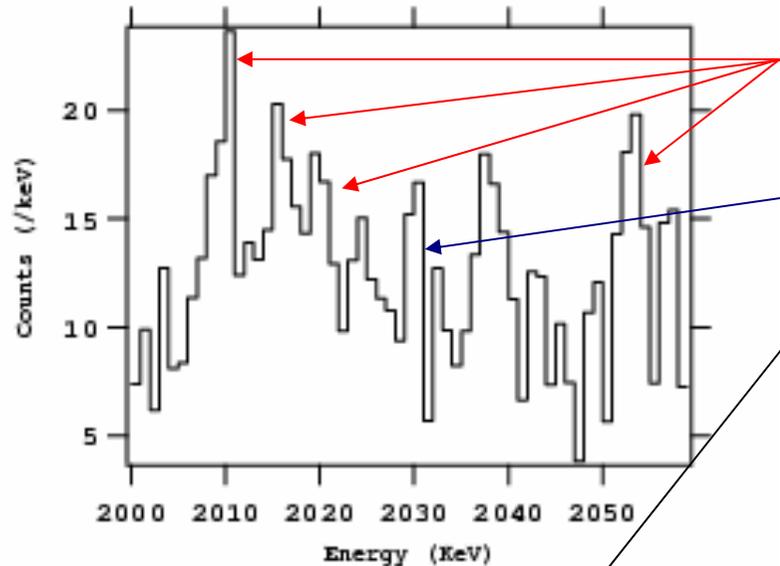
3) Cosmology. Upper bounds:

$\Sigma < \text{eV/sub-eV range}$,

depending on several inputs and priors. E.g.,



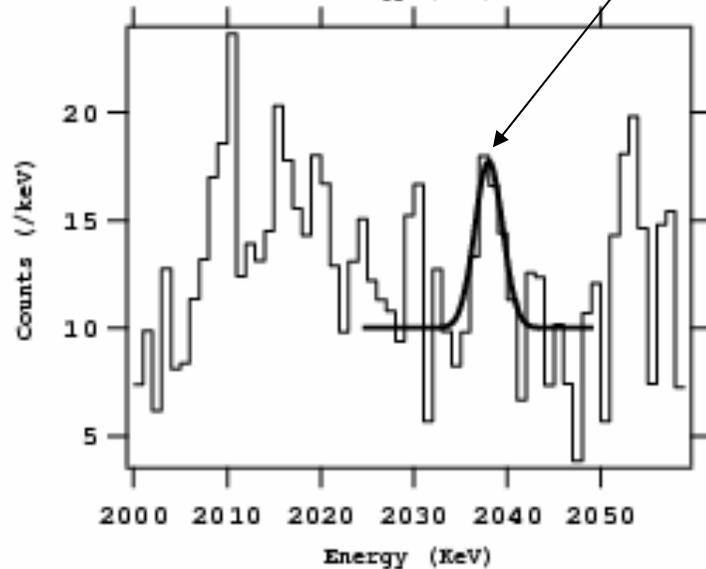
$0\nu 2\beta$ decay: Heidelberg-Moscow experiment final analysis (March 2004)



Four lines at 2010, 2017, 2022, 2053 keV are identified as due to ^{214}Bi decay

One possible line at 2030 keV is not identified

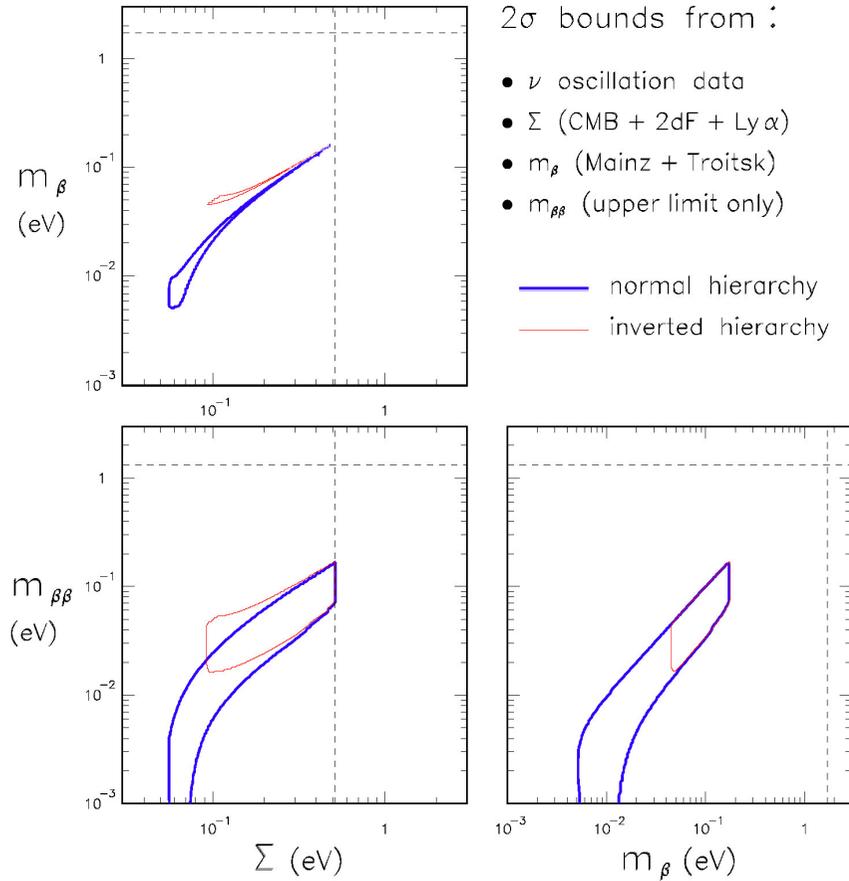
Claimed $0\nu\beta\beta$ line at ~ 2039 keV is now more clearly seen "by eye". Statistically, it emerges at about 4σ C.L. (~ 23 events)



We might have reached an "LSND-like" situation:

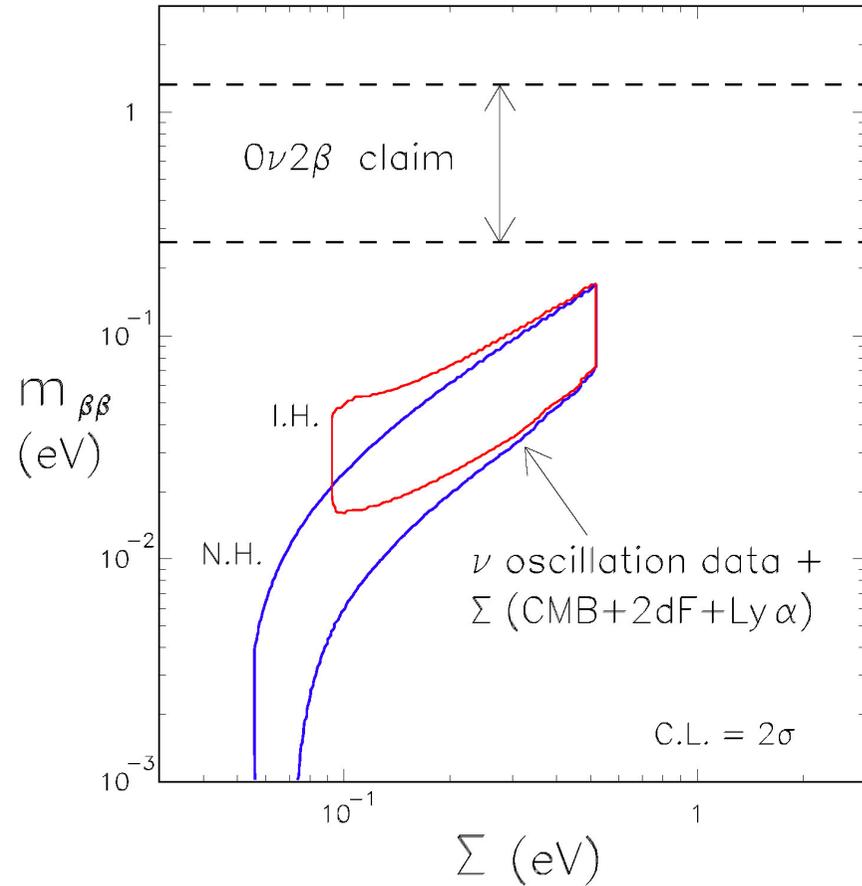
- Initial claim is rather controversial
- Then, further data/analysis strengthen it
- No current experiment can disprove it
- It will stay with us for a long time and will demand more sensitive expt. checks

$0\nu 2\beta$ claim rejected



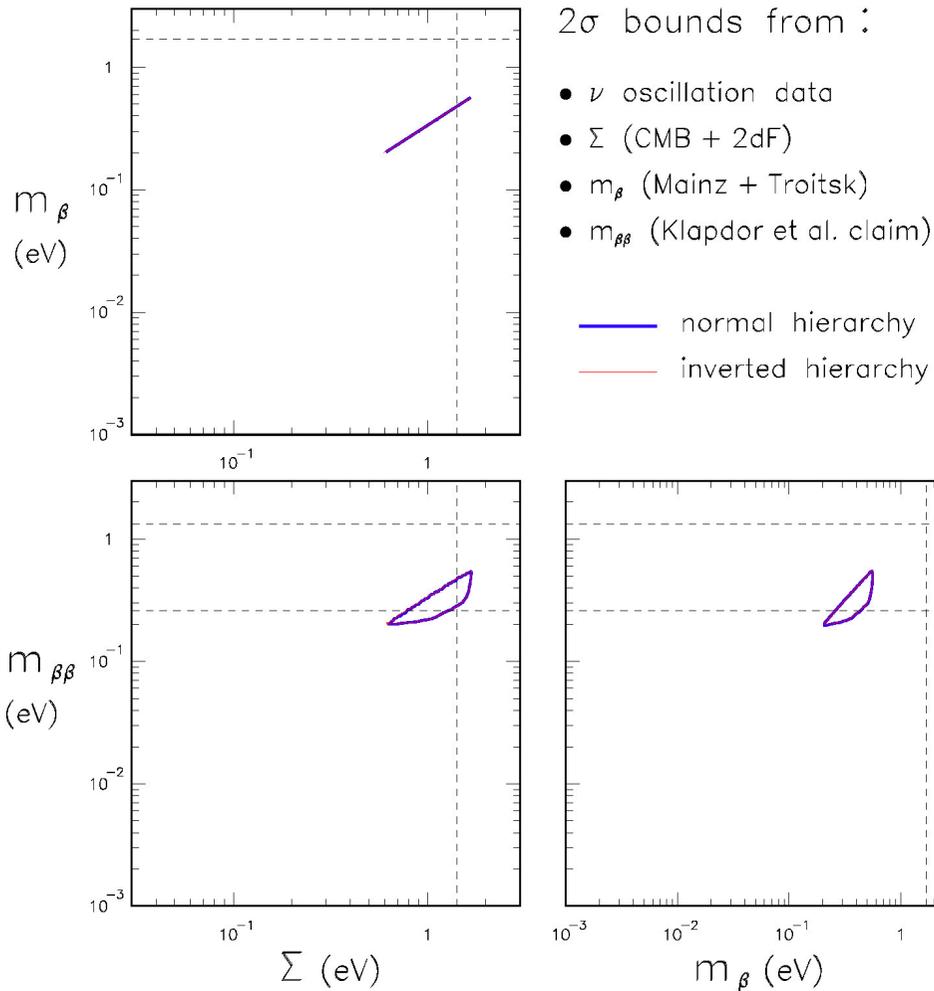
Cosmological bound dominates, but does not probe hierarchy yet

$0\nu 2\beta$ claim accepted



Tension with cosmological bound
(no combination possible at face value)
But: too early to draw definite conclusions

E.g., if $0\nu 2\beta$ claim accepted & cosmological bounds relaxed:



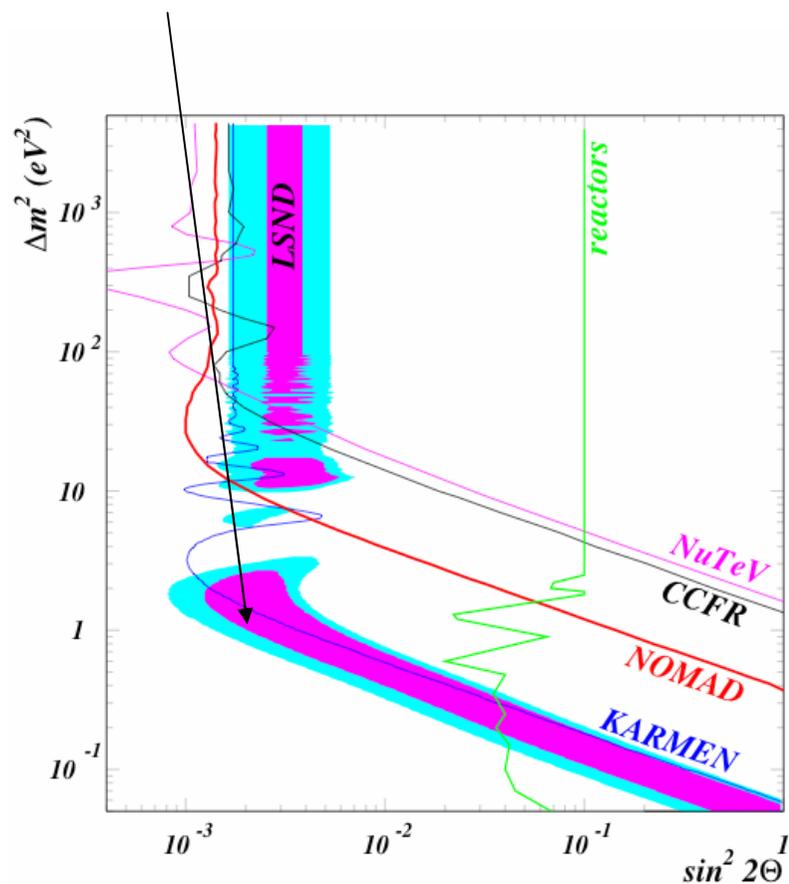
Combination of all data
(osc+nonosc.) possible

Complete overlap of
the two hierarchies
(degenerate spectrum
with "large" masses)

High discovery potential
in future ($m_\beta, m_{\beta\beta}, \Sigma$)
searches

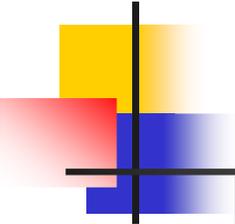
Beyond three-neutrino mixing: LSND

Many theoretical reasons to go beyond the standard 3ν scenario
A purely experimental reason: the puzzling LSND oscillation claim
 $\Delta M^2 \sim O(eV^2)$ with very small mixing?



Solutions invented so far
(new sterile states, new
interactions or properties)
seem rather "ad hoc"
and/or in poor agreement
with world neutrino data

If MiniBoone confirms
LSND this year (2005),
many ideas will be revised,



Question raised by neutrino data I

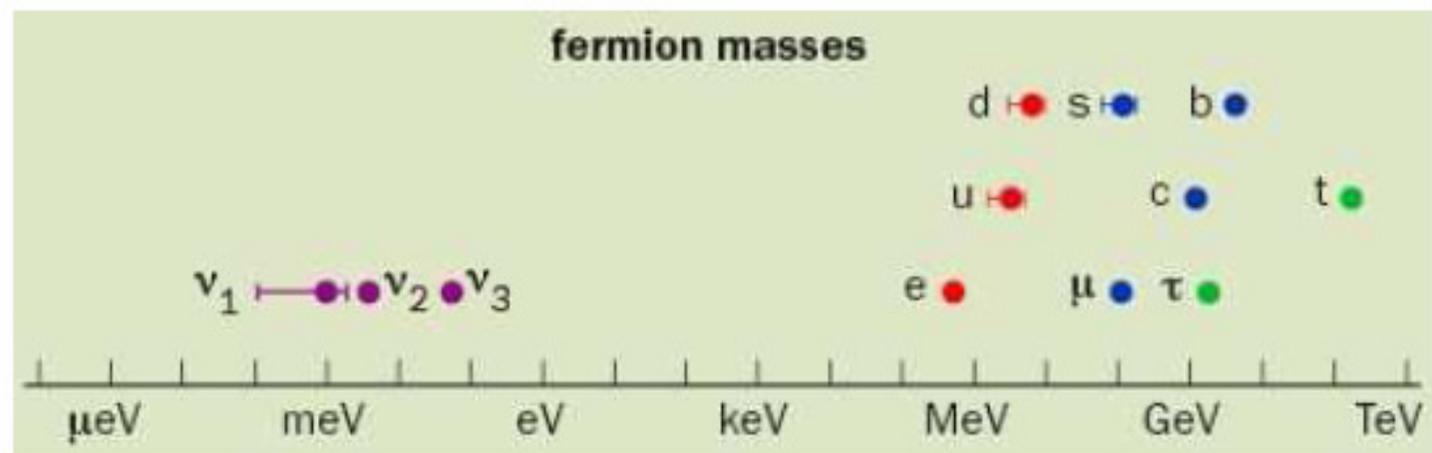
Quarks

$$U_{CKM} = \begin{pmatrix} 1 & 0.2 & 0.005 \\ 0.2 & 1 & 0.04 \\ 0.005 & 0.04 & 1 \end{pmatrix}$$

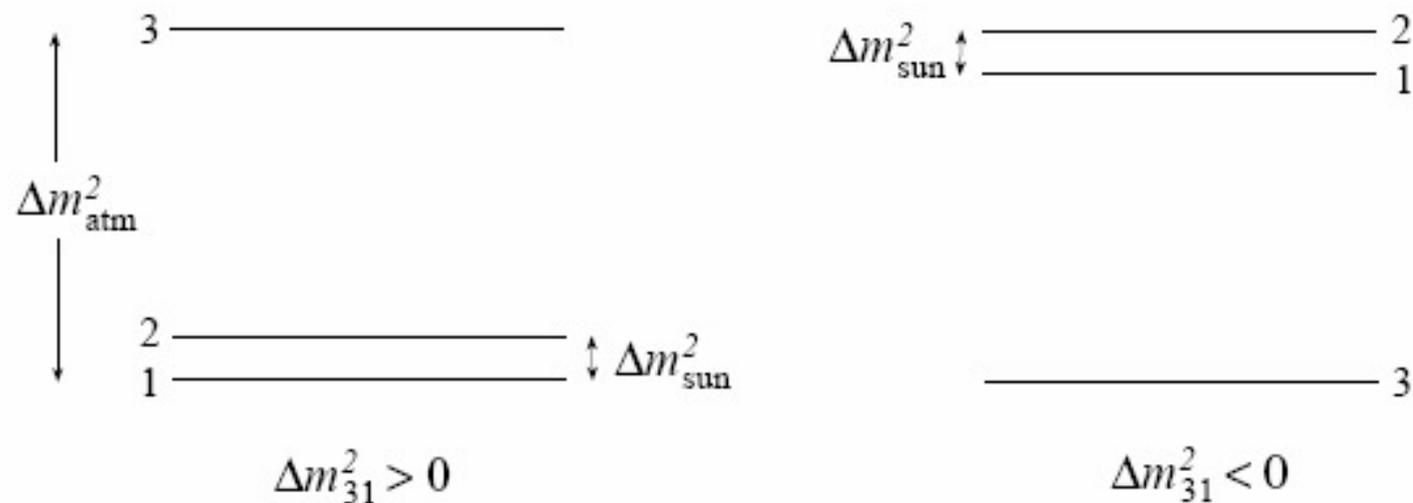
Why neutrino mixing are so large?

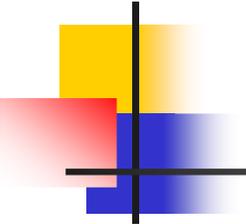
Question raised by neutrino data II

Mass hierarchy in the SM



What makes neutrinos so much lighter?

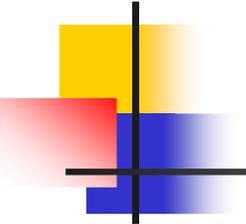




Question raised by neutrino data III

Origin of neutrino mass

- Neutrinos in the Standard Model (SM) are strictly massless, *ie.* there is no way to write a mass term for neutrinos with only SM fields which is gauge invariant and renormalizable
- Neutrinos are massive in reality - thus neutrino mass requires physics beyond the standard model
- One example of how to generate neutrino masses is the see saw mechanism: it introduces a heavy right handed neutrino (N_R), *ie.* a singlet under the SM gauge group



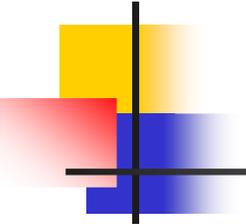
Question raised by neutrino data IV

Origin of baryons

- At the same time N_R can provide a mechanism for creating the observed tiny surplus of matter over anti-matter
- Leptogenesis requires the temperature of the Universe to be high enough that there is a thermal population of N_R . Their subsequent out-of-equilibrium decays are a new source of CP violation and lepton number

$$\Gamma(N_R \rightarrow LH) - \Gamma(N_R \rightarrow LH^*) \neq 0$$

which later on is converted to baryon number by nonperturbative processes

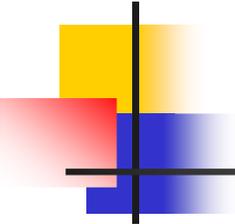


Key measurements

In the context of *GUT* scale right handed neutrinos it is very difficult to establish a one-to-one correspondence between high and low-energy observables.

A given model, however, usually has generic predictions for low energy observables. Therefore studying neutrinos allows to gain considerable insight into phenomena which otherwise would be inaccessible.

Neutrinos provide a unique window of observation on the *GUT* scale! And complementary to the energy frontier (e^+e^- and hadron colliders)



What still we have to observe or measure with higher precision

- The source of atmospheric oscillations (detect τ appearance)
- Three angles (θ_{12} , θ_{13} , θ_{23})
- Two mass squared differences (Δm^2_{12} , Δm^2_{23})
- The sign of the mass squared difference Δm^2 ($\pm \Delta m^2_{23}$)
- One CP phase (δ)
- The absolute mass scale
- Are neutrino Dirac or Majorana particles (or both)?
- Are there more - sterile - neutrinos?

All the underlined items can be
studied with LBL experiments

Discovery
Precision meas.

The accelerator/reactor based program

1st step: transition era

Ongoing: 2005-2010

- Improve the precision on the atmospheric parameters looking at ν_{μ} disappearance
- Confirm (atm. osc) = $(\nu_{\mu} \rightarrow \nu_{\tau})$ and first look at $\nu_{\mu} \rightarrow \nu_e$

2nd step: θ_{13} era

Approved/Proposed: 2008-2015

- Demonstrate visibility of sub-leading transitions:
 $\nu_{\mu} \rightarrow \nu_e, \nu_e \rightarrow \nu_e$
- Explore θ_{13} down to 2° (today $< 10^\circ$)

3rd step: precision era

To be prepared: 2015-2025

$\theta_{13} > 3^\circ$ ——— Known by 2011 ——— $\theta_{13} < 3^\circ$

- Existing facilities could reach it
- ... but with very small sensitivity to δ_{CP} and mass hierarchy

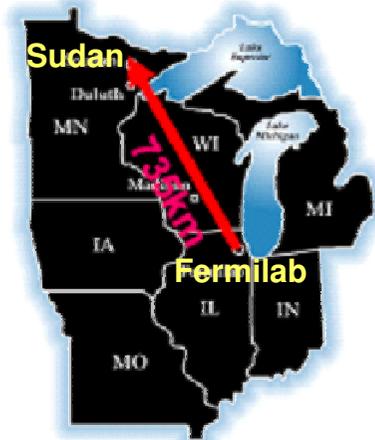
- No access for ongoing experiments at that time

Cleaner and more intense beams + bigger detectors

Transition era

- Conventional ν_μ beams from pion decay
- Long baseline experiments (such as K2K)
- Increased initial proton beam power: 0.01 (K2K) \rightarrow 0.4 MW

NUMI beam: *MINOS* (2005)



Magnetised
iron calorimeter

CNGS beam: *OPERA* (2006)

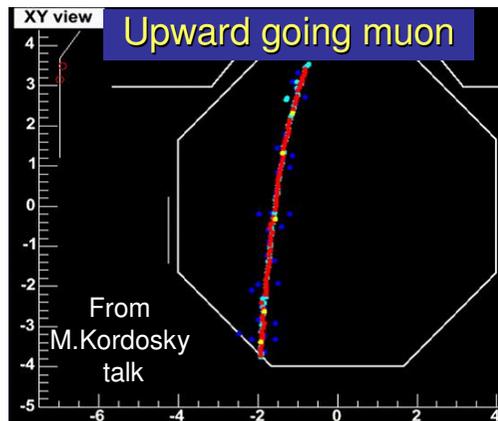
Hybrid emulsion
detector

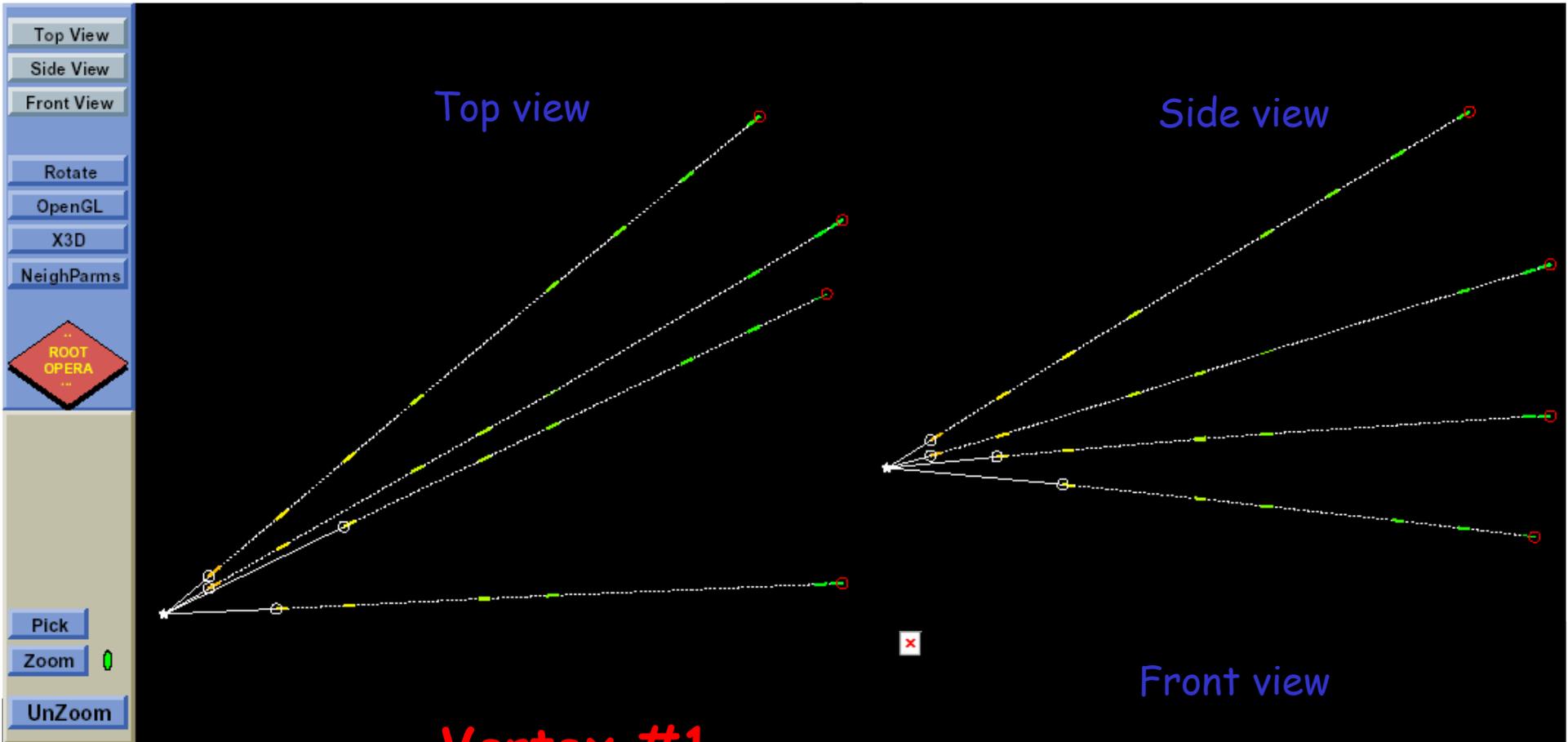


*Improve
atmospheric
parameters*

Confirm atm = $\nu_\mu \rightarrow \nu_\tau$

*First look at
 ν_e appearance*





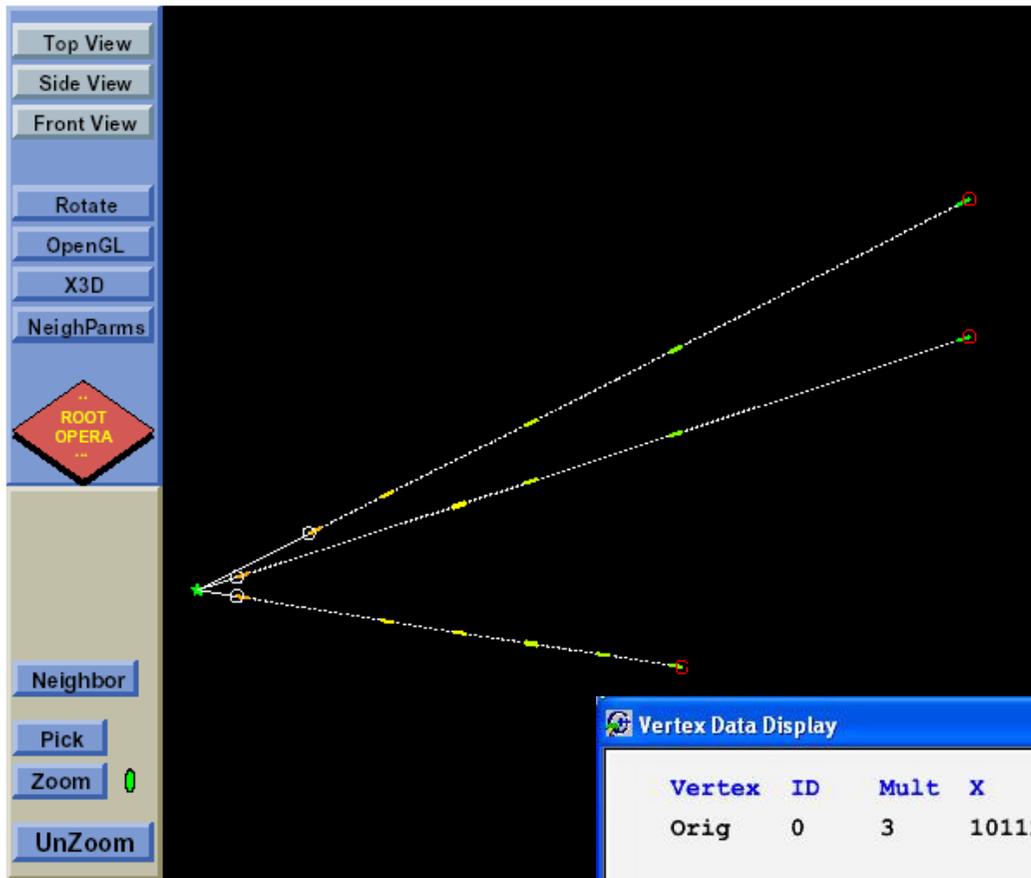
Vertex #1

Impact parameters: $[4,4,7,11] \mu\text{m}$



The "FIRST" neutrino interaction detected in an OPERA ECC has been found in Napoli Scanning Lab

D. Coppola, F. Di Capua, A. Marotta, C. Pistillo, L. Scotto Lavina, V. Tioukov



Vertex #2

Vertex Data Display

Vertex	ID	Mult	X	Y	Z	Dist	Chi2	Prob
Orig	0	3	10112.5	11555.9	-13879.3	5.9	0.2	0.92128

Tracks parameters for Original vertex

#	ID	Nseg	Mass	P	Chi2/ndf	Prob	Chi2Contrib	Impact	
0	986	5	0.1390	1.00	0.95	0.9970	0.000	1.23	Rem
1	1545	6	0.1390	1.00	0.71	0.9999	0.010	0.35	Rem
2	505	5	0.1390	1.00	2.03	0.8600	0.479	8.28	Rem

Vertex Data Display | VTX-Canvas-1 | x=0, y=-9.45823



θ_{13} era: Reactors

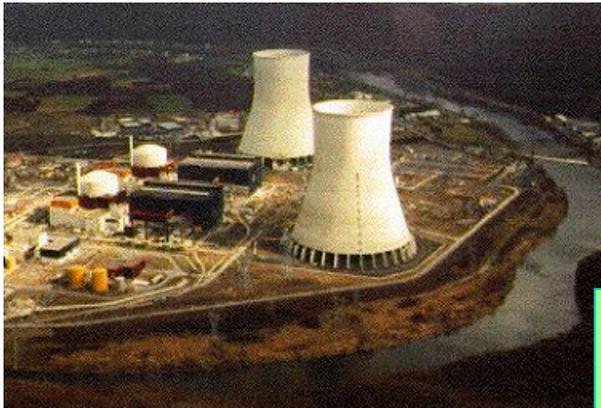
- High rate $\bar{\nu}_e$ by inverse beta decay
- Unambiguous determination of θ_{13}
- ... but cannot test mass hierarchy or CP violation

- Europe: Double-Chooz
- Others sites: Brazil, China, Japan, Russia, US, ...

Can new reactor experiments achieve the required low level of systematic errors ?



Double-Chooz (2008)



Collaboration

- France, Germany, USA, Russia
- Approved in France
- LOI's: hep-ex/0405032 & hep-ex/0410081
- <http://doublechooz.in2p3.fr>

*go down to $\theta_{13} \sim 4-5^\circ$
With $\bar{\nu}_e$ disappearance*



- Reduce systematic errors by a factor 5 with two identical detectors
- Still pending for full funding

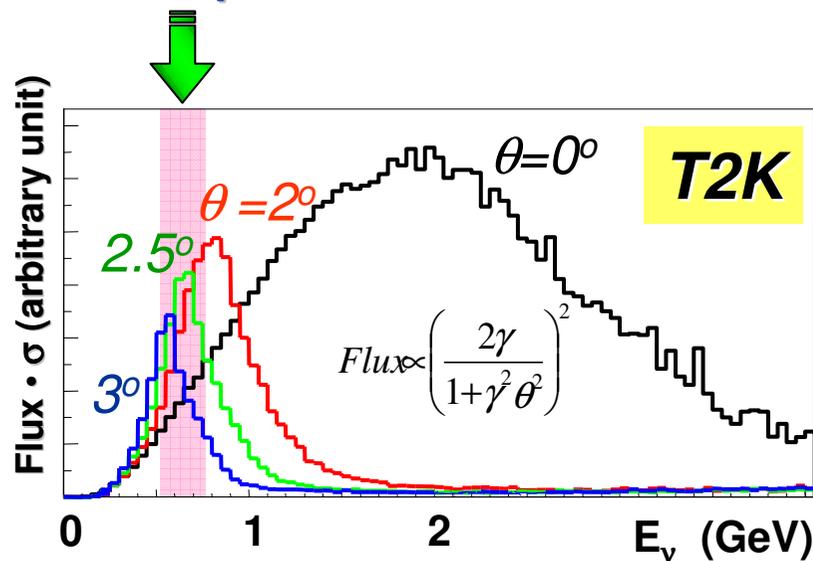
Chooz site (France)

- Agreement with EDF in 2005
- *Far site*: ready for integration (2007)
- *Near site*: 40 m shaft to build (2009)

θ_{13} era: Super-Beams

- Conventional ν_{μ} beams from pion decay
- Increased proton beam power: 0.4 \rightarrow 0.8 MW
- **Off-axis technique**: narrow band beam with purer composition
- **Tune L/E** to the oscillation maximum (L/E \sim 500 Km/GeV)

Oscillation peak at 295 Km



JPARC beam: T2K (2009)

- 0.4% ν_e
- L=295 Km
- Water Čerenkov (SK)

NuMI off-axis: NOvA (2011 ?)

- 0.5-1% ν_e
- L=810 Km
- Fully active calorimeter

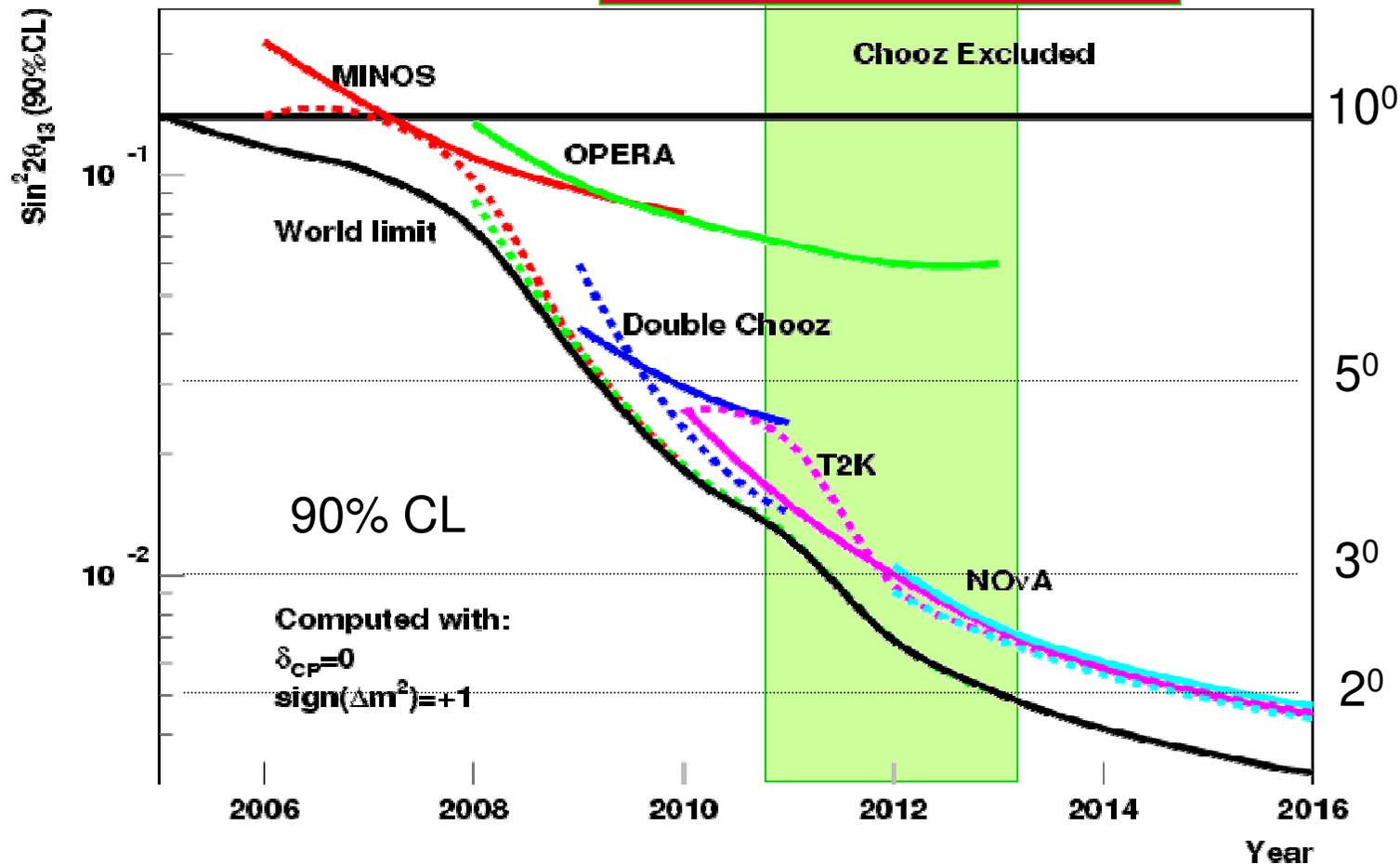
Further improve atmospheric parameters with ν_{μ} disappearance

go down to $\theta_{13} \sim 2-3^\circ$ with ν_e appearance

A first look at mass hierarchy (NOvA only)

Results for θ_{13}

Decision about 3rd step



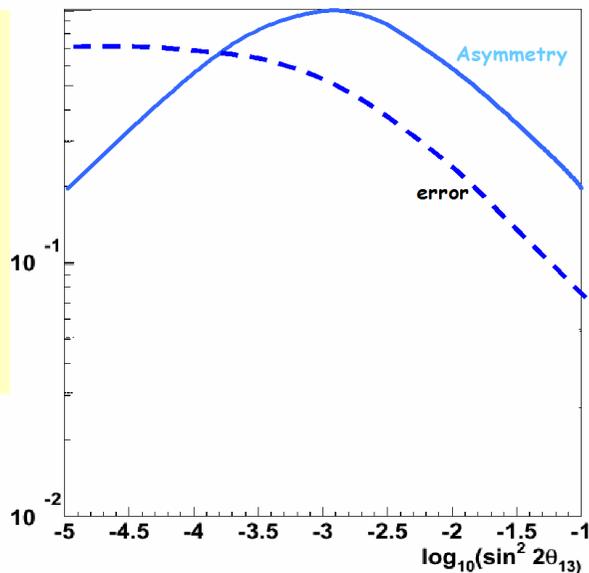
- If θ_{13} is not measured by ~ 2011 , the probability to measure it with ongoing experiments would be very small
- Building new facilities will take more than 5 years

CP violation and mass hierarchy

CP violating phase

$$A^{CP} = \frac{P(\nu_e \rightarrow \nu_\mu) - P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu)}{P(\nu_e \rightarrow \nu_\mu) + P(\bar{\nu}_e \rightarrow \bar{\nu}_\mu)} \propto \sin \delta_{CP} \cdot f(\theta_{13})$$

CP ASYMMETRY



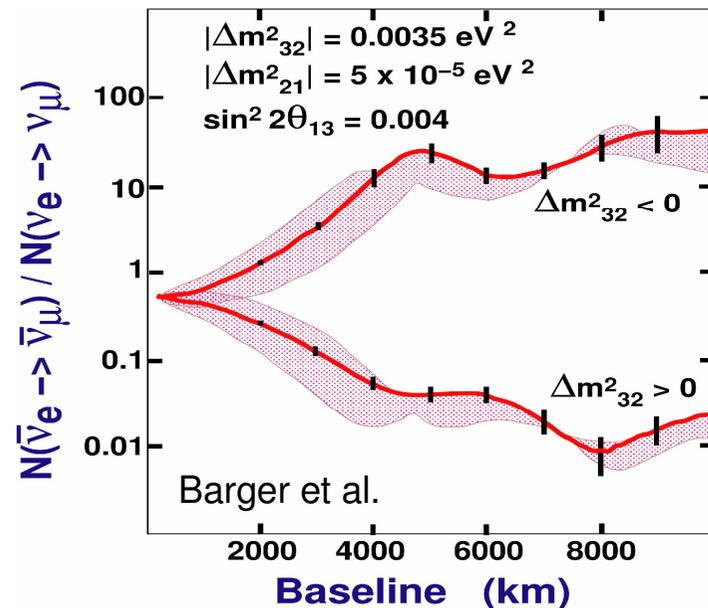
Difficulties

Correlations

- Several unknowns in the same Eq.
 - θ_{13}
 - δ_{CP}
 - Sign (Δm_{23}^2)

Mass hierarchy: sign(Δm_{23}^2)

- The oscillation probability depends on sign(Δm_{23}^2) through matter effects
- Sensitivity increases with L



Degeneracies

- Ambiguities due to lack of knowledge on:
 - δ_{CP}
 - Sign (Δm_{23}^2)
 - Octant: $\theta_{23} > \pi/4$ or $\theta_{23} < \pi/4$

Improved Super-beams

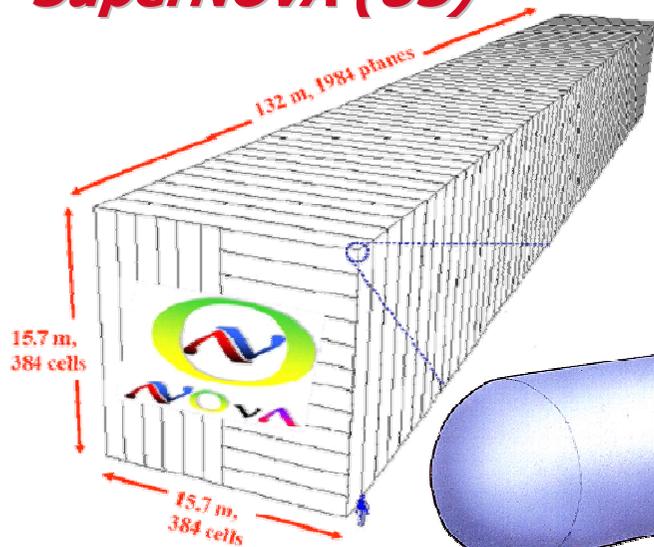
- Increase by one order of magnitude
 - beam power: $\sim 4\text{MW}$
 - detector mass
- Three proposals:

Systematics unchanged

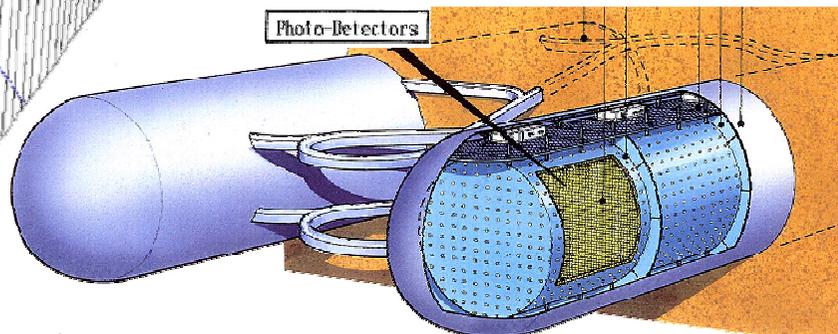
- Beam contamination
- Cross section
- Detector efficiency

T2HK (T2K-II)	Japan	0.6 GeV	295 Km	1000 KT Water Čerenkov
SPL-Memphys	Europe	0.25 GeV	130 Km	440 KT Water Čerenkov
NuMI-SuperNOvA	US	2 GeV	890 Km	130 KT fully active calorimeter

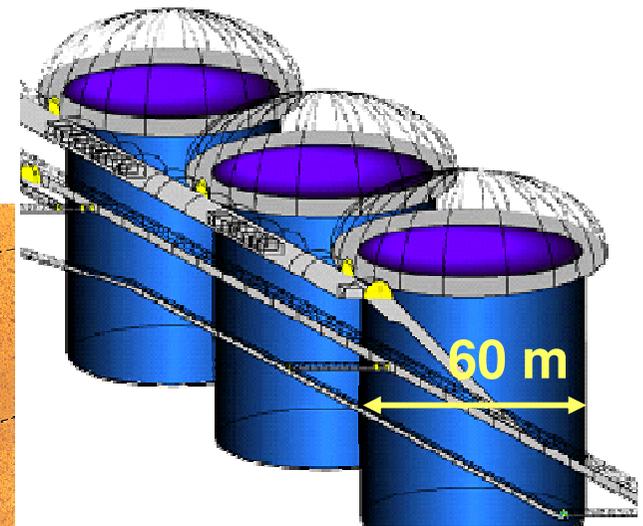
SuperNOvA (US)



Hyper-Kamiokande (Japan)

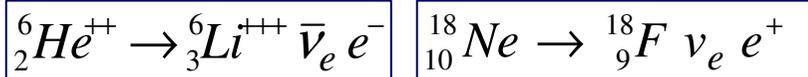


Memphys (Frejus)

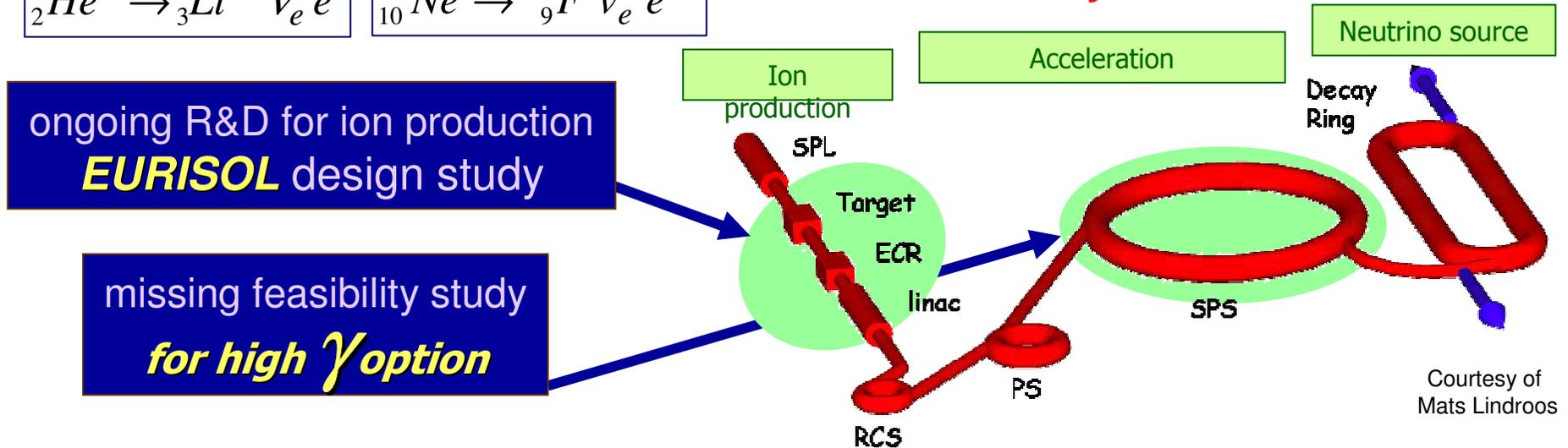


Beta-beam

Pure ν_e or $\bar{\nu}_e$ beam \rightarrow small beam systematics and backgrounds



CERN layout



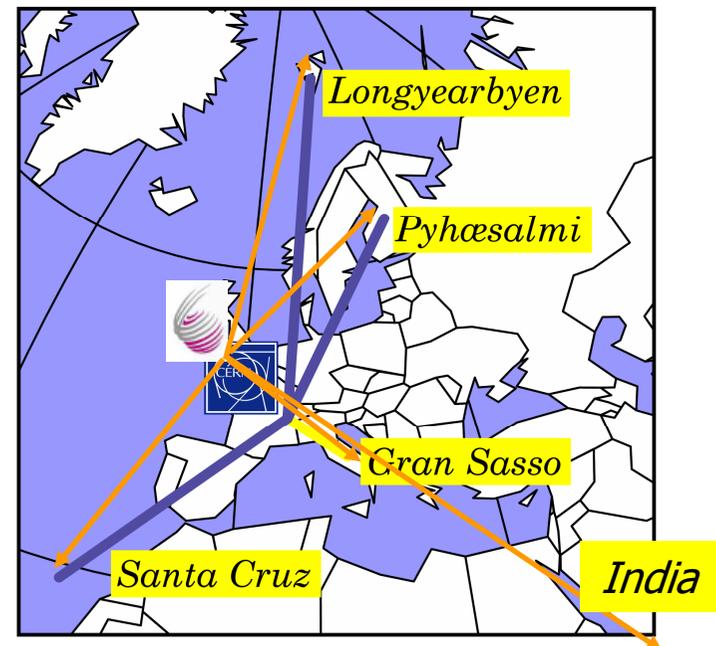
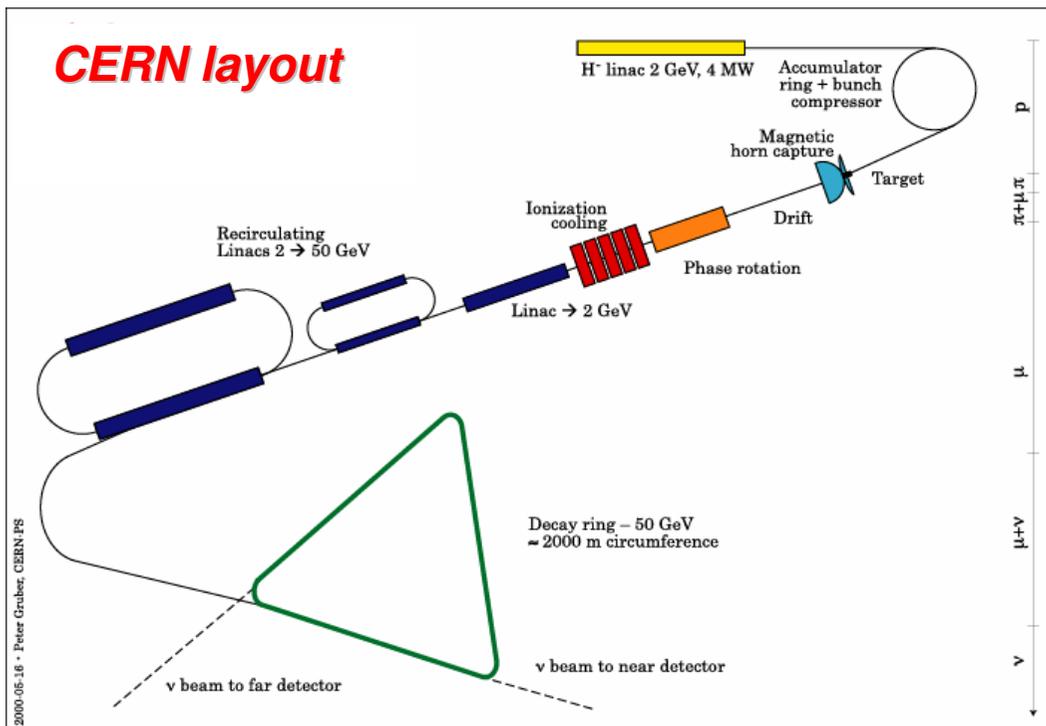
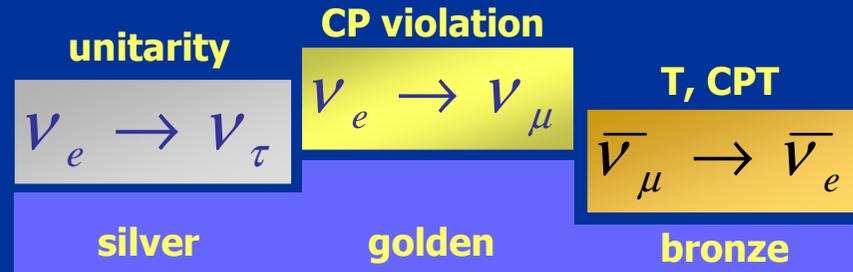
- Performance increases with beam energy if L/E is kept at oscillation max:
 - Higher flux and cross section. Better energy binning (no Fermi motion)
 - Smaller systematics from cross section and detector efficiency
- (Burget et al.)

Performance \uparrow	High γ	LHC	$\gamma \sim 1500$	7 GeV	3000 Km	0.1 MT TC	CERN-Canarias
		Tevatron or S-SPS	$\gamma \sim 350$	1.5 GeV	730 Km	1 MT WČ or TC	CERN-GS/Canfranc
	Low γ	SPS (max energy)	$\gamma \sim 150$	0.6 GeV	300 Km	1 MT WČ	?
		SPS	$\gamma \sim 100$	0.35 GeV	130 Km	1 MT WČ	CERN-Frejus

Neutrino factory

- 50% $\bar{\nu}_\mu$ 50% ν_e → small beam systematics ... but charge required
- High energy beam → small cross section systematics
- A wide variety of studies are possible:

and also: $\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$ $\bar{\nu}_\mu \rightarrow \bar{\nu}_\tau$
Atmospheric osc.



Detectors

High energy beams only: Nufact or high γ β -beam

Hybrid emulsion (4 KT)

- Experience from OPERA
- Silver channel
- Interesting to solve degeneracies
- Golden and bronze also

$$\nu_e \rightarrow \nu_\tau$$

CP asymmetry has opposite sign

Tracking Calorimeters (100 KT)

- Fully active with liquid scintillator: \sim NOvA
- Or sampling iron calorimeter: \sim MINOS
- Muon charge is crucial: B field !!!
- Golden channel

$$\nu_e \rightarrow \nu_\mu$$

Liquid Argon TPC (100 KT)

Both

- 3D active detector: Imaging, calorimetry, Čerenkov
- Challenging: ongoing R&D strategy

- GLACIER conceptual design
- ... also with magnetic field
- Could explore all channels

Low energy beam only:

- $\gamma < 350$ β -beam
- Super-beam

And also:

- Proton decay
- Supernovae neutrinos

Water Čerenkov (0.5-1 MT)

- Well known technique from Super-K
- Interesting for e/μ separation

$$\nu_\mu \rightarrow \nu_e$$

super-beam

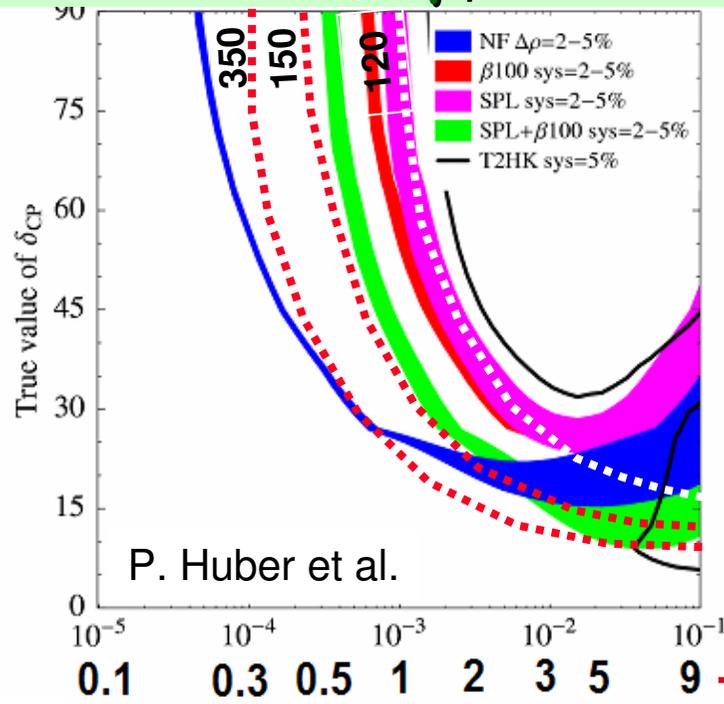
$$\nu_e \rightarrow \nu_\mu$$

β -beam

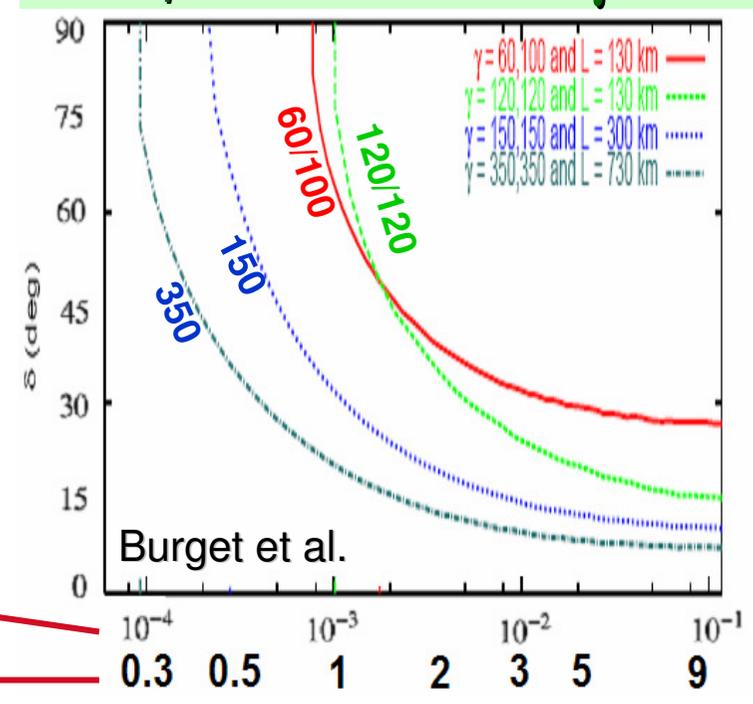
Neutrino Energy

Comparisons: $\delta_{CP}-\theta_{13}$

All except high- γ β -beam



β -beam for various γ

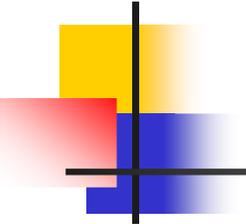


$\sin^2(2\theta_{13})$
 θ_{13} (°)

- The Nufact study is old (5 years). It should be revisited in order to make a fair comparison with β -beam
- *For large θ_{13}* systematic errors dominate. The picture is not clear yet: ongoing studies by several groups
- *For small θ_{13}* Nufact and High γ β -beam (350) outperforms all others
- The third option should be Low γ β -beam + Super-beam

Physics reach

	Systematic errors	E bins	B field	WČ	channels	Sensitivity to oscillation parameter $\theta_{13} < 3^\circ$ $\theta_{13} > 3^\circ$			R&D	Others
Nufact	Matter eff.	yes	yes	no	Golden Silver Bronze Atmos.	CP phase Mass hierarchy maximal θ_{23} ?	Good Good Good	Good Good Good	High	T violation Muon physics Muon collider
β -beam $\gamma > 1000$	Matter eff.	yes	no	no	Golden Silver	CP phase Mass hierarchy	Good Good	Good Good	Med. ?	
β -beam $\gamma \sim 350$		yes	no	yes	Golden Silver	CP phase Mass hierarchy	Good Small	Good Med.	Small	Supernovae Proton decay Atmos.
β -beam $\gamma < 150$	x-section Efficiency	poor	no	yes	Golden Silver	CP phase Mass hierarchy maximal θ_{23} ?	Med. None Med.	Good Small Med.	Small	Supernovae Proton decay Atmos.
Super beam	x-section Efficiency Beam	poor	no	yes	Bronze Atmos.					



Some open questions

Facilities

- *Neutrino Factory* and *high $\gamma\beta$ -beam* are the best options for osc. physics, but we do not understand yet all the elements, including their cost, feasibility and time scale
- *Low $\gamma\beta$ -beam* and improved *Super-Beams* are not separated options, they form a package. Their combined physics reach should be better understood in terms of neutrino fluxes and systematics uncertainties

Detectors

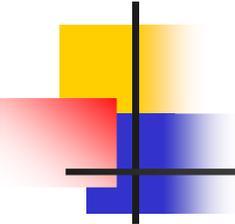
- The different options should be understood at the same level
- Detectors are not yet optimised for all possible θ_{13} values

Systematic errors

- For $\theta_{13} > 3^\circ$ systematic errors dominate for all scenarios
- Can we control them?

THE QUESTION

What is the best realistic scenario one could build in a reasonable time scale (10-15 years) to address CP violation and mass hierarchy ?



Outlook

Physics

- The CP violating phase and the mass hierarchy are crucial elements for the understanding of the leptonic sector
- Next generation neutrino facilities are required to assess these issues

Low γ β -beam +
Super-Beam +
Megaton detectors

High γ β -beam

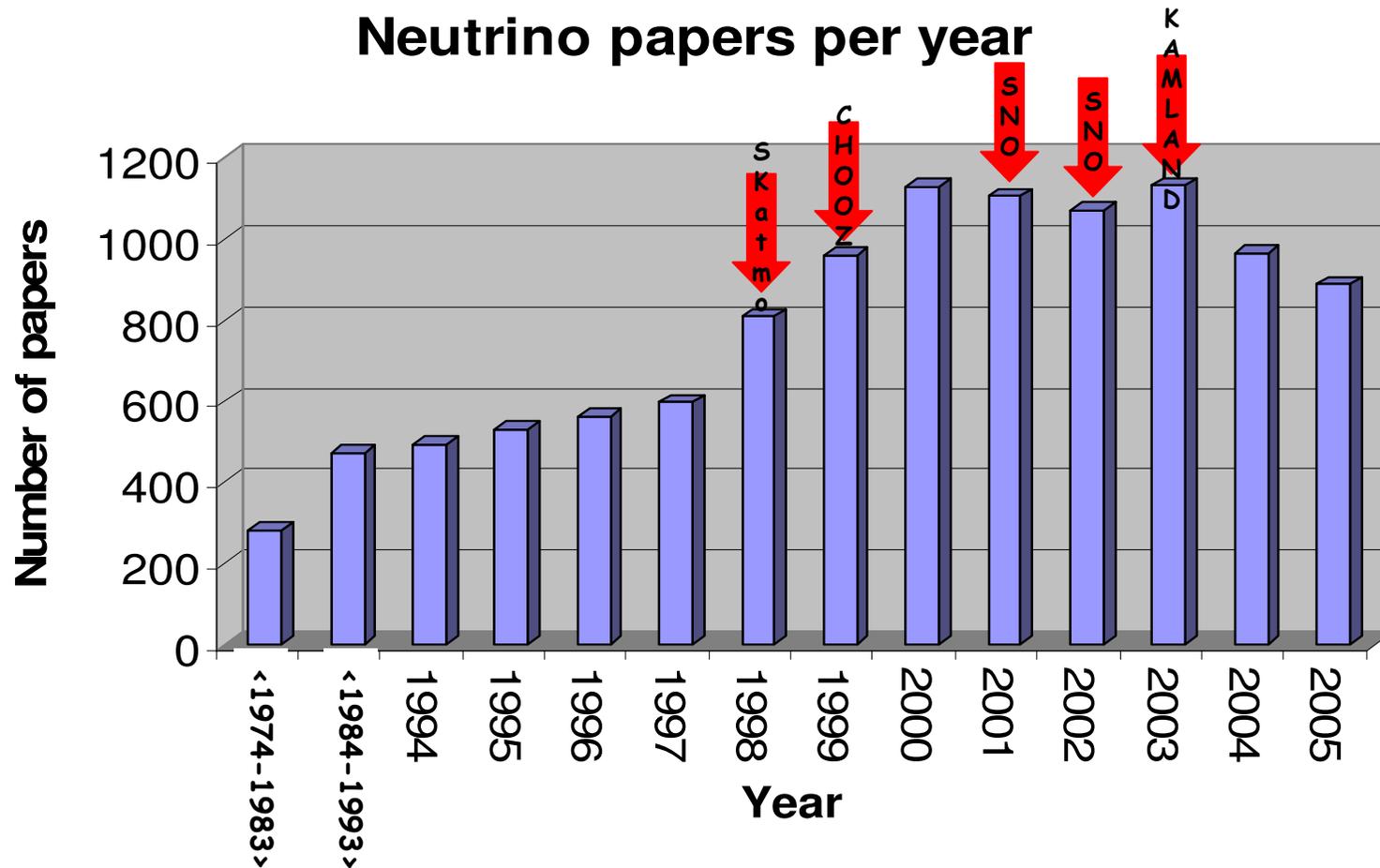
Neutrino Factory

Time scale

- We should enter the precision era in the second half of the next decade
- Meanwhile **priority is to perform an “International Design Study”**
 - Conceptual design and realistic cost estimate
 - Hardware R&D on accelerator and detectors
- ... in order to be able to compare cost, feasibility and performance
- ... to make the best choice by ~ 2011

\sim **G€**

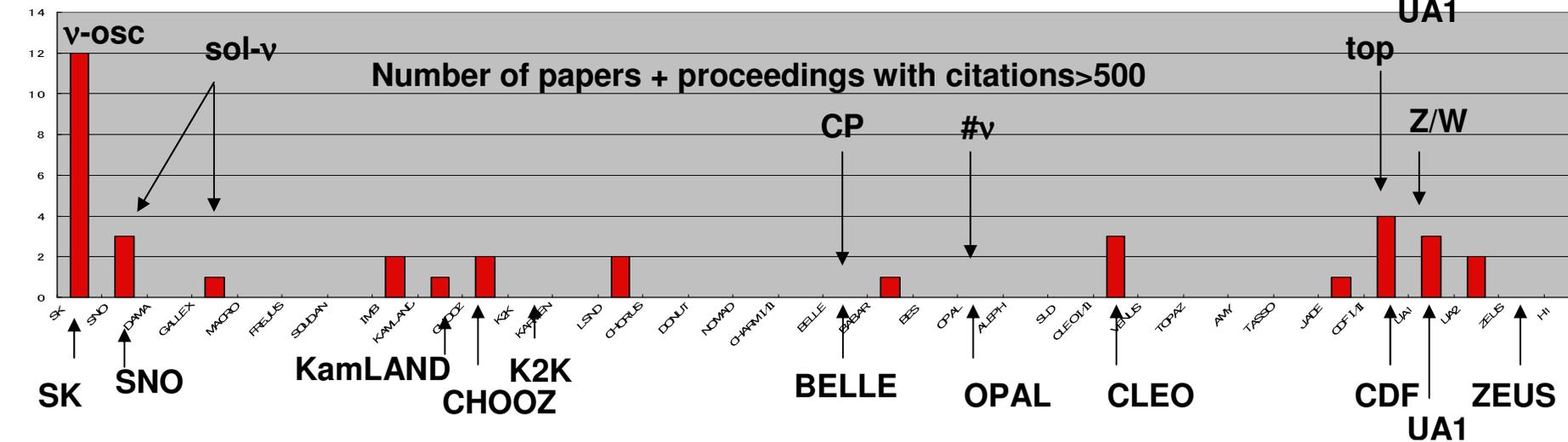
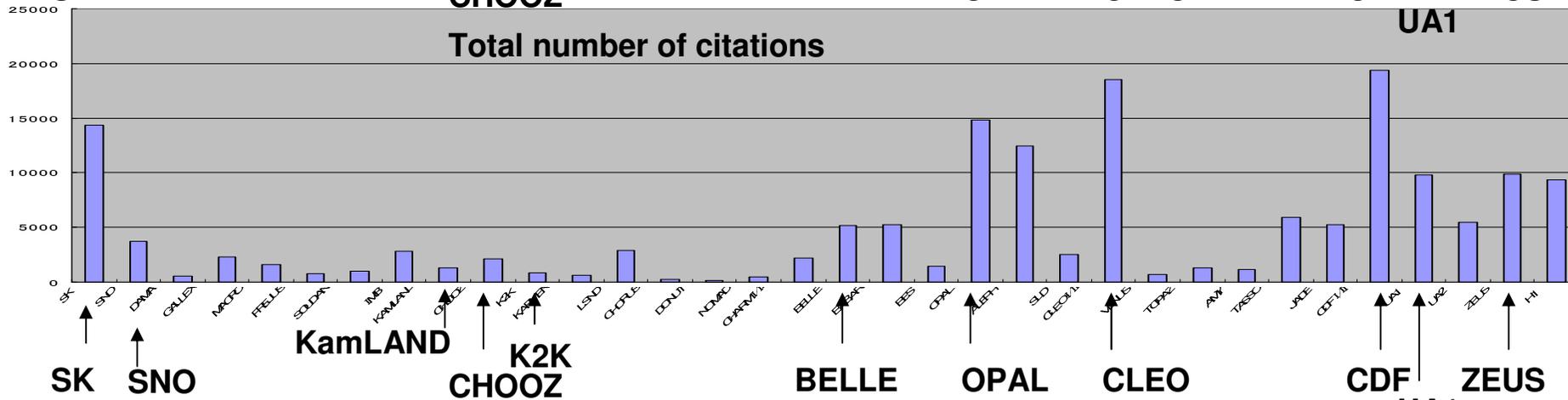
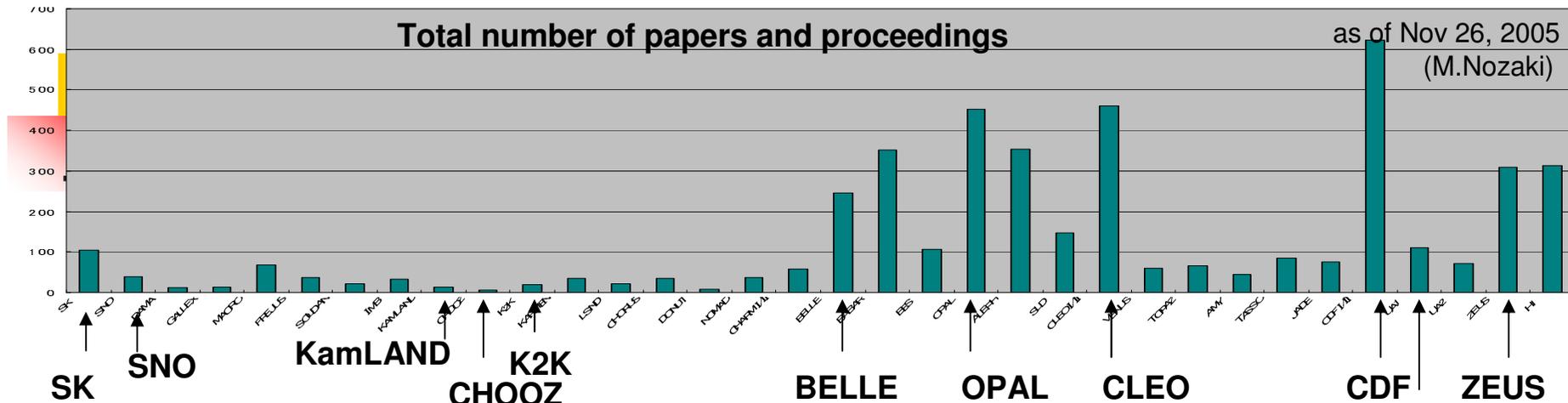
Neutrino papers per year

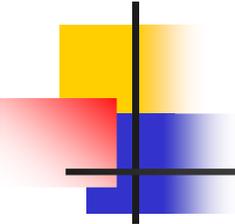


In HEP from 1997 on:

19 papers are topcite 1000+ \Rightarrow 6 (32%) involve neutrinos

73 papers are topcite 500+ \Rightarrow 16 (22%) involve neutrinos





Conclusion

All the places where we have looked for new physics we haven't found anything, but with neutrinos the first searches already were successful - it just took us a long time to believe it. Still, neutrinos are the least known of all fundamental Fermions and therefore even the most exotic things could be just around the corner